

Communication

Calculation of the Proportion of Free Water Molecules in Aqueous Solutions Using the Parameters of Their Dielectric Permittivity in the Terahertz Range, Based on the Onsager Theory

Nikita V. Penkov 

Institute of Cell Biophysics RAS, Federal Research Center "Pushchino Scientific Center for Biological Research of the Russian Academy of Sciences", Pushchino 142290, Russia; nvpenkov@rambler.ru

Abstract: The question of the structure of aqueous solutions is one of the most fundamental and complex, while it is relevant to all natural science disciplines. An important parameter of the dynamically equilibrium structure of an aqueous solution is the number of free water molecules. To date, there are no reliable and fully justified methods for determining this parameter. Recently, the terahertz time-domain spectroscopy (THz-TDS) method has been developing. It makes it possible to record the spectra of the complex permittivity in the THz region, where an orientation relaxation band of free water molecules is detected for aqueous solutions. The purpose of this work is to establish the relationship of the parameters of THz dielectric permittivity with the number of free water molecules. For this purpose, the process of polarization of water in the THz region was theoretically considered using the formalism of electrodynamics of continuous media. The Onsager theory is taken as a basis and its development is carried out for the case of high-frequency fields. As a result, an analytical ratio was obtained for calculating the proportion of free water molecules in solutions. A comparison with other well-known, more simplified and poorly substantiated approaches is presented. Calculations of the fraction of free molecules for a number of aqueous solutions have been carried out. It can be argued that the first theoretically justified approach to calculating the population of free water molecules in a solution, which does not contain internal contradictions, is presented.



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1. Introduction

The structure of water is one of the fundamental objects of research, which, despite the efforts of several generations of scientists, is still largely unexplored. In this context, first of all, it is worth noting the classical works of D.J. Bernal and R.H. Fowler [1], J.I. Frenkel [2], O.Ya. Samoilov [3], and I.Z. Fisher [4]. An important achievement was the understanding that the concept of “structure” of water contains not only a spatial aspect (the mutual arrangement of molecules), but also a temporal aspect. In [4] it is shown that the result of the analysis of the structure of water depends on the time of its registration. It can be instantaneous ($\ll 10^{-13}$ s), vibrational-averaged (10^{-13} – 10^{-11} s) and diffusion-averaged ($\gg 10^{-11}$ s) structure. The characteristic times may vary greatly for different structural experimental methods; therefore, the result will be different.

A statistical description of the structure of water in solution is often used in the form of decomposition of water molecules into fractions based on binding energy, mobility, and the number of hydrogen bonds. An important parameter in this context is the number of free water molecules. These are molecules that are not bound by any hydrogen bonds and are not in the ion field, which significantly affects their mobility. The proportion of free water molecules can be considered as an integral characteristic of the connectedness (or

structuredness) of water in solution [5]. This parameter was attempted to be determined in different ways, and the values obtained vary greatly. Thus, according to various data, free molecules in pure water at room temperature are less than 1% [6], about 2–3% [5], or even more than 10% [7]. Obviously, the discrepancies are related to the different specifics of the methods used and their sensitivity to the binding of water molecules.

The appearance of the terahertz time-domain spectroscopy (THz-TDS) method has opened up new possibilities in studying the structure of aqueous solutions. The energy and times corresponding to the terahertz range are characteristic of the intermolecular structure and dynamics of water. In addition, the THz-TDS method allows us to obtain not only absorption spectra, but also much more informative spectra of complex dielectric permittivity. A weakly pronounced relaxation band of water was found in the THz region [8–10], which differs from the well-known Debye relaxation band from the gigahertz region [11,12]. To date, it is understood that this high-frequency band refers to the orientation relaxation of free water molecules [5,13]. Obviously, the parameters of this band should be directly related to the number of free water molecules. Despite the technical ability to obtain high-quality spectra in the THz range, there are still no reliable approaches for calculating the number of free water molecules in solutions. This requires the development of an appropriate theory.

