


Article

Numerical Prediction of Forming Car Body Parts with Emphasis on Springback

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Abstract: Numerical simulation is an important tool which can be used for designing parts and production processes. Springback prediction, with the use of numerical simulation, is essential for the reduction of tool try-outs through the design of the forming tools with die compensation, therefore, increasing the dimensional accuracy of stamped parts and reducing manufacturing costs. In this work, numerical simulation was used for performing the springback analysis of car body stamping made of aluminium alloy AA6451-T4. The finite element analysis (FEM) based software PAM-STAMP 2G was used for performing the forming and springback simulations. These predictions were conducted with various combinations of material models to achieve accurate springback prediction results. Six types of yield functions (Barlat89, Barlat2000, Vegter-Lite, Hill90, Hill48 isotropic, and Hill48 orthotropic) were used in combination with the Voce hardening model. Springback analysis was conducted in three sections of the formed part; the numerical results were compared with the experimental values. It was found that the combinations of Barlat's yield functions and the Voce hardening law were most accurate in terms of springback prediction. Additionally, it was found that the phenomena that were investigated, which are required for the determination of the kinematic hardening model, such as the change of Young's modulus E , the transient behaviour, work-hardening stagnation, and permanent softening, were not observed in the aluminium alloy studied.

Keywords: springback; numerical simulation; yield function; aluminium alloy formability

1. Introduction

Automobile manufacturers have started to use new types of high strength steels (HSS, AHSS, and UHSS) at the end of the last century, with the aim of increasing the passive safety of vehicles and to reduce vehicle weight to decrease fuel consumption [1–3]. However, these types of steels have a lower formability in comparison with steels used for deep drawing. The main reason for this is the higher values of the yield strength and lower ductility of high strength steels [2]. In addition, aluminium alloys are now widely used in the automotive industry due to advantages, including the low density, high specific strength, good corrosion resistance, exceptional specific stiffness, and so forth [1]. The implementation of aluminium alloys in car body production can reduce fuel consumption and emissions [4]. Both high strength steels and aluminium alloys are more prone to wrinkling and springback than mild steels [1,5].

Springback in the present refers to a change of shape which is elastically driven. Springback occurs following a sheet-forming operation when the forming loads are removed from the workpiece—sheet metal blank. It is usually unwanted, causing problems in the next forming operations, in assembly,

and in the final product. These problems usually degrade the accuracy, appearance, and quality of the products being manufactured [2,3,6]. The most common counter measurement against the springback of car body parts is to design a forming tool with anticipation of springback, thus compensating springback by die design. However, the amount of compensation is a difficult question even for skilled tool designers. In practice, this compensation of die is still sometimes done by the “trial and error” method. This method can be replaced by FEA (finite element analysis)—numerical simulation. With the use of FEA, it is possible to achieve a more accurate prediction of springback [6–8]. There are other counter measurements against springback, for example, the stiffening of pressings (use of beads or embossing), crash forming with pressure pads, the use of variable blank holder force, and so forth [7].

In general, two types of methods are used for springback prediction—finite element analysis and the analytical model. For example, the analytical model for springback prediction of aluminium alloys can be found in the work by Gau and Kinzel [9]. Analytical methods usually use simplified models of real processes. Thus, analytical models are usually not as accurate in predicting springback as numerical simulations and their use is limited, especially for stampings with complex geometry [10]. The finite element method (FEM) is a well-known tool for the prediction and analysis of sheet metal deformation. Springback prediction with the use of numerical simulation is not limited by the geometrical complexity of the stamped part like in the case of the analytical model. However, the numerical simulation of springback is more sensitive to the accuracy of the input data than the analytical method. Thus, it is very important to choose the correct input and numerical parameters in the FEA analysis of springback [11].

Numerical parameters involve the through-thickness integration scheme (which can be implicit, explicit, or a combination of both), the number of integration points, the used elements (type, size, and count), and so forth. Trzepieciniski and Lemu [12] studied the effect of a number of integration points and integration rules on the springback amount. Their results indicate that at least 5 integration points must be used to achieve accurate springback prediction. The input parameters involve geometry (sheet thickness, tool and sheet dimensions, and so forth), process conditions (type of forming method, tribology, forming forces, forming temperature and speed, and so forth), and material characteristics (Young’s modulus, yield strength, hardening behaviour, yield function, and so forth) [10,11]. Slota, Siser, and Dvorak [13] studied the effects of yield functions (isotropic and orthotropic) on the springback prediction accuracy of aluminium alloys. Their results showed that the orthotropic yield function is more accurate in predicting springback than the isotropic function. In addition, the effect of the die radius on springback was studied. They found out that the increase of the bending radius caused a higher springback of the bend materials. Seo et al. [14] conducted a study to evaluate the effect of constitutive equations on the springback prediction accuracy. They used two yield functions, Hill48 and Yld2000, in combination with the Yoshida-Uemori hardening model in the finite element (FE) simulation to predict the springback of the U-bend part and drawn T-shape part. Both parts were made of TRIP steel. They found out that it is essential to choose the right yield function to get an accurate prediction of springback. Mulidran et al. [15] conducted a numerical simulation of the drawing hat-shaped part made of the DP600 and DC04 steels with the use of two forming methods: drawing with a blank holder and crash forming with a pressure pad. They studied the effect of forming methods and the various process parameters on springback amount. Their results indicated that the higher blank holder and pad pressure have positive effects on reducing springback. Additionally, crash forming with a pressure pad showed lower springback in comparison with drawing with a blank holder. The work by Jung et al. [16] aimed at studying anisotropic hardening behaviour and the springback of AHSS steels. They proposed the modified isotropic-kinematic hardening model, which they used in the simulation of U-bending. Their model showed better results in the predicting springback in comparison with the isotropic hardening model.

