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Rocking and Kinematic Approaches for Rigid Block Analysis of Masonry Walls: State of the Art and Recent Developments

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Abstract: The assessment of the rocking and overturning response of rigid blocks to earthquakes is a complex task, due to its high sensitivity to the input motion, variations in geometry and dissipation issues. This paper presents a literature review dealing with classical and advanced approaches on rocking motion with particular reference to masonry walls characterized by a monolithic behavior. Firstly, the pioneering work of Housner based on the concept of the inverted pendulum is discussed in terms of the most significant parameters, i.e., the size and slenderness of the blocks, the coefficient of restitution and ground motion properties. Free and restrained rocking blocks are considered. Then, static force-based approaches and performance-based techniques, mostly based on limit analysis theory, are presented to highlight the importance of investigating the evolution of the rocking mechanisms by means of pushover curves characterized by negative stiffness. From a dynamic perspective, a review of probabilistic approaches is also presented, evaluating the cumulative probability of exceedance of any response level by considering different earthquake time histories. Some recent simplified approaches based on the critical rocking response and the worst-case scenario are illustrated, as well.

Keywords: rocking; dynamics of rigid blocks; restrained blocks; out-of-plane behavior; force-based approaches; displacement-based approaches

1. Introduction

Masonry buildings are constituted of three-dimensional assemblies of walls, where the out-of-plane behavior of each wall is highly influenced by the type and strength of connection with the others. There are traditional (e.g., steel ties or buttresses [1]) and innovative (e.g., composites [2,3]) techniques to ensure a safe seismic behavior at the local level, by considering the typical biaxial stress state that could involve energy dissipation [4–6]. Nevertheless, when a global box-type behavior is not guaranteed, the walls, especially the peripheral ones, are more vulnerable to out-of-plane overturning, which is one of the main causes of damage or collapse induced by earthquakes on existing masonry structures. The main deficiencies are the lack of proper connections between orthogonal walls, the absence of connecting ties, insufficiently rigid floor diaphragms, low strength and deterioration of materials.

However, if a monolithic behavior can be assured for such walls, they can be regarded as rigid blocks, and their out-of-plane seismic response can be treated through two fundamental approaches

on which this paper is mainly focused: rocking dynamics and kinematic analysis. Discrete and finite elements can also be used to assess the response of masonry structures under earthquakes [7], but those involve many uncertainties mainly related to the definition of the constitutive laws of the materials and are more suitable to masonry walls far from monolithic behavior.

On the other hand, the apparent simplicity of the rocking of a single-degree-of-freedom (SDOF) rigid block hides many subtle phenomena, largely investigated in the literature, with more emphasis since the 1980s. The pioneering work of Housner [8] was the basis for the subsequent research in the field of rocking blocks, studied in structural dynamics and particularly in earthquake engineering. After him, the first modern contributions to the rocking issues are due to Aslam and Scalise [9], who considered the motion of a free rocking block subjected to ground motions. They analyzed the transition to the sliding phenomenon, as well. The work of Ishiyama [10] proposed analytical formulations for the transitions between different types of motion (sliding, rotation, translational or rotational jumping).

Despite the numerous studies carried out so far on rigid body rocking (an extensive review work on this subject is reported in [11]), little literature describing the state of the art on masonry walls treated as rigid bodies is available, let alone the rocking of restrained walls. Therefore, this paper proposes an extensive review on kinematic and rocking approaches for monolithic masonry walls, specifically regarded as rigid blocks on rigid foundations, also including recent developments and different models accounting for lateral constraints. These modeling approaches can be used for defining the non-linear static and dynamic response of masonry buildings, especially the historic ones, for which macro-elements and their connections play a crucial role in the seismic vulnerability assessment [12–16].

The kinematic approach includes static force-based and displacement-based approaches based on standard and non-standard limit analysis methods, taking into account the evolution of motion over time through incremental kinematic analysis, while dynamic effects are considered more appropriately by means of the dynamic approach, since it also accounts for the energy dissipation in the motion.

The aim of this work is to illustrate and discuss the issues involved in the classical and non-classical theories and to present the two above-mentioned approaches adapted to masonry structures. In addition, emerging design techniques are illustrated to clarify the direction in which the research is practically oriented. Section 2 describes the basics of the rocking of a free-standing single-degree-of-freedom block, whereas Section 3 introduces the issue of restrained blocks, with an extension to multi-degree-of-freedom (MDOF) systems in Section 4. Section 5 illustrates the static approaches, referring to the static force-based and displacement-based approaches. The last two sections comment on the influence of the input motion parameters on the dynamic response (Section 6) and the deterministic and probabilistic methods of the rocking analysis (Section 7).

2. Rocking of Free-Standing SDOF Block

2.1. The Classical Theoretical Model of Housner and the Main Geometric Parameters Influencing Rocking

Housner's formulation [8], referred to as classical theory in this paper, allows investigating the behavior of a rigid block on a rigid foundation under transient actions such as earthquakes. The considered single-degree-of-freedom (SDOF) block is a rigid prism with a rectangular cross-section, rocking about the two corners O and O' (Figure 1a) and supported by a flat rigid base. Neither bouncing nor sliding are considered in this formulation.

The main geometric parameters influencing the rocking response are two: (i) slenderness ratio α , namely the arctangent of thickness to height $2b/2h$; and (ii) radius vector R , which connects the pivot point to the block center of mass.

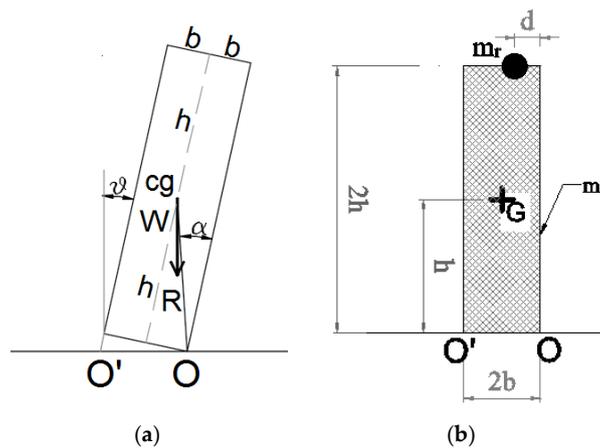


Figure 1. The rectangular free-standing rocking block of Housner's model (a) and the centroid G_c of the composed masonry block + the roof mass system with eccentricity d (b).