Due to the absence of such a theory, when calculating the proportion of free water molecules in solutions based on THz spectra, the ratio of the amplitudes of the relaxation bands of free and bound molecules is usually used [13–15]. However, this is a crude approach that does not take into account a number of important factors. For example, relaxation processes cannot be considered completely independent of other processes of molecular dynamics in an external electric field. This issue was considered in [16], which describes an approach for calculating the proportion of free water molecules, taking into account the influence of electric field shielding on their orientation polarization from faster polarization processes. However, this approach is criticized because it is based on the use of the Langevin–Debye theory, which is considered untenable for liquids with high dielectric permittivity [17], which is water.

In this paper, a theory is developed that, based on the analysis of water polarization processes in the THz frequency range, allows us to obtain the relationship of the parameters of the complex permittivity of an aqueous solution with the number of free water molecules.

2. Theoretical Background

2.1. Models of Water Polarization

A water molecule has an electric dipole moment. In an external electric field, dipoles tend to orient themselves along the field, while thermal motion disorders them. Based on such a simple consideration, it is possible to express the orientational polarizability of water molecules through a linear term of the Langevin function and associate it with the permittivity [18]:

$$\varepsilon_s - \varepsilon_\infty = \frac{Nd^2}{\varepsilon_0 3kT}, \quad (1)$$

where N is the number of molecules per unit volume, d is the electric dipole moment of a water molecule, ε_0 is the vacuum permittivity, k is the Boltzmann constant, T is the absolute temperature, ε_s is the static permittivity. The parameter ε_∞ is called the high-frequency permittivity, which describes in the sum all faster polarization processes than the orientation polarization of molecules (electronic, deformation, etc.).

The polarization of a substance is an integral characteristic associated with both the dipole moment of individual molecules and their intermolecular interaction. Since Equation (1) does not take into account the interaction of molecules in any way, it can be used only for rarefied gases, but not for an aqueous solution.

A deeper consideration of the polarization of dipole molecules in a substance requires taking into account the local field. The local field differs from the external field by the

presence of an additional polarization field. Usually, this field is called the Lorentz local field, accounting for which leads to the well-known Langevin–Debye equation [18]:

$$\frac{\varepsilon_s - 1}{\varepsilon_s + 2} = \frac{N4\pi}{3} \left(\alpha_\infty + \frac{d^2}{4\pi\varepsilon_0 3kT} \right). \quad (2)$$

Using the well-known Clausius–Mossotti equation for high-frequency polarizability:

$$\frac{\varepsilon_\infty - 1}{\varepsilon_\infty + 2} = \frac{N4\pi}{3} \alpha_\infty, \quad (3)$$

the Langevin–Debye equation can be rewritten in the following form:

$$\frac{\varepsilon_s - \varepsilon_\infty}{(\varepsilon_s + 2)(\varepsilon_\infty + 2)} = \frac{Nd^2}{\varepsilon_0 27kT}. \quad (4)$$

As it was later shown, Equation (4) contains a fundamental defect called the Mossotti catastrophe [17], since it leads to ferroelectric properties of water at room temperature, which is not observed in reality. This approach is considered applicable for substances with a small dielectric constant. Nevertheless, it is still often taken as the basis for describing the polarization of water.

Further development of the theory of water polarization was made by Onsager [19], showing that the inconsistency of the Langevin–Debye approach lies in the inaccuracy of the expression for the Lorentz local field. In Onsager’s theory, a water molecule is considered as a spherical object with a permittivity ε_∞ , having a dipole moment \vec{d} . The molecule is surrounded by a continuous medium with a dielectric constant ε_s . The theory takes into account not only the influence of the polarization of the environment on the dipole, but also the fact that the polarization of the environment itself changes under the influence of the dipole field. This change in the environment leads to the appearance of an additional member of the local field, which also affects the polarization of the dipole molecule. This additional field was called the Onsager reaction field. The Onsager formula can be written as follows:

$$\frac{(\varepsilon_s - \varepsilon_\infty)(2\varepsilon_s + \varepsilon_\infty)}{\varepsilon_s(\varepsilon_\infty + 2)^2} = \frac{Nd^2}{\varepsilon_0 9kT}. \quad (5)$$