The novelty of this work lies in findings which indicate that the isotropic hardening model is sufficient for predicting the springback of formed parts made of aluminium alloy. This model is simpler, and does not need cyclic shear tests. In addition, isotropic hardening models do not use as

many parameters as kinematic hardening models. The accuracy of springback prediction with use of the isotropic hardening model was high. In addition, we found out that the phenomenon of the degradation of Young's modulus is not present for aluminium alloys which are precipitation hardened, and that the degradation is not as significant as in AHSS steels (max. 2% degradation of Young's modulus for aluminium alloys). Villuendas et al. [17] and Roca et al. [18] studied the effect of plastic deformation on the changes of Young's modulus of metallic alloys. They reported that, in aluminium alloys, there were no appreciable changes in the E value. This conclusion is consistent with our findings. These changes are related to the dislocation density changes. However, even though the dislocation density is high, the values of parameter l (length of dislocations) are very low, due to the interaction between nanometric precipitates and dislocations in the aluminium alloys. The Mott model then shows that the change of Young's modulus E is very small. Kinematic hardening models also take into account other phenomena, such as the transient behaviour, work-hardening stagnation, and the permanent softening. These phenomena were not observed in the material studied in this work.

These findings have a significant financial impact. For example, it is not necessary to conduct time-consuming tests on special (expensive) equipment, which are used to determine the parameters for kinematic hardening models.

In addition, the detailed analysis of a complex shaped part made of aluminium alloy with a significant thickness of 3 mm, mainly used in car production, was conducted. In most of the studied literature, the research was done on simply shaped parts.

In this research work, a FEM was used to predict the springback of a car body part made of aluminium alloy AA6451-T4. The finite element analysis (FEA) was conducted to investigate the influence of the used yield functions in the numerical simulation on springback prediction accuracy. Three sections were used for springback evaluation; in these sections, the thickness and part profile were measured and compared with the experimental results. The experimental results were given by the automobile manufacturer. Additionally, the computation times for the various yield functions were compared.

2. Materials and Methods

In the presented work, aluminium alloy AA6451-T4 with a thickness of 3.00 mm was used as the blank. Mechanical properties were measured by uniaxial and biaxial tensile tests. To obtain the required data for the FEM model, the specimens for the uniaxial tensile test were cut in three different orientations (0° , 45° , and 90° to the rolling direction). Specimens for the uniaxial tensile test were produced according to the EN 10002-1:2002-11 standards. Several specimens were tested for each orientation, and the average values of the basic mechanical properties (displayed in Table 1) were obtained by the formula

$$X_{av} = \frac{X_0 + 2X_{45} + X_{90}}{4} \quad (1)$$

where X is the mechanical parameter, and the subscripts denote the orientation of the specimen with respect to the rolling direction of the sheet. The elastic mechanical properties of the aluminium alloy are shown in Table 2.

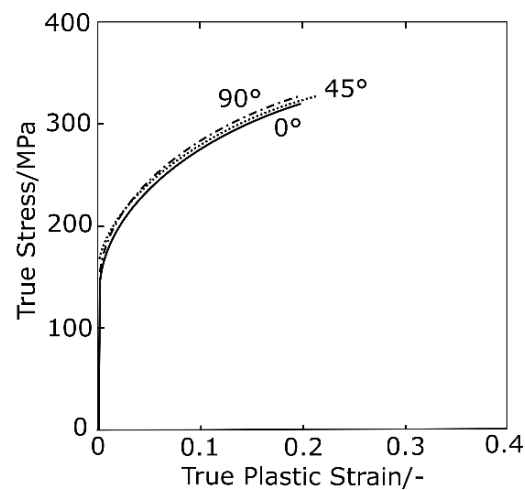
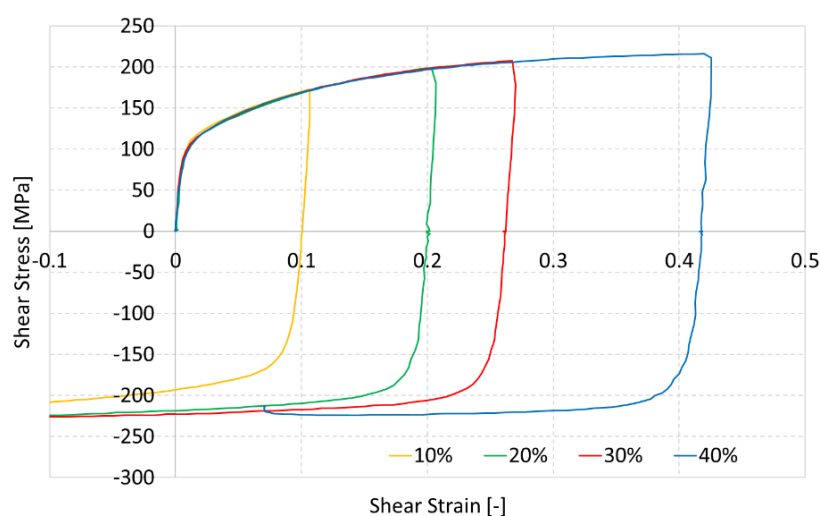
Table 1. The uniaxial tensile test data of the AA6451-T4 sheet [19].

Orientation	Yield Strength σ_y (MPa)	Normal Anisotropy r (–)
0°	151.28	0.62
45°	171.20	0.33
90°	163.60	0.80
Average value	164.32	0.52

Table 2. The elastic mechanical properties of AA66451-T4 [19].