Considering the rotation ϑ (>0 if counter-clockwise) as a Lagrangian coordinate, the equation of motion takes the form:

$$I_0 \ddot{\vartheta} + \operatorname{sgn}(\vartheta) m g R \sin(\alpha - \operatorname{sgn}(\vartheta) \vartheta) - m \ddot{u}_g R \cos(\alpha - \operatorname{sgn}(\vartheta) \vartheta) = 0 \quad (1)$$

where I_0 is the inertia moment, m the mass and \ddot{u}_g the acceleration time-history (in g) of the mass. Equation (1) can be modified for homogeneous prismatic blocks as follows:

$$\ddot{\vartheta} + \operatorname{sgn}(\vartheta) p^2 \sin(\alpha - \operatorname{sgn}(\vartheta) \vartheta) - p^2 \ddot{u}_g \cos(\alpha - \operatorname{sgn}(\vartheta) \vartheta) = 0 \quad (2)$$

where p is the frequency parameter, equal to $\sqrt{\frac{3g}{4R}}$.

The equation of motion governing the rocking phenomenon of a free-standing rigid block is analogous to that used for investigating the response of a free-standing block with additional mass. A mass placed at the top of the block could simulate a roof mass, whose effect is one of changing the centroid of the composed system block + mass (Figure 1b). Thus, the inertia moment accordingly changes, with multiple values of radius and slenderness, but the equation of motion is basically the same as that of Housner's model (see Section 3) [17].

A relevant finding of the classical theory is that, between two rigid rectangular blocks of the same slenderness, the block with the higher radius vector is more stable. Moreover, between two blocks of the same width, the taller block is more stable than the shorter one [18]. However, the shorter block is related to a smaller amplitude of ground acceleration than that of the taller one. In addition, Kounadis [18] specifies that this behavior is not valid in the case of shocks, when the resisting forces (including inertia) do not have enough time to react, and the taller block results in being less stable than the shorter one (when overturning occurs without impact).

2.2. Geometry Influence and the Formulation of Energy Dissipation

The classical theory considers a rectangular block, but normally, asymmetric blocks or blocks with different shapes are to be modeled. Many authors studied the changes in the equations of motion caused by the modified geometry. These changes affect the following elements:

- moments of inertia and radius vector;
- energy loss at each impact and/or restitution coefficient;
- 3D motion to be considered instead of 2D motion;
- bouncing effect in the case of stocky blocks;
- additional terms in the equation of motion such as those due to damping and springs.

The energy dissipation or damping over rocking motion occurs when the block hits the base and the pivot point suddenly changes from O to O' (Figure 1). In the classical theory, the energy loss is expressed as a reduction of kinetic energy after each impact. By equating the moment of momentum about O' immediately before impact to that immediately after impact, Housner found the following expression, valid for slender and rectangular blocks:

$$e_H = 1 - \frac{3}{2} \sin^2 \alpha = \frac{2\mu^2 - 1}{2(\mu^2 + 1)} \quad (3)$$

where $\mu = h/b$ is the height to thickness ratio. The so-called "restitution coefficient" e_H depends on the slenderness ratio α : the higher the slenderness, the higher the number of impacts to get the same energy loss for the same input action. A value of $e_H = 0$ means totally inelastic impact, whilst $e_H = 1$ indicates a perfectly elastic impact. Generally, the energy loss during impact is lower in experiments than that predicted by Housner's model. In other words, the experimental value of the restitution coefficient is higher than the theoretical one e_H .

Experimental tests were performed to identify the values of restitution coefficients for different unreinforced masonry (URM) specimens of several construction materials and slenderness ratios. Liberatore et al. [19] tested a marble block rocking on a marble foundation, whereas Aslam et al. [9] tested a concrete block with an aluminum base rocking on a steel foundation. They noticed discrepancies of the restitution coefficients' values between experimental and analytical results. A relevant contribution was also given by Lipscombe and Pellegrino [20], who conducted four free rocking tests on steel blocks with $\mu = 1, 2, 4, 8$, colliding on a steel base. The blocks, initially tilted almost to the point of overturning and then released, exhibited values of restitution coefficients ranging between 0.88 and 0.93, although it was respectively $e_H = 0.25, 0.70, 0.91, 0.98$ (given by Equation (3)) for the mentioned μ values. Peña et al. [21] performed experimental tests on granite stones with the height to thickness ratio varying from 4–8 under free-vibrations, harmonic and random motions. The restitution coefficient was found to be slightly higher than the experimental one, as the real case does not fully comply with the hypotheses of no bouncing or an ideal block. Indeed, the body not being fully rigid, the moment of momentum was not conserved. In this line, also Sorrentino et al. [22] carried out an experimental campaign on URM solid clay brick or tuff specimens for two-sided and one-sided motion, with height to thickness ratios varying between 6.5 and 14.6. The authors found values of the restitution coefficient equal to 95% of the theoretical value. On the other hand, the lower energy loss during impact in most experiments is related to the unevenness of the colliding surfaces, generating consecutive impacts that reduce the energy loss. That aspect was recently analytically modeled with additional pseudo-bumps at any position of the section. These bumps change the evolution of motion as more impacts occur; the modified equations of motion are proposed in [23]. Recent experimental tests observed relevant results on the evaluation of the coefficient of restitution and motion decay [24–27]. It was shown how the interface material strongly influences the dissipation in the free rocking behavior. Blocks tested on rubber were seen to have the fastest energy dissipation followed by concrete and timber bases [24]. An improved estimation of the coefficient of restitution was proposed in [26] to accurately quantify the energy dissipation of free rocking members. In addition, even for symmetric structures with uniaxial shaking, multiple modes and three-dimensional responses are likely to occur [25].

2.3. Bouncing: Sliding Phenomena and the Role of Vertical Ground Motion

In the classical theory, bouncing and sliding are neglected, but these phenomena have to be sometimes necessarily considered. Indeed, for stocky blocks, say $\mu = 1, 3$, bouncing was shown to occur, and the response was seen to be very sensitive to the restitution coefficient value [20]. For these blocks, such as cubes, a two-dimensional or three-dimensional bouncing model is therefore required. In this case, the determination of a restitution coefficient is then crucial to get a reliable response. The impact response of short blocks is therefore complex to predict for these stocky structures. In [20],

a probabilistic approach was suggested to solve this issue. As concerns sliding, the absence of it is acceptable when the static coefficient of friction is larger than the ratio b/h [22,28]. Generally, masonry walls have this ratio much lower than the static friction coefficient, usually equal to 0.6–0.7 [28].