The most developed theory of polarization of polar liquids is the Kirkwood theory [20], which takes into account the correlations of the mutual orientation of the dipole moments of molecules:

$$\frac{(\varepsilon_s - \varepsilon_\infty)(2\varepsilon_s + \varepsilon_\infty)}{\varepsilon_s(\varepsilon_\infty + 2)^2} = \frac{Nd^2}{\varepsilon_0 9kT} g. \quad (6)$$

The Kirkwood Equation (6) differs from the Onsager Equation (5) only by the coefficient g , which is determined based on models of the near-order structure of water or experimentally. Its value for water at room temperature is approximately 2.5 [21]. This theory is considered to be the most justified in describing the polarization of water and allows us to obtain an almost perfect correspondence between the theoretical and experimental values of ε_s .

2.2. Polarization in an Alternating Electric Field

So far, the polarization of water in an electrostatic field has been considered, and the equations for ε_s have been given. In the electrodynamics of continuous media, polarization in alternating fields is often considered. In this case, the static permittivity ε_s is replaced by the complex permittivity $\varepsilon^*(\omega)$, which depends on the frequency ω . The function $\varepsilon^*(\omega)$ makes it possible to describe the polarization of matter at different frequencies, including in the dispersion regions where the absorption of the electromagnetic field is observed. The

form of all equations obtained for the electrostatic field remains unchanged for the case of alternating fields under certain conditions [22,23]. For example, the Langevin–Debye Equation (2) for an electrostatic field passes into the Clausius–Mossotti Equation (3) or the Lorentz–Lorentz equation [24], describing polarization in the infrared and optical frequency ranges. The form of the equations remains the same, only the contribution of the orientation polarization of the molecules disappears due to the fact that at such high frequencies the molecules do not have time to reorient themselves.

2.3. Complex Permittivity of Water in the THz Range

The dielectric characteristics of water in the THz range are determined by several reliably established molecular dynamics processes. Firstly [25], it is the orientational relaxation of bound water molecules (Debye relaxation) with characteristic times of ~10 ps and the orientational relaxation of free water molecules (sometimes called “fast relaxation”) at ~0.3 ps. Intermolecular oscillations also appear on the high-frequency side of the THz range [26] and some others. With this in mind, the permittivity of water in the THz range can be described in a standard way for dielectric spectroscopy [9,14]:

$$\varepsilon^* = \frac{\Delta\varepsilon_1}{1 - i\omega\tau_1} + \frac{\Delta\varepsilon_2}{1 - i\omega\tau_2} + \varepsilon_\infty + i\frac{\sigma_0}{\varepsilon_0\omega}, \quad (7)$$

where $\tau_{1,2}$ is the relaxation time of bound and free water molecules; $\Delta\varepsilon_{1,2}$ is the contribution of the corresponding relaxation processes to the overall dielectric response; ε_∞ is the high-frequency dielectric constant; ω is the cyclic frequency; i is an imaginary unit, and the term $i\frac{\sigma_0}{\varepsilon_0\omega}$ describes the dielectric losses associated with ionic conductivity (σ_0 is dc-conductivity, ε_0 is the vacuum permittivity). If necessary, additional terms describing high-frequency processes of molecular dynamics [26,27], which in this case are included in ε_∞ , can be explicitly included in Equation (7).

The calculation of the parameters of the model complex permittivity (7) of aqueous solutions can be performed on the basis of fitting experimental spectra. The spectra of the complex permittivity can be measured using the terahertz time-domain spectroscopy (THz-TDS) method, the principles of which are described, for example, in [28]. This method makes it possible to simultaneously determine the transmission spectra $Tr(\nu)$ and the refractive index $n(\nu)$. From these two spectra, the spectra of the real and imaginary parts of the complex permittivity can be calculated using the following equations:

$$\varepsilon'(\nu) = n^2(\nu) - \left[\frac{\ln Tr(\nu)}{4\pi\nu l} \right]^2, \quad \varepsilon''(\nu) = \frac{-n(\nu)\ln Tr(\nu)}{2\pi\nu l} \quad (8)$$

where ε' and ε'' are the real and imaginary parts of complex permittivity, l is sample thickness, ν is wavenumber ($\nu = \omega/2\pi c$).