Sample	Density, ρ (g·cm ⁻³)	Young's Modulus, E (GPa)	Poisson's Ratio, ν
AA6451-T4	2.7	70.0	0.3

The true stress–strain curves obtained in three different orientations (0°, 45°, and 90° with respect to the rolling direction) are shown in Figure 1. The tension-compression test (Figure 2) was started by the tension load as the first part of the full cycle. After a specific crosshead stroke corresponding to a defined pre-strain level, the load was reversed to compression until it reached the crosshead displacement according to a given compression strain. Next, reloading in tension direction was introduced until the crosshead stroke was equal to that in the first tension. The investigated phenomena, such as the change of Young's modulus E , transient behaviour, work-hardening stagnation, and permanent softening, were not observed in the material studied in this work, as we can see from Figure 2.

**Figure 1.** The experimental stress-strain curves from the tensile test.**Figure 2.** The cyclic tension-compression experimental curve of AA6451-T4 [19].

The equal biaxial tensile yield stress and the biaxial anisotropy r_b are given in Table 3. The parameters obtained from this test are necessary to determine the advanced yield functions.

Table 3. The equal biaxial tension test data [19].

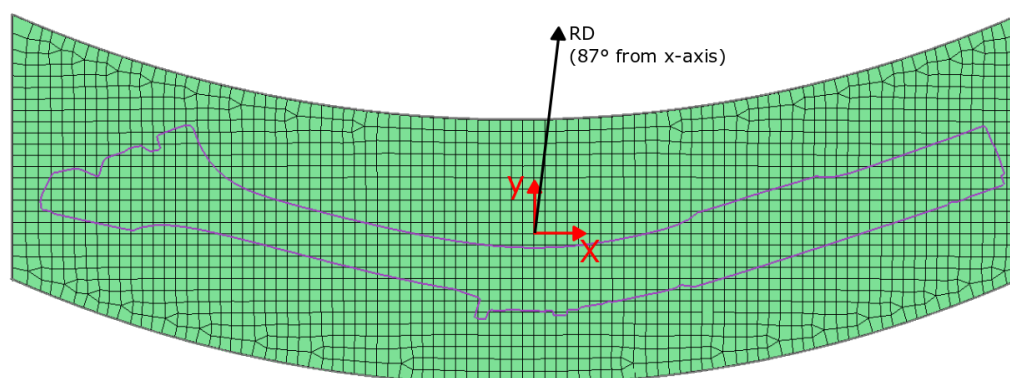
Biaxial Yield Stress σ_b (MPa)	Biaxial Anisotropy r_b (-)
153.60	0.55

3. Numerical Model

The springback computation was performed using the dynamic explicit code in the PAM-Stamp 2G software. The tool setup imported in the simulation software is shown in Figure 3. The tool consists of punch, blankholder, and die. The tool is aligned with the global z- axis without a plane of symmetry. The blank was positioned between the die and blankholder.

**Figure 3.** The tool geometry in the PAM-Stamp 2G software.

The blank was meshed by the quadrilateral shell elements which were 23 mm in size. The refinement level of the elements was set to 4 so that the smallest elements after refinement had a size of 2.875 mm. The number of integration points was set to 11, which is recommended for springback computation. The friction coefficient was set to 0.08, which responded to the grease lubrication. The initial meshed blank with a rolling direction is illustrated in Figure 4.

**Figure 4.** The rolling direction on the initial meshed blank.

The obtained values of the mechanical properties were used as the basic input for the material model in the FEM simulation. The accuracy of the springback prediction for several yield functions was investigated by the FEM simulation. The yield function describes the material transition from the elastic state to the plastic state. It can be described as a function of the area that limits the elastic area of the multi-axis stress plane. In 1948, Hill introduced the concept of material anisotropy in yield functions. According to Hill's plasticity conditions [20], in case of uniaxial load, a local thickness reduction occurs in a direction sensitive to the sample load. Hill assumed that the direction of the compression is in line with the direction of zero extension and, therefore, the deformation of the

narrowed areas is only reflected as a reduction in thickness. This is assumed for the plane strain (σ_1 —major stress, σ_2 —minor stress, and $\sigma_3 = 0$). If we assumed that the anisotropy axes are identical with the main guideline strain tensor ($\sigma_x = \sigma_1$, $\sigma_y = \sigma_2$, $\tau_{xy} = 0$), it is possible to describe the Hill48 yield function by the formula

$$2f(\sigma) = (G + H)\sigma_{xx}^2 + (F + H)\sigma_{yy}^2 - 2H\sigma_{xx}\sigma_{yy} + 2N\sigma_{xy}^2 = 1 \quad (2)$$

where σ_{xx} , σ_{yy} , and σ_{zz} are stresses in the RD (x), TD (y), and thickness (z) directions, respectively; σ_{xy} , σ_{yz} , and σ_{zx} are the shear stresses in xy, yz, and zx directions. Parameters F , G , H , and N are material parameters that describe the anisotropy of the material. If $F = G = H = 1$, and $N = 3$, the Hill48 function is reduced to the von Mises criterion, or as it is called in FEM code, the Hill48 isotropic criterion. A more common description is based on normal anisotropy in the 0° , 45° , and 90° directions to the rolling direction. Then, the material parameters F , G , H , and N can be described by

$$F = \frac{r_0}{r_{90}(r_0 + 1)}, \quad G = \frac{1}{r_0 + 1}, \quad H = \frac{r_0}{r_0 + 1}, \quad N = \frac{(r_0 + r_{90})(1 + 2r_{45})}{2r_{90}(1 + r_0)}. \quad (3)$$