As regards the vertical ground acceleration, this was shown to have a marginal effect on the stability of a free-standing rocking column [29]. As the authors highlight, this is due to the fact that the ground acceleration enters the equation of motion after being multiplied with $\sin(\alpha - \vartheta) \ll 1$, whereas the horizontal acceleration enters the equation of motion after being multiplied by $\cos(\alpha - \vartheta) \approx 1$.

3. Restrained Rocking Blocks

In real conditions, rocking blocks such as masonry walls are restrained by flexible diaphragms. Several authors studied the influence of these boundary conditions on the dynamic response through analytical models [30], showing that diaphragm and shear-wall accelerations might increase with the flexibility of the diaphragm.

Housner's equation of motion was modified for considering the effect of single or smeared horizontal restraint [31] following a variational formulation. A single spring (or bed spring) was considered to have axial stiffness K (or K' per unit of length) simulating an element with a stabilizing effect, such as a strengthening device (tie-rod), transverse walls or flexible diaphragms. The stiffness K or K' can assume different values depending on the type of the roof and on the type of the roof-wall connection.

Due to the eccentricity of the roof mass, the radius vector changes depending on the rotation sign, $R_{r.cw}$ for clockwise and $R_{r.cc}$ for counterclockwise rotations, respectively (Figure 2, [17]). However, for slender blocks, the difference between $R_{r.cw}$ and $R_{r.cc}$ is negligible, and the radius vector can be assumed as the double of the radius vector R .

If R_r is the current roof radius vector and the block, subjected to a horizontal thrust H_r , is connected to a single horizontal restraint with stiffness K , the equation of motion reads [17]:

$$I_{0c}\ddot{\vartheta} + \operatorname{sgn}(\vartheta)m g R \sin \hat{\alpha} + \operatorname{sgn}(\vartheta)K R_r^2 \cos \hat{\alpha}_r [\sin \alpha_r - \sin \hat{\alpha}_r] + \operatorname{sgn}(\vartheta)m_r g R_r \sin \hat{\alpha}_r - H_r R_r \cos \hat{\alpha}_r - (m R \cos \hat{\alpha} + m_r R_r \cos \hat{\alpha}_r) \ddot{u}_g g = 0 \quad (4)$$

where the polar moment I_{0c} includes the contribution of the roof mass, $\hat{\alpha} = \alpha - \operatorname{sgn}(\vartheta)\vartheta$ and $\hat{\alpha}_r = \alpha_r - \operatorname{sgn}(\vartheta)\vartheta$, m is the block mass and \ddot{u}_g is the acceleration time-history (in gravity acceleration g units) of the mass.

The direction and magnitude of the vector H_r may be assumed constant during motion in the hypothesis of small displacements or acting up to a fixed displacement given, for instance, by the Mohr–Coulomb law. A higher centroid position is due to the presence of the roof mass. Indeed, the roof mass is assumed to participate in the rocking motion by rotating together with the masonry block. In the case of no eccentricity ($d = b$, Figure 1b), the distance of the centroid of the composed system G_c from the base is:

$$h_{G_c} = \frac{(m + 2m_r)h}{m + m_r} \quad (5)$$

and the inertia moment of the system is obtained with the following expression:

$$I_{0c} = I_0 + m_r R_r^2. \quad (6)$$

where I_0 is the polar inertia moment with respect to the base corner O , $I_0 = \frac{4}{3}mR^2$.

The effect of vertical restrainers with elastic pre-stressed tendons was also investigated [25]. These systems were shown to be effective in improving the response of rocking frames with small columns subjected to long-period excitations. As the size of the columns, the frequency of the excitations or the weight of the cap-beam increases, the vertical tendons do not apply any beneficial effect. Indeed, the resistance of tall rocking frames primarily originates from the mobilization of the rotational inertia of the columns.

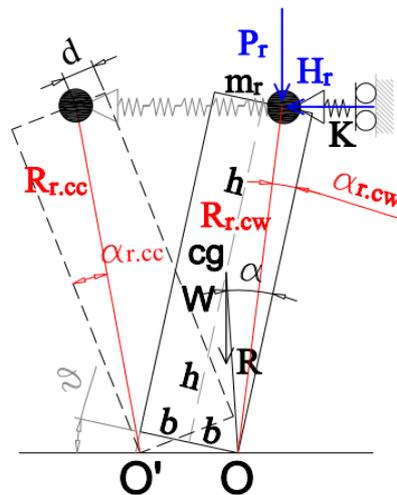


Figure 2. The rocking block with horizontal elastic restraint and roof mass.

4. Extension of SDOF to MDOF Rocking Blocks

Significant equivalence methods allowed extending the results obtained for an SDOF free-standing block to more complex structures, such as rocking frames and masonry arches. Sometimes, flexural bending movements can occur in the out-of-plane failure modes of masonry walls, and the degree of freedom cannot always be assimilated to one. In such cases, the multi rigid body dynamics could be more suitable to represent the structural response [32,33]. In [33], several tests on scale models of masonry 3D assemblies were performed on a shaking table, subjected to sinusoidal input of varying frequency and displacement amplitude, but constant acceleration, and a simplified multi-degree-of-freedom (MDOF) model was presented based on the relative rocking of two portions of walls, each one made of a finite number of bricks. A two-degree-of-freedom (2DOF) model was also presented by Simsir et al. [32] to compute the out-of-plane response of a wall that cracks at the bed joints. This model accounts for the flexibility of the diaphragm, the stiffness of the wall and the possibility for horizontal cracks developing under combined flexural moments and axial forces.

However, in contrast to the single rocking block applied to masonry that has been examined extensively, the dynamic behavior of masonry structures regarded as multi-block structures has not, to date, been exhaustively studied because of its complexity. Therefore, in order to prompt the scientific research in this line, it is here briefly illustrated how the rocking analysis can be applied to MDOF blocks.