2.4. Methods for Calculating the Proportion of Free Water Molecules in Aqueous Solutions Based on Permittivity Parameters in the THz Range

When calculating the proportion of free water molecules using the permittivity parameters of a solution in the THz range, as a rule, an approach based on the ratio of the amplitudes of the relaxation bands $\Delta\varepsilon_1$ and $\Delta\varepsilon_2$ is used, or a similar approach with minor variations [13–15]:

$$n = \frac{\Delta\varepsilon_2}{(\Delta\varepsilon_1 + \Delta\varepsilon_2)} \quad (9)$$

The number of free water molecules is considered to be proportional $\Delta\varepsilon_2$, and the total number is proportional $\Delta\varepsilon_1 + \Delta\varepsilon_2$. Equation (9) implies that the process of orientation polarization of free water molecules does not depend on other polarization processes, which is not quite true. In [16], the orientational polarization of free water molecules was considered on the basis of the combined Langevin–Debye and Clausius–Mossotti theory, and the need to take into account the shielding of the field by faster polarization

processes was shown. The relationship of the proportion of free water molecules with the parameters $\Delta\epsilon_2$ and $\Delta\epsilon_1$, as well as with the parameters of high-frequency processes (bending, stretching intermolecular modes, etc.) which, however, can be combined into one parameter ϵ_∞ , is established:

$$n = \frac{3\Delta\epsilon_2}{(\Delta\epsilon_2 + \epsilon_\infty + 2)(\epsilon_\infty + 2)} * \frac{\epsilon_0 9kT}{Nd^2}, \quad (10)$$

where the parameters' definitions are the same as in Equations (1) and (7). This approach, however, is imperfect, since the underlying Langevin–Debye theory, as applied to the polarization of water, leads to the above-mentioned Mossotti catastrophe. As a result, despite the existence of reasonable aspects of this theory, it is unclear what the error of calculations according to Equation (10) is.

3. Calculation of the Proportion of Free Water Molecules in Aqueous Solutions Based on the Permittivity Parameters in the THz Range Using the Onsager Theory

As shown in Section 2.3, several types of molecular dynamics can be distinguished in water, described by Equation (7): orientation relaxation of bound molecules, orientation relaxation of free molecules and other high-frequency processes. Each of these types of molecular dynamics is associated with a corresponding type of polarization of water: orientation polarization of bound molecules, orientation polarization of free molecules and other, faster, polarization processes. Each type of polarization is realized for its characteristic times and gives a certain contribution to the overall dielectric response. Therefore, as is known, static permittivity is determined by the sum of the contributions of all polarization processes:

$$\epsilon_s = \Delta\epsilon_1 + \Delta\epsilon_2 + \epsilon_\infty \quad (11)$$

Despite the additivity of the polarization processes for the electrostatic case, considering each polarization process separately at its characteristic times requires, as mentioned above, taking into account the shielding from faster polarization processes. For the orientational polarization of free water molecules, all faster processes are described by the value ϵ_∞ . At the same time, the orientation polarization of bound molecules is much slower and does not have time to shield the external electric field during the orientation polarization of free molecules. Indeed, at the peak of the orientation relaxation band of free molecules (at the frequency $\omega_2 = \frac{1}{\tau_2}$), the contribution of the polarization of bound molecules is $\frac{\Delta\epsilon_1}{[1+(\tau_1/\tau_2)^2]} \approx 0.09$ with the total permittivity of water at this frequency $Re[\epsilon^*(\omega_2)] \approx 5.5$ ($\tau_1 = 8.3$ ps [27], $\tau_2 = 0.3$ ps and $\Delta\epsilon_1 \approx 70$ [29]).