For orthotropic hardening law and the values of anisotropy under 1.0, the Hill90 yield function is more suitable. This function is considered to be more suitable for aluminium alloys, and it is based on a non-quadratic transition function. In order to construct this function, the values from the uniaxial tensile test are deficient. For a complete description of this function, the biaxial test data are also required. The function can be described as

$$\left(\frac{\sigma_1}{\sigma_2}\right)^2 + \left(\frac{\sigma_2}{\sigma_{90}}\right)^2 + \left[(p + q + c) - \frac{p\sigma_1 + p\sigma_2}{\sigma_b}\right] \left(\frac{\sigma_1\sigma_2}{\sigma_0\sigma_{90}}\right) = 1, \quad (4)$$

where σ_0 is uniaxial tensile stress in the rolling direction, σ_{90} is uniaxial tensile stress in the direction normal to the rolling direction, σ_b is the stress under the balanced biaxial stress, and the c , p , and q parameters are defined as follows [21]:

$$c = \frac{\sigma_0}{\sigma_{90}} + \frac{\sigma_{90}}{\sigma_0} - \frac{\sigma_0\sigma_{90}}{\sigma_b^2}, \quad (5a)$$

$$\left(\frac{1}{\sigma_0} + \frac{1}{\sigma_{90}} - \frac{1}{\sigma_b}\right)p = \frac{2R_0(\sigma_b - \sigma_{90})}{(1 + R_0)\sigma_0^2} - \frac{2R_{90}\sigma_b}{(1 + R_{90})\sigma_{90}^2} + \frac{c}{\sigma_0}, \quad (5b)$$

$$\left(\frac{1}{\sigma_0} + \frac{1}{\sigma_{90}} - \frac{1}{\sigma_b}\right)q = \frac{2R_{90}(\sigma_b - \sigma_0)}{(1 + R_{90})\sigma_{90}^2} - \frac{2R_0\sigma_b}{(1 + R_0)\sigma_0^2} + \frac{c}{\sigma_{90}}, \quad (5c)$$

where R_0 is the anisotropy value for the uniaxial tension in the rolling direction and R_{90} is the anisotropy value for the uniaxial tension in the in-plane direction, perpendicular to the rolling direction.

The Barlat's material models describe the plastic behaviour of a material in a more detailed way than Hill's functions, but the higher number of parameters increases the calculation time. The Barlat89 model needs three parameters for its complete formulation, by which it is possible to describe the plane strain behaviours. Those parameters are defined in Table 4. The formulation is the following:

$$f = a|k_1 + k_2|^M + a|k_1 - k_2|^M + (2 - a)|2k_2|^M = 2\sigma_e^M, \quad (6)$$

where M is the exponent related to the crystallographic structure of the material, and k_1 and k_2 can be described as

$$k_1 = \frac{\sigma_x + h\sigma_y}{2}, \quad k_2 = \left[\left(\frac{\sigma_x - h\sigma_y}{2}\right) + p^2\tau_{xy}^2\right]^{1/2}, \quad (7)$$

where a , h , and p are the material model parameters.

Table 4. The material constants for the Barlat89 yield function ($m = 8.0$).

a	c	h	p
1.3033	0.9556	0.9247	0.8465

A more precise function was presented by Barlat in 2003, called Barlat2000, where the linear transformation method was used. In the FEM software, this function is described by the eight parameters shown in Table 5. The formulation for this model is as follows:

$$\phi = \phi'(X') + \phi''(X'') = 2\bar{\sigma}^a, \quad (8)$$

where a is an exponent related to the crystallographic structure of the material and ϕ' and ϕ'' are two isotropic functions described as follows [22]:

$$\phi' = |X_1' - X_2'|^\alpha; \phi'' = |2X_2'' + X_1''|^\alpha + |2X_1'' + X_2''|^\alpha. \quad (9)$$

Table 5. The material constants for the Barlat2000 yield function ($a = 8.0$).

a ₁	a ₂	a ₃	a ₄	a ₅	a ₆	a ₇	a ₈
1.065173	0.841891	0.960059	0.958652	1.034037	1.027112	0.838988	0.877033

According to several works [23–25], the Vegter yield function should be more suitable for special steels and aluminium alloys due to its more convenient results. The Vegter criterion describes the yield locus more accurately from a series of physically tested points. According to Vegter, it is possible to establish the first quadrant of the yield function on the basis of the basic experimental measurement. To construct the ellipses, the Bezier interpolations between each point need to be performed. Every point requires three parameters to be defined, the main stresses σ_1 and σ_2 , and the strain vector $\rho = d\varepsilon_2/d\varepsilon_1$. For a complete description of planar anisotropy, it is necessary to obtain 17 parameters from 9 mechanical tests. The mathematical expression of this function is

$$\begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix} = (1 - \lambda)^2 \begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix}_i^r + 2\lambda(1 - \lambda) \begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix}_i^h + \lambda^2 \begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix}_{i+1}^r, \quad (10)$$

where λ is the parameter for the Bezier interpolation subscript, i refers to the first reference point, r and h refer to a reference point and hinge point, respectively [26].

It is possible to use a simplified formula—Vegter-Lite. For this optional model, only 7 parameters from three mechanical tests (uniaxial tensile test, hydraulic bulge test, and the measurement of anisotropy) need to be defined. In this model, the second order Bezier interpolation is replaced by the second order NURBS interpolation, and the weight factor w —that controls the position of the curve between the points—is introduced. The formula for this model is

$$\begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix} = \frac{(1 - \lambda)^2 \begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix}_i^r + 2\lambda(1 - \lambda) \begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix}_i^h + \lambda^2 \begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix}_{i+1}^r}{(1 - \lambda)^2 + 2w\lambda(1 - \lambda) + \lambda^2}. \quad (11)$$

To fully define the material behaviour, the hardening curve of the material is also required. The Voce hardening curve gives the best correlation with the experimental results at an orientation of

0° from the rolling direction. This law provides a sufficient description of the elastic behaviour for aluminium alloys. The Voce hardening law is given by the equation

$$\sigma_y(\varepsilon_p) = A - Be^{-C\varepsilon_p}, \quad (12)$$

where A , B , and C are parameters defined in Table 6.