When the degree of freedom is one, even with more complex structures, the response is determined by carrying out the same analysis as the free-standing block, but by assuming a modified value of the frequency parameter p . The square of the frequency parameter is the ratio of restoring the moment to the polar moment of inertia $p^2 = mgR/I_0$. For an SDOF rectangular block, the frequency parameter is equal to $\sqrt{\frac{3g}{4R}}$. The dynamic equivalence between a free-standing rocking frame and a free-standing rocking column was first shown by Makris and Vassiliou [34]. It refers to a rocking frame, made by a set of N columns of mass m_c and a cap beam of mass m_b . The frequency ratio of a rocking frame is given by $\hat{p} = p\sqrt{(1+2\gamma)/1+3\gamma}$, where $\gamma = m_b/(Nm_c)$ [34].

The results were extended to the case of an asymmetric rocking frame and the hinging masonry arch, defining proper values of p for each structural system [35]. When the Lagrangian parameter is not only one rotation, namely the system is MDOF, the analysis becomes much more challenging. The pure rocking of two or three-degree-of-freedom systems was investigated in [36]. The number N_u of configuration patterns that may lead to overturning instability for an MDOF with n blocks is:

$$N_u = 3^n - 1 \quad (7)$$

with corresponding total number of nonlinear differential equations $N \times n$.

Two-DOF systems were first studied in [37]. A relevant contribution to the definition of the dynamics of such systems was given by [38]: in this work, the equations of motion are obtained for each of the four + four modes (Figure 3), and criteria for the initiation of rocking and the transition between them are illustrated. In particular, the discussion of the impact between the lower block and the ground surface, together with the impact with the higher block and the lower one is performed. More recently, other works faced the more complex problem of 3DOF systems, generally solved only with analytical approaches, e.g., in [39], equations of motion are proposed to describe the response of multi-drum ancient Greek columns as 3DOF systems, to which 26 configuration patterns are associated.

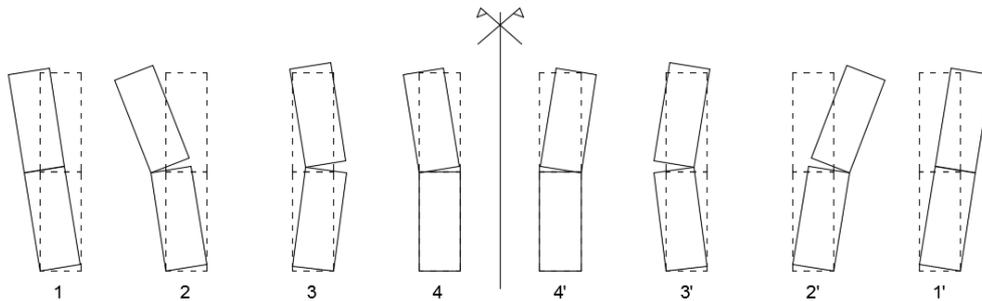


Figure 3. Eight ($N = 8 = 4 + 4$) possible configuration patterns of a two-rigid block ($n = 2$) system (adapted from [39]).

5. Static Approaches

5.1. Static Force-Based Approaches: Standard and Non-Standard Limit Analysis Methods

The static force-based (FB) approach is the most common and the simplest tool to estimate seismic design forces on rocking rigid blocks. It basically consists of the application of the plasticity theory, as first formulated by Kooharian [40] and Heyman [41]. According to this method, recognized as standard limit analysis, the application of the static theorem provides a lower-bound or safe solution of the collapse load factor, based on equilibrium equations, while the application of the kinematic theorem leads to an upper bound multiplier. The solution that satisfies the hypotheses of both theorems, equilibrium, compatibility and material conditions is the correct solution and provides the collapse load multiplier for the specific problem.

The simple mechanical model is based on the following assumptions: blocks with infinite compressive strength (rigid blocks), joints with zero tensile strength and sliding failures not allowed (friction at interfaces is sufficiently high to prevent sliding/twisting).

Based on such assumptions, the maximum horizontal force F_0 that a rigid block can undergo at the onset of rotation can be obtained by imposing equilibrium conditions. For instance, the overturning equilibrium of a cantilever SDOF wall about the pivot point O (Figure 1a) can be used to determine F_0 :

$$F_0 h = m g b \quad (8)$$

where m is the block mass. If F_0 is seen as the fraction of the block self-weight $F_0 = \alpha_0 \times m g$, it is possible to calculate the collapse multiplier or the maximum horizontal acceleration (in g) to which the block can be subjected in equilibrium conditions:

$$\alpha_0 = \frac{b}{h} \quad (9)$$

Similar equilibrium conditions can be imposed for blocks with different support conditions. For example, a load-bearing SDOF simply-supported wall (vertical-spanning wall) with a horizontal hinge at mid-height has a higher F_0 than that of a cantilever non-loadbearing wall. In particular, the

force of the load-bearing simply-supported wall is $4(1 + \Psi)$ times the force of the free-standing wall, where Ψ is the ratio of overburden weight and self-weight of the upper-half of the wall above mid-height [42].

In real conditions, rocking blocks such as masonry walls are restrained by horizontal diaphragms and transverse walls, as already mentioned above. In particular, in masonry buildings without a box-type behavior and subjected to seismic loadings, in-plane and out-of-plane failure mechanisms can take place where frictional resistances might play a predominant role. Hence, it is often necessary to consider the more realistic assumptions of the presence of sliding during rocking mechanisms and also to calibrate the role of friction with reference to other systems of resistance in masonry buildings (e.g., insertion of tie-rods, rigid diaphragms, etc.), as addressed in [13].

On the other hand, it is well known that when treating such so-called non-standard materials, the bounding theorems of plastic limit analysis do not generally provide unique solutions for collapse loads, due to the non-associated flow rules imposed by friction.

Drucker [43] was perhaps the first to point out that whilst the exact solution to a problem involving Coulomb friction interfaces could be bounded from above and below, unfortunately, such bounds will often be too wide to be of practical use. In this class of problems, in fact, the lower bound is generally the condition that assumes no friction and cannot obviously be proposed for the analysis of most masonry structures. A novel modeling strategy, based on a simplified macro-block approach, is capable of providing a closer range of solutions for the ultimate load factor [44] and a reliable solution falling within the range [45]. However, for the simple out-of-plane mechanism of a front wall interlocked with sidewalls involving vertical cogged cracks along the corners (laterally-constrained rigid block), the reliable solution can be computed directly by considering the full action of friction on the bed joints along the cracks. In fact, the frictional resistances can be expressed as [13]:

$$F_P = 2 \sum_1^n S_i = \gamma t h_b l_b \frac{n(n+1)}{2} f \quad (10)$$

where n is the number of rows crossed by the vertical crack line ($H = n \times h_b$), t is the thickness of the sidewalls, f is the friction coefficient, h_b and l_b are the height and length of the single masonry unit, respectively, and S_i is the single limiting shear force due to friction at contact interface i , obeying the cohesionless Coulomb's law and stepwise increasing from the top to the bottom of the wall (Figure 4).