Consider the process of polarization of water at a frequency of $1/\tau_1 \ll \omega_0 \ll 1/\tau_2$ ($\omega_0 \approx 0.6$ THz), that is, on a spectral scale between two relaxation bands of water (Figure 1). At this frequency, the orientation polarization of free water molecules is almost completely realized, since $\tau_2 \ll 1/\omega_0$, as well as the polarization of all faster processes described by ϵ_∞ . At the same time, the polarization of bound water molecules practically does not have time to be realized, since $\tau_1 \gg 1/\omega_0$.

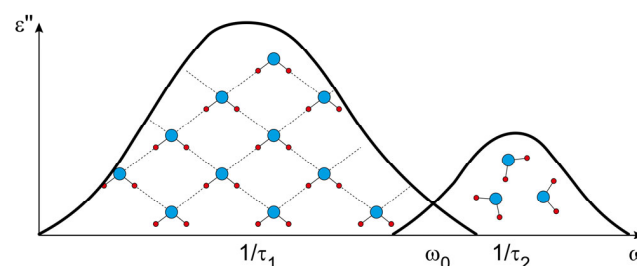


Figure 1. Schematic description of the location of orientation relaxation bands of bound and free water molecules. ω_0 is the selected frequency at which the process of orientational polarization of free water molecules is considered.

By analogy with the static permittivity (11), at a frequency of ω_0 , the permittivity of an aqueous solution can be represented as the sum of $\varepsilon(\omega_0) \approx \Delta\varepsilon_2 + \varepsilon_\infty$. To describe the orientation polarization of free water molecules at the frequency ω_0 , we will then use the Onsager Equation (5), where instead of the static permeability ε_s , we need to substitute $\Delta\varepsilon_2 + \varepsilon_\infty$. In addition, taking into account the selected frequency domain, where the orientation polarization of only free molecules is realized, it is necessary, instead of the total number of molecules N , to substitute the number of free water molecules N_{free} (in unit volume):

$$\frac{\Delta\varepsilon_2(2\Delta\varepsilon_2 + 3\varepsilon_\infty)}{(\Delta\varepsilon_2 + \varepsilon_\infty)(\varepsilon_\infty + 2)^2} = \frac{N_{free}d^2}{\varepsilon_0 9kT} \quad (12)$$

As a rule, it is not the absolute number of free water molecules that is of interest, but their fraction n of the total number:

$$n = \frac{N_{free}}{N} \quad (13)$$

Based on Equations (12) and (13), we obtain an expression for calculating the proportion of free water molecules:

$$n = \frac{\Delta\varepsilon_2(2\Delta\varepsilon_2 + 3\varepsilon_\infty)}{(\Delta\varepsilon_2 + \varepsilon_\infty)(\varepsilon_\infty + 2)^2} * \frac{\varepsilon_0 9kT}{Nd^2}, \quad (14)$$

where the parameters' definitions are the same as in Equations (1) and (7).

4. Discussion

Equation (14) for calculating the proportion of free water molecules in solution was obtained on the basis of theoretical consideration of the process of orientation polarization of free molecules with the involvement of the Onsager theory. Despite the fact that the Onsager theory has been substantially modified by us for application in the case of high-frequency fields, its main aspects have remained unchanged. Firstly, water molecules are considered not to interact with each other in any way other than Coulomb, that is, there are no correlations of the relative location of neighboring molecules. Secondly, a water molecule is considered to be a sphere filled with a dielectric medium with a permittivity of ε_∞ .

As is known, not taking into account the correlations of the relative position of water molecules leads to an underestimation of the magnitude of the static permittivity of water, which was corrected in Kirkwood's theory [20] by the introduction of the correlation g-factor (6). However, when considering the orientation polarization of the fraction of free water molecules at their characteristic times, the Onsager approach turns out to be the most appropriate. Taking into account the Kirkwood g-factor in this case would be erroneous, since it refers to bound water molecules whose orientation polarization is not realized at the considered times.

The spherical symmetry of the water molecule in the Onsager approach may also seem like a rough approximation. However, when considering the processes of polarization, it is not the physical sphericity that is important, but the spherical symmetry of the polarizability of the water molecule. In this case, this is quite justified, since it has been experimentally shown that the polarizability of a water molecule at THz frequencies is close to isotropic [30].