Table 6. The parameters for the Voce hardening curve.

A (MPa)	B (MPa)	C (–)
359.093260	169.310139	9.374256

To determine the failure criteria, Keller's and Goodwin's forming limit curve (FLC) model was used [27]. This empirical formula was obtained from experimental trials, and requires only two parameters: the thickness of the material and the strain hardening coefficient. The formula can be written as follows:

$$\varepsilon_{10} = \frac{(23.3 + 14.13t_0)n}{0.21}, \quad (13)$$

where t_0 is the initial thickness of the sheet and n is the strain hardening coefficient.

The simulation process consisted of three operations (stamping, trimming, and springback). Stamping was carried out as one continuous process in which the die moved at a speed of 100 mm/s. The blank was positioned between the die and the blankholder during holding. The die movement was set in the $-z$ -direction at 300 mm until the blank was clamped. Subsequently, a blankholding force of 1900 kN was applied. The die and blankholder moved in the $-z$ -direction until the tool was closed. After the part was fully formed, the trimming operation was performed. The trimming curve is shown in Figure 5.



Figure 5. The trimming curve on the punch.

4. Results

The results obtained from the numerical simulation were compared with the experimental ones. The springback was measured using locked nodes of the model. The stamped part profile and thickness were evaluated in three sections displayed in Figure 6. Section 1 is located on the right side of the part. This section passes through the hole, and due to the complex shape of the stamped part, the effect of the springback is quite significant in this area. The hole also runs through Section 2; this section is located on the left side of the part. The third section is located between section one and two. Figure 6 also shows the centres of the coordination systems used for profile evaluation. These centres were key for assembly purposes.

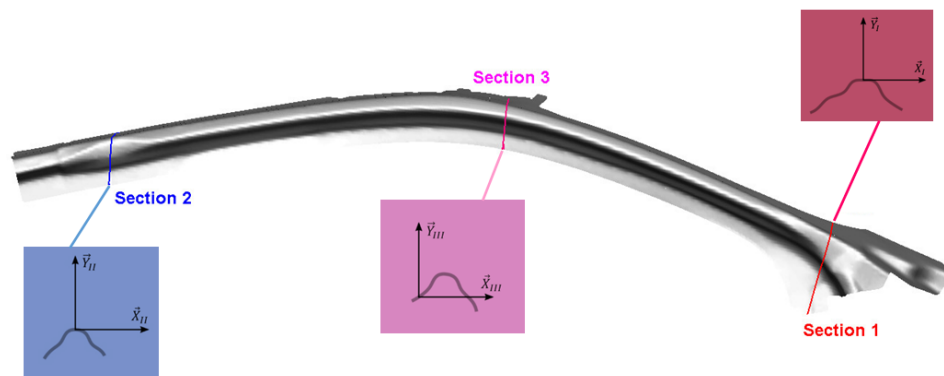


Figure 6. The sections used for profile and thickness evaluation.

4.1. Profile Analysis

The profile of the part was measured before and after springback. Since the stamped part copies the shape of the die, the profiles after stamping were almost identical for every material model. Good correlation of the experimental and numerical results can be observed after stamping. The comparison of the experimental results and the FEM simulation after springback for each section is shown in Figure 7. Subsequently, the FEM results for each yield function after springback were also compared.

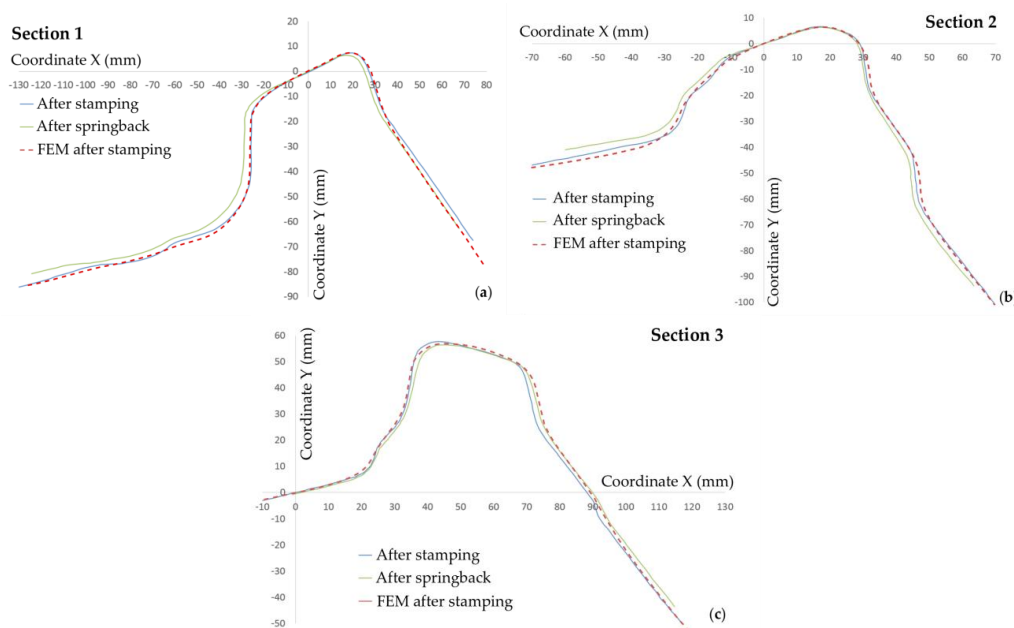


Figure 7. The comparison of the part profile from FEM simulation after springback with the experimental results after stamping and springback in (a) Section 1; (b) Section 2; (c) Section 3.

