



Article The Dynamical Analysis of a Biparametric Family of Six-Order Ostrowski-Type Method under the Möbius Conjugacy Map

Xiaofeng Wang * D and Xiaohe Chen

School of Mathematical Sciences, Bohai University, Jinzhou 121000, China; 2020010009@qymail.bhu.edu.cn

* Correspondence: wangxiaofeng@qymail.bhu.edu.cn

Abstract: In this paper, a family of Ostrowski-type iterative schemes with a biparameter was analyzed. We present the dynamic view of the proposed method and study various conjugation properties. The stability of the strange fixed points for special parameter values is studied. The parameter spaces related to the critical points and dynamic planes are used to visualize their dynamic properties. Eventually, we find the most stable member of the biparametric family of six-order Ostrowski-type methods. Some test equations are examined for supporting the theoretical results.

Keywords: Möbius conjugacy; strange fixed points; stability; free critical points; parameter space; dynamic plane

MSC: 65H05; 65B99



Citation: Wang, X.; Chen, X. The Dynamical Analysis of a Biparametric Family of Six-Order Ostrowski-Type Method under the Möbius Conjugacy Map. *Fractal Fract.* 2022, *6*, 174. https://doi.org/ 10.3390/fractalfract6030174

Academic Editors: Ángel Alberto Magreñán, Iñigo Sarría and Ravi P. Agarwal

Received: 21 December 2021 Accepted: 18 March 2022 Published: 21 March 2022

Publisher's Note: MDPI stays neutral with regard to jurisdictional claims in published maps and institutional affiliations.



Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/).

1. Introduction

In many fields including mathematics, physics and engineering, the problem of solving nonlinear equations f(x) = 0 is involved where $f : D \subseteq R \rightarrow R$ is a differentiable function defined in an open interval D. It is almost impossible to find the exact solutions of nonlinear equations, so an iterative method is used to obtain the as accurate as possible approximate solutions. The most classic iterative scheme is Newton's iterative method [1]. In recent years, many researchers have devoted themselves to designing a class of high-order iterative methods to solve nonlinear equations by using different techniques to improve Newton's scheme. Rational linear function method [2], combination method [3] and weight-function method [4] are often used to generate new iterative methods or to improve the convergence order of known methods.

Ostrowski's method [5] is the optimal fourth-order method under the Kung–Traub's conjecture [6]. A variant of Ostrowski-type methods using the weight-function procedure is proposed by Chun [7], which is given by

$$\begin{cases}
y_n = x_n - \frac{f(x_n)}{f'(x_n)}, \\
w_n = y_n - \frac{f(x_n)}{f(x_n) - 2f(y_n)} \frac{f(y_n)}{f'(x_n)}, \\
x_{n+1} = w_n - H(u_n) \frac{f(w_n)}{f'(x_n)},
\end{cases}$$
(1)

where $u_n = \frac{f(y_n)}{f(x_n)}$ and $H(u_n)$ represents a weight function. In view of Theorem 1 of Reference [6], if function $H(u_n)$ satisfies the conditions H(0) = 1, H'(0) = 2, and $|H''(0)| < \infty$, then the above iteration scheme is sixth-order. In the current study, we will select the concrete form of $H(u_n)$ with two-parameter $t, k \in C$: $H(u_n) = \frac{t+(2t+k)u_n}{t+ku_n}$.