Thus, the Onsager theory, which is unacceptable for describing the polarization of water in the electrostatic case, after appropriate modification turned out to be much more suitable for describing the orientation polarization of free molecules in water at terahertz frequencies than other well-known Langevin–Debye or Kirkwood models.

Table 1 presents a comparison of the results of calculating the number of free water molecules in various solutions using the three approaches described above: according to Equations (9) and (10) and according to Equation (14) obtained in this work.

Table 1. Calculation of the proportion of free water molecules (in % of the total number) using three approaches: based on the ratio of the amplitudes of relaxation processes (9), based on the Langevin–Debye theory (10), based on the theory developed in this paper (14).

Aqueous Solution	Calculation by the Equation (9)	Calculation by the Equation (10) ¹	Calculation by the Equation (14) ¹
CaCl ₂ 0.5 M solution at 25 °C [5]	2.01	1.81	1.99
KCl 1 M solution at 25 °C [5]	2.41	2.12	2.38
CsCl 1 M solution at 25 °C [5]	2.81	2.52	2.90
BSA 48 mg/mL solution at pH = 6, 25 °C [31]	3.75	3.99	5.00
Suspension of DPPC liposomes 20 mg/mL in the liquid crystalline phase at 45 °C [32]	2.72	3.10	3.64
DNA solution 25 mg/mL at 25 °C [33]	4.77	4.01	5.14
Glucose solution 50 mg/mL at 25 °C [34]	4.50	3.90	4.95
Water at 25 °C [34]	3.76	3.78	4.72

¹ $d = 6.187 \times 10^{-30} \text{ C} \cdot \text{m}$ [35], $N = 6 \times 10^{23} \times 55.56 \times 10^3 \text{ molecules/m}^3$.

It follows from the data in Table 1 that simplified approaches (9) and (10), as a rule, underestimate the number of free water molecules. Interestingly, the seemingly crude approach (10), based on the Langevin–Debye theory, underestimates the number of free water molecules by only 15–20% relative to the values obtained in this work.

The results of calculating the proportion of free water molecules based on the theory developed in this paper are consistent with modern concepts of hydration in aqueous solutions. In the Hoffmeister series for Ca²⁺, K⁺, Cs⁺ cations, an increase in the number of free water molecules is observed. This is explained by a decrease in the charge density of the ion on the surface, which leads to a weakening of the binding of water in the solution. In fact, the increase in the number of free molecules in this series is even greater if we take into account the effect of water displacement by ions, since the total volume of ions in the solution increases in this series. The presence of protein, sugar and DNA in the solution leads to a slight increase in the proportion of free molecules. This phenomenon is observed for the vast majority of biomolecules during their hydration [29,31,36]. The explanation of this fact was first proposed in [31]. Energetically strong binding of water molecules to hydrophilic sites of biomolecules in primary hydrate shells leads to distortion of the structure of water and does not allow the formation of undisturbed (“normally bound”) structure of water outside the primary hydrate layer for steric reasons. At the same time, beyond the primary hydrate layer, there is no longer an energetically strong binding of water molecules. That is, the hydrate shell of the biomolecule, roughly speaking, consists of two layers: a strongly bound primary and a partially destroyed secondary, outside of which the structure of undisturbed water is formed. It should be particularly noted that the calculation of the proportion of free water molecules in the solutions of biomolecules indicated in Table 1 was carried out on the basis of the dielectric permittivity of the aqueous phase of these solutions. The dielectric contribution of the biomolecules themselves was excluded from the experimental permittivity of solutions using effective medium models [22,37] and did not affect the result.

5. Conclusions

A consistent theory of orientation polarization of the fraction of free water molecules in an aqueous solution is presented. This theory is based on the theory of Onsager polarization, modified for use in the case of high-frequency fields. Based on the presented theory and experimental spectra of the complex permittivity in the THz range, it is possible to calculate the proportion of free water molecules in any aqueous solution. The calculations show that the previously known similar approaches are greatly simplified and underestimate the number of free water molecules.

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