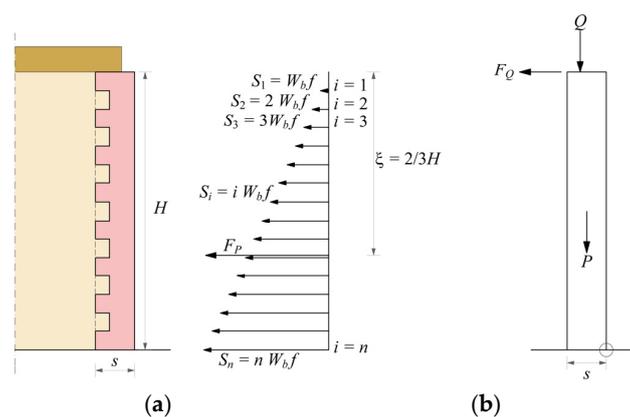


Figure 4. Frictional resistances transmitted to the façade wall by the shear walls (a) and the simply-supported horizontal diaphragm (b).

On the other hand, the presence of a horizontal diaphragm that is simply supported on the façade wall implies that beams can slip out, acting on the wall with a horizontal friction force, which is (cohesionless Coulomb's law):

$$F_Q = Q f_d \quad (11)$$

where f_d is the beam-wall friction coefficient and Q is the vertical load transmitted to the wall. Thus, by imposing equilibrium, the activation load multiplier α_0 for incipient rotation around the hinge at the base of the façade wall in Figure 4 is given by the expression:

$$\alpha_0 = \frac{P \frac{s}{2} + Q \frac{2}{3} s + F_Q H + F_P \frac{H}{3}}{QH + \frac{PH}{2}} \quad (12)$$

where s is the thickness of the front wall and the vertical resultant of the overload is assumed, e.g., applied at $s/3$ from the internal edge.

In general, the static force-based approach is very useful for the identification of the failure mechanisms, especially when not self-evident, since it is based on minimization routines. However, the evaluation of the static multiplier is only related to the onset of rocking, not to the overturning that may occur under dynamic actions. In fact, the evolution of the dynamic system over time is neglected, and the reserve capacity of the rocking phenomenon is not considered.

5.2. Displacement-Based Approaches

As dynamically-loaded walls can sustain accelerations well in excess of their ‘quasi-static’ capabilities [46], the displacement-based (DB) approach is revealed to be a more realistic tool to assess the seismic design forces on a rocking wall. A DB approach is based on a force-displacement $F - \Delta$ law, which is different according to whether the model is ideal or has imperfections. The force-displacement function (analogous to a pushover curve) can be obtained by determining the total horizontal reaction (or base shear) F at different displacements by using simple static equilibrium principles or the principle of virtual work. In the latter most adopted case, a non-linear kinematic analysis is applied by considering kinematic varied configurations of the mechanism, in large displacements. The ideal block has a moment-rotation or force-displacement law with negative stiffness (Figure 5a) once the force F_0 at the onset of rotation has been attained.

By contrast, geometrical imperfections cause the force-displacement law to be non-linear. In this case, the block is said to be “semi-rigid”. Generally, a tri-linear simplified force-displacement law with a finite initial stiffness is adopted (Figure 5b) [42]. The real block has therefore an initial branch with a finite stiffness, a plateau phase (between Δ_1 and Δ_2) and a third branch with negative stiffness up to an ultimate displacement Δ_u (Figure 5). The ratios Δ_1 / Δ_u and Δ_2 / Δ_u are related to the material properties and the state of degradation of the mortar joints at the points of rotation. Some experimental tests were performed on URM specimens to obtain some values of those ratios [42]. A relevant parameter of the force-displacement law is the secant stiffness, which is amplitude dependent.

One method commonly used is to obtain the secant stiffness from the non-linear force-displacement curve corresponding to the point of maximum (permissible) displacement. For ductile systems, the point of maximum displacement is often associated with a point on the post-peak softening section of the non-linear force-displacement curve where the force has reduced to a fraction (generally 70–80%) of the peak force value.

For non-ductile structures such as masonry, the definition of secant stiffness is difficult due to material strength variability and a lack of yield and or unique softening points. The stiffness corresponding to a line going from the origin to the point of the curve with $\Delta = \Delta_2$ may be considered reasonably consistent [28]. A displacement-based procedure for rocking masonry structures was recently proposed [16]. It is compatible with the PERPETUATE (PERformance-based aPproach to Earthquake proTECTION of cUlturAlHeriTage in European and mediterranean countries) methodology for the seismic performance-based assessment of cultural heritage [47] and is based on the following steps: (1) definition of the rocking mechanism (by considering rigid blocks, constraints, internal and external elasto-plastic links, constructive features and masonry quality); (2) evaluation of the pushover curve, by the incremental equilibrium limit analysis performed on varied kinematic configurations; (3) definition of performance levels (PLs), in terms of displacement thresholds and related values of

the equivalent viscous damping); (4) evaluation of the capacity curve, through the conversion to an equivalent SDOF system; (5) definition of the seismic demand, in terms of an overdamped elastic acceleration-displacement response spectrum (ADRS), modified from the seismic input at the ground level in the case of local mechanisms placed at the higher levels of the structure; (6) evaluation of the values of the intensity measure (IM) that are compatible with the different PLs. It is worth noting that the displacement demand is obtained by the classical capacity spectrum method [48], through the intersection with the overdamped acceleration-displacement spectrum, without the need of defining a proper secant stiffness.