Hence, we obtain the following biparametric family of sixth-order method (OM):

$$\begin{cases} y_n = x_n - \frac{f(x_n)}{f'(x_n)}, \\ w_n = y_n - \frac{f(x_n)}{f(x_n) - 2f(y_n)} \frac{f(y_n)}{f'(x_n)}, \\ x_{n+1} = w_n - \frac{tf(x_n) + (2t+k)f(y_n)}{tf(x_n) + kf(y_n)} \frac{f(w_n)}{f'(x_n)}. \end{cases}$$
(2)

The fixed point operator of (2) or iteration function is:

$$R_f(z) = z - \frac{f(z)}{f'(z)} - \frac{f(z)}{f(z) - 2f(y)} \frac{f(y)}{f'(z)} - \frac{tf(z) + (2t+k)f(y)}{tf(z) + kf(y)} \frac{f(w)}{f'(z)},$$
(3)

where:

$$y = z - \frac{f(z)}{f'(z)}, w = y - \frac{f(z)}{f(z) - 2f(y)} \frac{f(y)}{f'(z)}.$$
(4)

For simplicity, we mark *w* as w(z).

For arbitrary members (except for t = 0 and k = 0) of the OM family (2), the speed of convergence is similar. Our main interest is the stability of the OM family in this paper. Since the dynamic properties of the iterative method provide us with important information about the stability and reliability of members. In these respects, several scholars described the dynamical behavior of different iterative families (see [4,6,8–11]). For example, Maroju et al. [11] analyzed the dynamical behavior of Chebyshev–Halley family. Magreñán [12] studied some anomalies in the fourth-order Jarrat family. These iterative methods are all one-parameter iterative schemes. We will investigate the dynamic properties of rational functions $R_f(z)$ on low-degree polynomials associated with the biparametric family (2) with the help of complex dynamics tools [13]. This analysis allows us to reject poorly behaved members and select the most stable ones.

Now, we review some basic dynamical concepts (see [8,9,14]) which will be used in paper. Given a rational function $R : \overline{C} \to \overline{C}$, where \overline{C} is the Riemann sphere, the orbit of a point $z_0 \in \overline{C}$ is defined as

$$\Big\{z_0, R(z_0), R^2(z_0), \cdots, R^n(z_0), \cdots\Big\}.$$

Moreover, a fixed point z_0 can be divided into the following four types:

(a) Attractor if
$$|R'(z_0)| < 1$$
.

- (b) Superattractor if $|R'(z_0)| = 0$.
- (c) Repulsor if $|R'(z_0)| > 1$.
- (d) Parabolic if $|R'(z_0)| = 1$.

The basin of attraction of an attractor z^* is defined as

$$\mathcal{A}(z^*) = \{ z_0 \in \overline{C} : R^n(z_0) \to z^*, n \to \infty \}.$$

This article is divided into seven sections. Section 2 briefly describes the preliminary results of the dynamics on \bar{C} . In Section 3, we obtain the strange fixed points of $R_f(z)$ and investigate its stability. In Section 4, the relevant properties of the free critical points are described. Section 5 presents the parameter space associated with free critical points and the dynamic plane of family (2) elements. In Section 6, we conduct numerical experiments on the proposed method and the existing methods. In Section 7, we make a short summary.

2. Preliminary Results

Theorem 1 (The Scaling Theorem). Let f(z) be an analytic function on the Riemann sphere, and let $\Gamma(z) = az + b$, $a \neq 0$, be an affine map. If $h(z) = f \circ \Gamma(z)$, then the fixed points operator (3) R_f is analytically conjugated to R_h by Γ , that is $\Gamma \circ R_h \circ \Gamma^{-1}(z) = R_f(z)$.

Proof. With the iteration function R(z), we have:

$$R_{h}(\Gamma^{-1}(z)) = \Gamma^{-1}(z) - \frac{h(\Gamma^{-1}(z))}{h'(\Gamma^{-1}(z))} - \frac{h(\Gamma^{-1}(z))}{h(\Gamma^{-1}(z)) - 2h(\Gamma^{-1}(z) - \frac{h(\Gamma^{-1}(z))}{h'(\Gamma^{-1}(z))})} \frac{h(\Gamma^{-1}(z) - \frac{h(\Gamma^{-1}(z))}{h'(\Gamma^{-1}(z))})}{h'(\Gamma^{-1}(z))} - \frac{th(\Gamma^{-1}(z) + (k+2t)h(\Gamma^{-1}(z) - \frac{h(\Gamma^{-1}(z))}{h'(\Gamma^{-1}(z))})}{th(\Gamma^{-1}(z) + kh(\Gamma^{-1}(z) - \frac{h(\Gamma^{-1}(z))}{h'(\Gamma^{-1}(z))})} \frac{h(w(\Gamma^{-1}(z)))}{h'(\Gamma^{-1}(z))}.$$
(5)