The deviations from the experimental results were measured on the left and on the right side of the sections. For assembly purposes, the deviation was measured in mm. Due to the difficulty of using conventional methods for springback measurement, the MATLAB system was implemented into the evaluation process. A simple program in the MATLAB environment was developed by the authors, which made it possible to measure the springback more accurately and easily. At first, a section was imported into the program, and the coordinate system was defined. Then, by selecting one point on the arm, a straight line parallel to the X-axis was created. Further points could only be created on this line, and the distance between each point was evaluated.

Figure 8 shows the experimental results and results from the numerical simulation in Section 1. Deviation from the experimental results for every yield function is shown in Table 7. A positive value means that the numerical simulation predicted a higher springback value than the experimental ones, and a negative value means that the simulation value was lower than the experimental ones. From these values, the average value of springback was calculated by the following equation:

$$X_{av} = \frac{|X_L| + |X_R|}{2}, \quad (14)$$

where X_L is the offset for the left side of the part and X_R is the offset for the right side of the part.

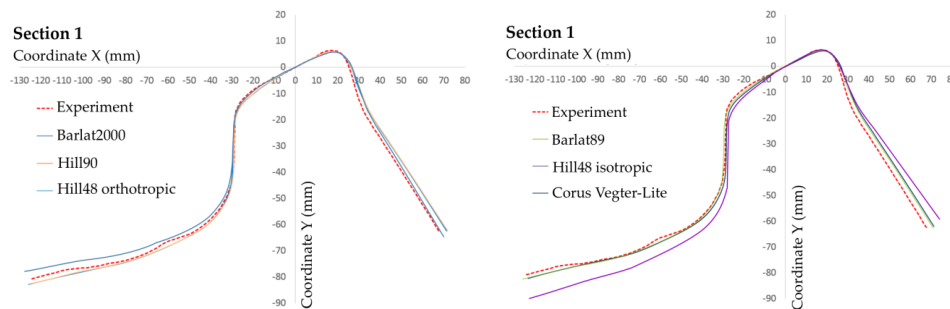


Figure 8. The comparison of the yield functions in Section 1.

Table 7. The results of springback in Section 1.

Yield Function	Left Side Offset (mm)	Right Side Offset (mm)	Yield Function	Left Side Offset (mm)	Right Side Offset (mm)
Barlat2000	+3.36	+0.88	Barlat89	−1.28	+1.15
Hill90	−2.14	+3.77	Hill48 isotropic	−9.29	+3.53
Hill48 orthotropic	−1.85	+4.51	Vegter-Lite	−1.67	+1.59

The lowest average deviation of 1.21 mm was measured when the Barlat89 yield function was used. Yield functions Barlat2000 and Vegter-Lite had average deviations of 2.12 and 1.63, also showing very good correlation with the experimental results. The highest average deviation of 5.44 mm was measured for the isotropic Hill48 yield function. In this function, the anisotropy of the material was not considered.

In Section 2 (Figure 9), greater deviation of the numerical springback values from the experimental ones can be seen on the right side of the profile. The reason is the shape of the part. Springback did not appear so significantly on the left side of the part where the material is compressed. The results are displayed in Table 8. In this section, the material model with the Vegter-Lite yield function showed the lowest average value of deviation: 2.65 mm. The Barlat89 yield function with 2.68 mm average deviation also shows good correlation. The other material models show very similar results, but the isotropic Hill48 method shows the highest deviation (3.51 mm) from the experimental results.

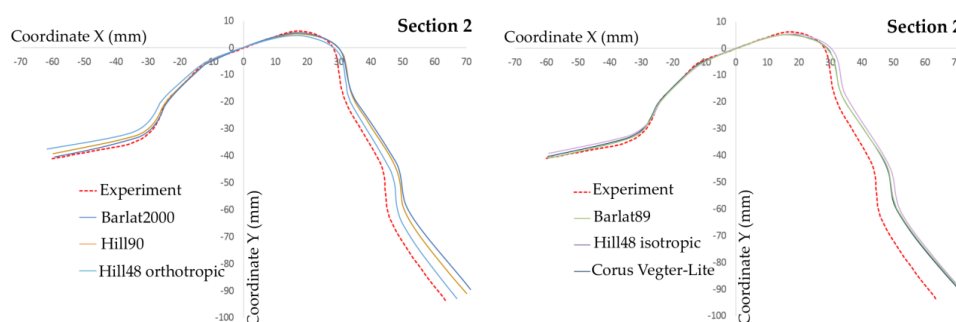
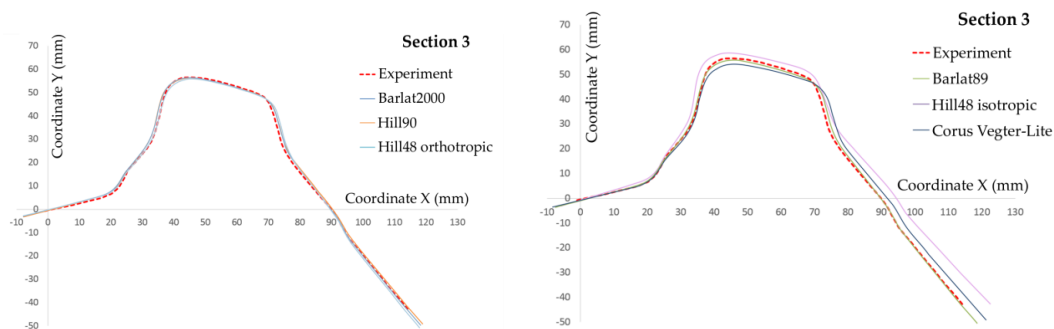


Figure 9. The comparison of yield functions in Section 2.