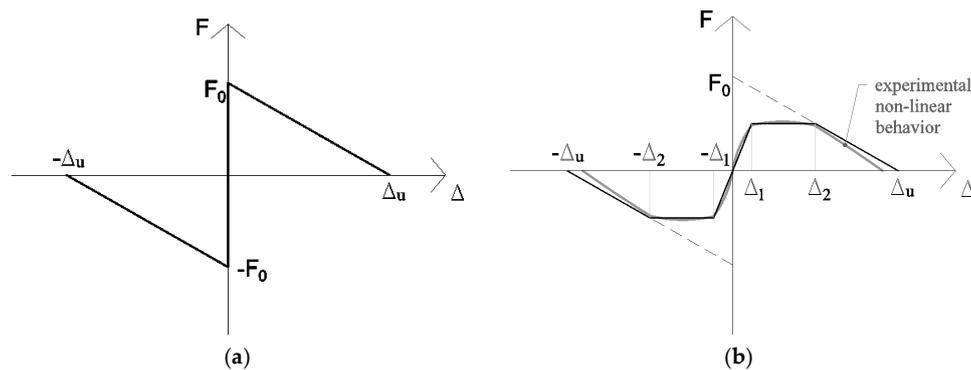


Figure 5. Force-displacement law for the rocking block: ideal model (a); model with imperfections (b).

Within this framework, some recent developments have highlighted the role of friction stresses, due to interlocking of the rocking rigid block with transverse walls, along the complete cycle of evolution of the mechanism [13,49]. In this case, a greater static multiplier is needed to activate rocking, as described in the previous section, while this effect gradually decreases after a certain displacement, due to the progressive detachment of the façade wall. The variation of the frictional forces acting along the corners can be represented by a stepwise function of the decreasing number of involved rows. In fact, the effectiveness of the frictional forces on the whole height of the corners is guaranteed as long as the first rows at the top of the sidewalls are detached from the rocking wall (first threshold displacement). In Figure 6, two examples of the pushover curve of a façade interlocked with transverse walls are reported together with that of the free wall condition. It is worth noting that the load factor can also be increased by about an order of magnitude for slender units (Figure 6b).

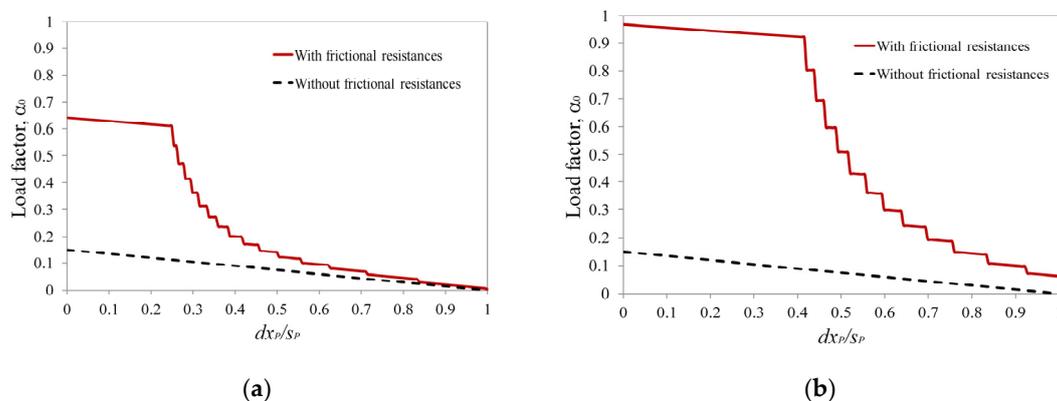


Figure 6. Influence of the frictional resistance due to interlocking between orthogonal walls after the hinge formation [13,49]. Pushover curves with unit aspect ratio $h_b/l_b = 1/3$ (a) and $h_b/l_b = 1/5$ (b).

6. Influence of Input Motion

6.1. Free-Vibrations, Harmonic Pulses and Real Accelerograms

Small changes in input or geometry (or both) can create large changes in system response for SDOF blocks subjected to random motion, at least for near collapse configurations [8,37]. As introduced in Section 2, the classical theory pointed out the main features of the model highlighting that the restoring mechanism is not elastic, but is governed by gravitational energy; consequently, the frequency of the free motion is amplitude dependent, and it is variable at each half-cycle. In other words, rocking structures do not resonate under a constant frequency, as their effective frequency depends on the response amplitude.

Casapulla and Maione [50,51] recently observed that the free-vibrations of slender rigid blocks can be represented by a succession of uniformly-decelerated motions with decreasing initial rotational velocities at each half-cycle occurring till stopping. Such a formulation was revealed to be in good agreement with Housner's model, not only with reference to the fundamental parameters governing the motion (maximum rotational angle and duration of the cycles of motion), but also with reference to the time histories of the rotational acceleration, velocity and angle.

However, while in the case of free-vibration or harmonic input forces, the response is more controllable, in the case of random motions, such as earthquakes, very large variations are found in the response of similar blocks, both in experimental and in numerical tests [21,24,52–54]. This is due to the combination of different frequency contents, duration and seismic actions parameters. Thus, in order to simplify the analysis, Housner [8] described the base acceleration as a rectangular or a half-sine impulse and derived some expressions for the minimum acceleration required to overturn the block, as a function of the duration of the impulse.

Following this pioneering work, many authors [9,55–57] examined the response of the block using both real or simulated accelerograms and harmonic loadings. These studies showed that, in contrast with the response to a single pulse, the response to more irregular but simplified accelerograms is very sensitive to the geometrical parameters of the block, as well as to the details of ground motions and the coefficient of restitution. These results allowed increased attention to various sorts of impulses or harmonic shaking and to relationships between pulses magnitude and toppling [58–61].

6.2. Critical Impulse Input and Resonance Conditions

As an alternative to standard approaches, increasing attention has been more recently focused on the identification of the worst input conditions that can imply the resonant response of the blocks [50,51,62,63].

Casapulla et al. [62,64,65] introduced an artificial limit accelerogram as a sequence of instantaneous Dirac impulses to represent the most unfavorable effects on the rocking block of the intense phase of an earthquake. A secondary sequence of intermediate impulses was also considered to reduce the resonance effects and to cover a broad range of conditions. The critical response of the block to the proposed accelerogram, represented by the achievement of the overturning, resulted in being mainly affected by the following issues: the effect of the amplitude of the main impulses; the role of the size and slenderness of the rigid block; the opposite role played by the intermediate impulses; the influence of the secondary sequence on the duration of the cycles of motion. The analysis enables, in particular, distinguishing the actions that cause overturning from those that cause the stabilization of the motion in relation with the influencing parameters.