Table 8. The results of springback in Section 2.

Yield Function	Left Side Offset (mm)	Right Side Offset (mm)	Yield Function	Left Side Offset (mm)	Right Side Offset (mm)
Barlat2000	+0.45	+5.79	Barlat89	−0.20	+5.15
Hill90	+1.62	+4.48	Hill48 isotropic	+1.56	+5.46
Hill48 orthotropic	+3.78	+2.21	Vegter-Lite	+0.04	+4.90

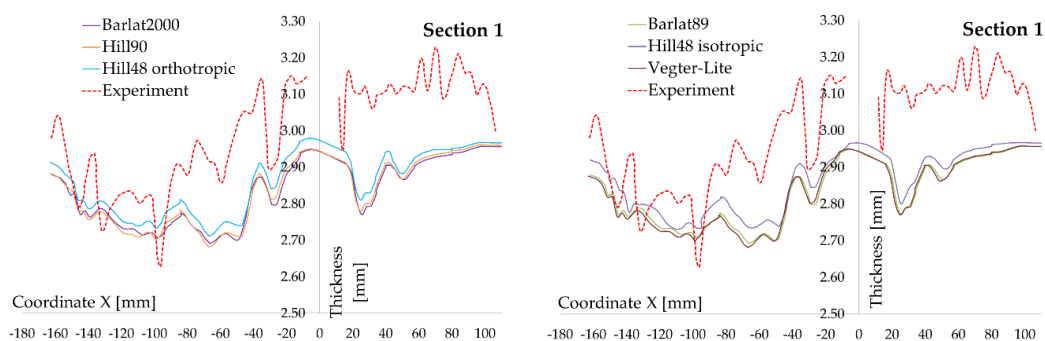
In the third section, the deviation was measured only on the right side because, on the left side, the deviation was too low, as shown in Figure 10. In this section, the best correlation was achieved with the Barlat yield functions, where the Barlat2000 deviation was 0.21 mm and the Barlat89 deviation was 0.32 mm. The highest deviation of 5.26 mm was measured for the isotropic Hill48 yield function. The results for all the material models are shown in Table 9.

**Figure 10.** The comparison of yield functions in Section 3.**Table 9.** The results of springback in Section 3.

Yield Function	Right Side Offset (mm)	Yield Function	Right Side Offset (mm)
Barlat2000	−0.21	Barlat89	−0.32
Hill90	−0.75	Hill48 isotropic	+5.26
Hill48 orthotropic	+0.37	Vegter-Lite	+1.97

4.2. Thickness Analysis

The next investigated parameter in this work was thickness, which was also measured in mentioned sections. The thickness of the material can significantly influence the accuracy of the FEM prediction. Since shell elements, which were used in this work, are suitable for thicknesses up to 1 mm, the volume element should give better results. However, in the FEM software, it is possible to define the volume elements for only the Hill48 and Barlat2000 yield functions. For comparison purposes of all previously mentioned yield functions, the models must have shell elements. Figures 11–13 show comparisons of the experimental and FEM results in each section.

**Figure 11.** The thickness distribution in Section 1.

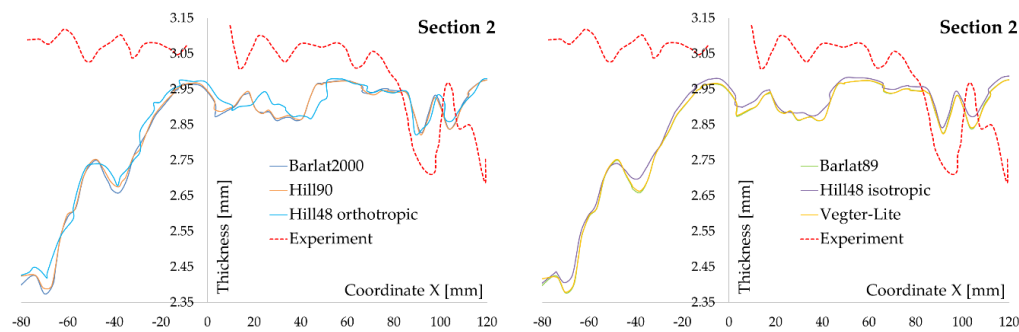


Figure 12. The thickness distribution in Section 2.

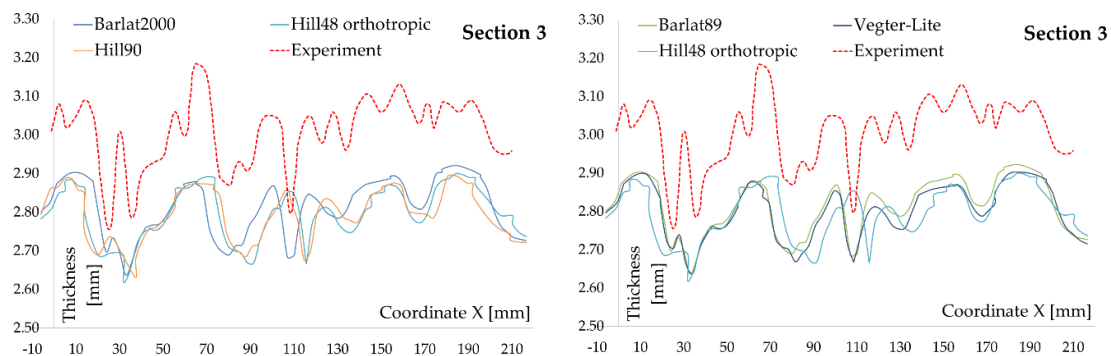


Figure 13. The thickness distribution in Section 3.