In this line, DeJong [66] defined a theoretical accelerogram with the condition of maximization of the input energy. It deals with a step function with an alternate sign, always opposite the current rotational velocity of the block; he also assumed a constant value for the amplitude of the ground acceleration and an increasing duration of each step of the accelerogram as the half-cycle of the rocking motion increases. This reference seismic input is then transformed in a sinusoidal signal obtaining a decreasing rate of energy input and more oscillations before the overturning of the block.

The attempt to simplify the representation of the seismic signals through mathematical functions based on the selection of a few meaningful parameters has characterized in particular the research on the near-fault ground ([67,68]). Recently, Kojima et al. [69] and Kojima and Takewaki [70] developed a simple approach that schematizes the impulsive component of a near-fault earthquake as a double impulse. This simplification has also been used to represent long duration earthquakes as a multiple impulse input [69]. In order to achieve a good approximation, the amplitude of the impulse was modulated so that its maximum Fourier amplitude coincides with that of the corresponding one-cycle sinusoidal input, and a relation of proportionality between the velocity of the double impulse and the velocity amplitude of the sinusoidal waves was introduced. The proposed approach has firstly been applied to the identification of the resonant condition of elasto-plastic systems and then extended to obtain the closed form overturning limit for rigid blocks [71].

6.3. Input Motion as Pulse Type Earthquakes

Several earthquakes show dominant pulses, for which rocking structures are particularly vulnerable to pulse-type earthquakes. Zhang and Makris [59] observed the detrimental role of long-period pulses inherent in near-fault ground motions. Pulse-like earthquakes can be approximated by instantaneous Dirac impulses, as introduced in Section 6.2, or more classically, by simple trigonometric functions with specific amplitude and circular frequency. In this latter case, they can be properly idealized with cycloidal pulses or through Ricker wavelets [72]. The latter have the advantage of capturing the slight asymmetry inherent in near fault pulses better than cycloidal pulses and are expressed as [73]:

$$\psi(t) = a_p \left(1 - \frac{2\pi^2 t^2}{T_p^2} \right) e^{-\frac{1}{2} \left(\frac{2\pi^2 t^2}{T_p^2} \right)} \quad (13)$$

where a_p is the amplitude of the acceleration pulse. The value of $T_p = \frac{2\pi}{\omega_p}$ is the period that maximizes the Fourier spectrum of the wavelet, ω_p being the circular frequency of the acceleration pulse. A relevant value correlated with those parameters is the characteristic length scale of the ground excitation. This parameter, $L_p = a_p T_p^2$, gives a measure of the persistence of the most energetic pulse to generate inelastic deformation [74,75]. This parameter is more relevant than the simple acceleration amplitude a_p . Indeed, between two pulses with different acceleration amplitudes (say $a_{p1} > a_{p2}$) and different pulse durations (say $T_{p1} < T_{p2}$), the inelastic deformation does not scale with the peak pulse acceleration, but with the strongest length scale L_p [75]. By performing numerical tests on rectangular blocks of several slenderness and size, Apostolou et al. [73] found the parameters having the strongest influence on the overturning potential. Among them, there are not only the dominant frequency, but also the nature and especially the asymmetry of the seismic input. Therefore, also for masonry panels, the practice of estimating ground shaking levels by analyzing the observations of overturned and not overturned slender blocks after an earthquake is not sufficient and needs to be completed by the analysis of the mentioned parameters.

7. Types of Dynamic Analysis for Rocking Structures

7.1. Deterministic Methods Based on the Critical Rocking Response

The rocking response can be easily obtained through the integration of motion (Equations (1) and (2) for the SDOF free-standing block or Equation (3) for the horizontally-restrained block), e.g., with Runge–Kutta methods, by defining the deterministic geometric parameters and the type of excitation [76,77]. This integration allows obtaining the peak value of the normalized amplitude (also called the amplitude ratio ϑ_{max}/α), which defines the stability of the block, if the rotation is (theoretically) lower than $\pi/2$. The maximum rotation generally attained for a stable rocking block is however about 1.0. Makris et al. [78] introduced a useful representation in terms of overturning spectra of blocks, which describe the stability of a given block in terms of the maximum acceleration amplitude

as a function of the block size or the excitation frequency (Section 7.2). These kinds of spectra can also be developed to perform interesting strategies for the seismic protection of monolithic art objects, as the semi-active control proposed by Ceravolo et al. [79,80] or the addition of a tuned pendulum to the rocking block with the aim of controlling the oscillations presented by Collini et al. [81].

This deterministic method is affected by some limitations, mainly related to the need for using significant acceleration time-histories. Indeed, for a set of earthquakes, the block can result in being stable, but for others not. To avoid this situation, a probabilistic approach could be more reliable, by stochastically varying the input motion and the restitution coefficient. In deterministic methods, it can be useful to fix threshold values of amplitude ratios to define limit states, such as incipient rocking ($\frac{\vartheta_{max}}{\alpha} = 0.0$), limited rocking ($\frac{\vartheta_{max}}{\alpha} = 0.2 - 0.4$) or rocking up to collapse ($\frac{\vartheta_{max}}{\alpha} = 1$) [82]. These values are naturally generic and can be modified depending on the importance class of the element.

7.2. Rocking Structures and Elastic Oscillators in Practical Applications

The problem of the stability against overturning of rigid blocks appears still far from finding a general settlement. The rocking behavior is still a mostly unknown issue in the professional practice of civil engineers. Indeed, currently for rocking structures as masonry walls, standard response spectra are used. These demand functions derive from elastic systems, and therefore, they are not suitable for purely rocking structures. The main difficulties are still related to the description of the seismic input and the great sensitivity of the response to small variations in both system parameters and ground motion details. Priestley et al. [83] presented early experimental studies on a model suitable for slender structures and developed a practical methodology to compute displacements of the center of gravity of the structure due to rocking motion by using standard displacement and acceleration response spectra for an elastic SDOF oscillator. This was then adopted by the FEMA 356 document [84]. Makris and Konstantinidis [85] demonstrated that this methodology is oversimplified and does not take into account the fundamental differences in the dynamical structure of the SDOF systems. They showed that the rocking spectrum is a distinct and valuable intensity measure of earthquakes and offers information on the earthquake shaking that is not identifiable by the response spectrum of an elastic SDOF oscillator. Therefore, rocking structures cannot be replaced by 'equivalent' SDOF oscillators. Nevertheless, up to the formation of a plastic hinge, if the system behaves as elastic, the use of response spectra is justified and recommended. Nowadays, in the common practice of the seismic vulnerability assessment of masonry walls, the use of response spectra is predominant. In particular, according to the Italian code procedure [86], the analysis is performed in terms of displacement capacity. The value of the displacement demand is taken from the response spectra corresponding to the secant period obtained from a capacity curve (Section 5.2). The secant (also called "effective") period can also be used to predict the response of bilinear systems and rocking systems with negative stiffness, as proposed by Makris and Kampas [29]. Nevertheless, this procedure, similar to that suggested by Priestley et al. [83], completely neglects the evolution of motion over time and the behavior of a rocking block [85]. For these reasons, a pure rocking analysis is strongly recommended together with a kinematic analysis for practical applications, as well.