From the results of the thickness distribution in the individual sections, it is clear that the material models with the Barlat (Barlat89 and Barlat2000) yield functions, and the Vegter-Lite yield function shows very similar results. Although the description of the yield function and the amount of data needed for their definition is different between those models, the results of the thickness distribution were practically the same. The Hill48 model, either isotropic or orthotropic, exhibited significant deviations. The average values of thickness are shown in Table 10. The experimental results show a higher average thickness than the thickness data obtained by the FEM simulation.

Table 10. The comparison of the average thickness of each section.

	Barlat2000	Barlat89	Vegter-Lite	Hill90	Hill48 Orthotropic	Hill48 Isotropic	Experiment
Section 1	2.825	2.828	2.826	2.828	2.834	2.858	2.980
Section 2	2.856	2.847	2.853	2.852	2.847	2.859	2.991
Section 3	2.809	2.810	2.792	2.798	2.797	2.809	2.990

4.3. Computation Time

With the increased complexity of the yield function formulation and, thus, with the increased number of necessary variables, the calculation time was increased. For the Hill48 isotropic model, where the anisotropy of material was not considered, the computation took around 13 h and 47 min. For the Hill48 orthotropic and Hill90 models, where the anisotropy of the material was considered, the computation time took 14 h. In the Barlat yield functions, where the material's crystallographic structure was considered, the computation time increased significantly to around 15 h and 20 min. The Vegter-Lite yield function achieved a similar accuracy and computation time as the Barlat material models. Figure 14 shows the comparison of the computation times for each yield function.

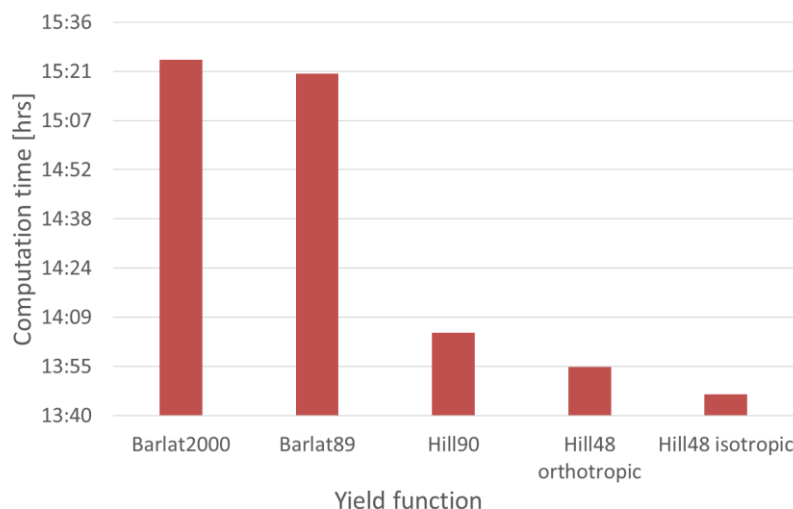


Figure 14. The computation time comparison.

5. Conclusions

The accuracy of the springback prediction is one of the most challenging problems in the numerical simulation of forming processes. In the present article, the influence of the yield function on the accuracy of springback prediction with the use of numerical simulation was investigated. Three sections were defined on the formed part. In these sections, the thicknesses of the profile of the stamped part after springback calculation were evaluated. After the stamping operation, for all the examined yield functions, the sheet metal copied the shape of the die. Visible differences can be seen after cutting and springback calculation. The results of the Barlat2000, Barlat89, and Vegter-Lite yield functions were in good correlation with the experimental results. Hill's yield functions (Hill90, isotropic Hill48, and anisotropic Hill48) were not as accurate as the yield functions mentioned above. Barlat's yield functions takes into account the material's crystallographic structures. The Barlat89 yield function, with an average deviation of 1.40 mm from experimental results of springback, is not suitable for materials with strong anisotropy [26]. Additionally, this model cannot capture the change of yield stress and the Lankford coefficient values. However, the advantage of this function lies in its simple mathematical description, and in the ability of the accurate plastic behaviour prediction (yield locus) of aluminium alloys, thus, the results of the Barlat89 model are more accurate than the results obtained with the use of Hill's yield functions. The Barlat2000 yield function is an improved version of the Barlat89 model, but the description of this improved model in the numerical simulation is more difficult. This is the reason why this function is not used as much in industrial practice. The experimental thickness values were higher than the predicted ones in all cases. From the results of the thickness distribution in the individual sections, it is clear that the yield functions of the Barlat's family and Vegter-Lite yield function show very similar results. Although the description of the yield function and the amount of data needed for their definition is different between those models, the results of the thickness distribution were practically the same. The Hill48 model, either isotropic or orthotropic, exhibited significant deviations. This can be attributed to the shell elements used in the numerical simulation or to the Voce isotropic hardening law. However, the Voce hardening law exhibited good correlation with the experimentally measured FLC (forming limit curve). One of the results of the work is that the combination of the isotropic hardening law with the isotropic yield function did not achieve accurate springback prediction results. The combination of more advanced yield functions (Barlat2000, Barlat89 and Vegter-Lite) with the isotropic Voce hardening law improved the accuracy of the springback prediction, but the computation time was increased by approximately an hour. Additionally, it was found out that the investigated phenomena which are required for the determination of the kinematic hardening model, such as the change of Young's modulus E , transient

behaviour, work-hardening stagnation, and permanent softening, were not observed in the aluminium alloy studied in this work. Our research confirmed that in aluminium alloys, there are no appreciable changes in the Young's modulus value.

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