7.3. Methods of Analysis Based on Probabilistic Approaches

Rocking models are extremely sensitive to small variations of many parameters, such as the restitution coefficient, boundary conditions and particularly of the input motion, at least for near collapse configurations, as stated before. Aslam et al. (1980) [9] attributed the main reason for these aspects to the dependency of the vibration period upon the displacement amplitude. This influence addressed the research in studying probabilistic approaches capable of widening the results of rocking analyses in a stochastic perspective. Recently, a contribution was oriented toward defining fragility curves for rocking structures [82] by considering near-fault excitations. In particular, it was seen that by considering only peak ground acceleration (PGA) and not also, e.g., the ratio PGV/PGA (where PGV is the peak ground velocity), leads to unconservative results and not reliable results.

This study offers the conclusion that the use of bivariate intensity measures (IMs) can lead to superior fragility curves compared with conventional univariate IMs. This result confirmed previous studies related to the seismic risk [87], where through various statistical techniques, it is shown that the use of more than one IM leads to a better prediction of the damage state of a building than just a single IM. Other authors [88] performed a seismic reliability assessment of classical columns, by using synthetic ground motions that contain a high- and a low-frequency component. Generally, the demand parameter EDP (engineering demand parameter) is the maximum amplitude ratio, whereas among the most significant IMs, the peak ground velocity is one of the most reliable IM [82]. Nevertheless, the other authors proposed as EDP for columns (i) the maximum displacement at the capital normalized by the base diameter and (ii) the relative residual dislocation of adjacent drums normalized by the diameter of the corresponding drums at their interface. These assumptions can be justified by a different structural configuration, analyzed through the discrete element method (DEM). Performance levels can be then assigned to each EDP (Section 7.1), together with the values of the corresponding thresholds, to generate fragility curves. These tools are then more reliable than a simple deterministic analysis, mainly for the extremely high sensitivity of the response to the input motion, as was seen in Section 6. For these reasons, especially when uncertainties of the model are many (geometry, boundary conditions, input motion), a probabilistic analysis is recommended, by considering the univariate IMs most influencing the response, such as velocity-based parameters or bivariate IMs including acceleration and energy-based parameters.

8. Conclusive Remarks

In this paper, a wide literature review on the methods available to analyze the out-of-plane behavior of masonry walls regarded as rigid blocks subjected to seismic excitation was illustrated, mainly considering rocking and kinematic analysis. In addition, insights on recent developments were discussed, by considering the parameters most influencing the overturning potential. The main uncertainties characterizing this kind of analysis are related to the high sensitivity of the rocking blocks to small variations in system parameters and ground motion details, with possible sudden amplifications due to resonance effects. The difficulties in the analysis are increasing when the rigid block model is taken as a basic reference for the seismic analysis of the out-of-plane mechanisms either of isolated masonry blocks and walls or masonry façades poorly connected to orthogonal walls. In these cases, further uncertainties related to specific aspects of the structural behavior of masonry need to be accounted for.

Following the order in which the different methodologies were presented throughout the text, some conclusions concerning the main issues addressed could be pointed out.

As far as the rocking analysis of the free-standing SDOF block is concerned, intense research was carried out since Housner's pioneering work (1963), leading to important achievements. These have been mostly aimed at highlighting the great influence of geometric parameters and dissipation issues on the rocking motion, both analytically and experimentally. Some recent and interesting developments were also presented as more focused on the influence of boundary conditions on the dynamic response, e.g., considering masonry walls horizontally restrained by flexible diaphragms. These approaches have the potential to be extended to more complex rocking elements, e.g., involving not only simple/complex overturning, but also flexural mechanisms, where MDOF systems could be more suitable to represent the structural response.

However, the problem of the stability against overturning of rigid blocks appears still far from finding a general settlement, and the rocking behavior is a mostly unknown issue in the professional practice of civil engineers. In fact, the static force-based approach is still the most common strategy, usually preferred by practitioners. Indeed, staying in the static field, while the force-based approach can be used to analyze the onset of the rocking motion, the displacement-based approach can be recognized as a more realistic tool to assess the out-of-plane response of masonry structures under dynamic actions. Therefore, some interesting contributions focused on linear and non-linear kinematic analysis were

here illustrated as considered relevant for explaining the design and assessment techniques of masonry walls in a static perspective. Incremental kinematic approaches based on standard and non-standard limit analysis methods were discussed, including some recent developments highlighting the role of friction resistances, due to interlocking of the rocking rigid block with transverse walls, along the complete cycle of evolution of the rocking mechanism.

Moreover, the influence of the input motion parameters on the dynamic response was shown to be a crucial issue, for which stochastic analyses are recommended to obtain significant results for a wide range of ground motions of different amplitudes and frequency contents. As an alternative to probabilistic approaches, increasing attention has been more recently focused on the identification of the worst input conditions that can imply the resonant response of the blocks. Related modeling approaches have the potential to schematize the effects of the impulsive component of a near-fault earthquake or of the strong phase of an earthquake.

Lastly, a wide literature has shown that the rocking spectrum is a distinct and valuable intensity measure of earthquakes and offers information on the earthquake shaking that is not identifiable by the response spectrum of an elastic SDOF oscillator. Therefore, rocking structures cannot be replaced by 'equivalent' SDOF oscillators. On the other hand, a useful representation was given in terms of overturning spectra of blocks, which describe the stability of a given block in terms of the maximum acceleration amplitude as a function of the excitation frequency. These kinds of spectra can also be developed to perform interesting strategies for the seismic protection of monolithic art objects.

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