since:

$$h \circ \Gamma^{-1}(z) = f(z), \tag{6}$$

$$(h \circ \Gamma^{-1})'(z) = \frac{1}{a}h'(\Gamma^{-1}(z)), \tag{7}$$

we obtain:

$$h'(\Gamma^{-1}(z)) = a(h \circ \Gamma^{-1})'(z) = af'(z),$$
(8)

and

$$h''(\Gamma^{-1}(z)) = a^2 f''(z).$$
(9)

Using (6) and (8), we obtain:

$$R_{h}(\Gamma^{-1}(z)) = \Gamma^{-1}(z) - \frac{f(z)}{af'(z)} - \frac{f(z)}{f(z) - 2h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})} \frac{h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})}{af'(z)} - \frac{tf(z) + (k+2t)h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})}{tf(z) + kh(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})} \frac{h(w(\Gamma^{-1}(z)))}{af'(z)}.$$
(10)

Therefore:

$$\begin{split} R_{f}(z) &= \Gamma \circ R_{h} \circ \Gamma^{-1}(z) = \Gamma(R_{h}(\Gamma^{-1}(z))) = aR_{h}(\Gamma^{-1}(z)) + b \\ &= a\Gamma^{-1}(z) - \frac{f(z)}{f'(z)} - \frac{f(z)}{f(z) - 2h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})} \frac{h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})}{f'(z)} \\ &- \frac{tf(z) + (k + 2t)h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})}{tf(z) + kh(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})} \frac{h(w(\Gamma^{-1}(z)))}{f'(z)} + b \\ &= z - \frac{f(z)}{f'(z)} - \frac{f(z)}{f(z) - 2h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})} \frac{h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})}{f'(z)} \\ &- \frac{tf(z) + (k + 2t)h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})}{tf(z) + kh(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})} \frac{h(w(\Gamma^{-1}(z)))}{f'(z)}. \end{split}$$

We then just need to prove that $h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)}) = f(y)$ and $h(w(\Gamma^{-1}(z))) = f(w(z))$. First prove that $h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)}) = f(y)$. By using the Taylor expansion of $h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})$ about $\Gamma^{-1}(z)$ and (9), we have:

$$h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)}) = h(\Gamma^{-1}(z)) - h'(\Gamma^{-1}(z)) \frac{f(z)}{af'(z)} + h''(\Gamma^{-1}(z)) \frac{f(z)^2}{2a^2 f'(z)^2} -h'''(\Gamma^{-1}(z)) \frac{f(z)^3}{6a^3 f'(z)^3} + \cdots = f''(z) \frac{f(z)^2}{2f'(z)^2} - f'''(z) \frac{f(z)^3}{6f'(z)^3} + \cdots = f(z - \frac{f(z)}{f'(z)}) = f(y).$$
(12)

The prove that $h(w(\Gamma^{-1}(z))) = f(w)$. Because:

$$h(w) = h\left(z - \frac{h(z)}{h'(z)} - \frac{h(z)}{h(z) - 2h(z - \frac{h(z)}{h'(z)})} \frac{h(z - \frac{h(z)}{h'(z)})}{h'(z)}\right).$$
(13)

Replacing *z* with $\Gamma^{-1}(z)$ in (13), we obtain:

$$\begin{split} h(w(\Gamma^{-1}(z))) &= h(\Gamma^{-1}(z) - \frac{h(\Gamma^{-1}(z))}{h'(\Gamma^{-1}(z))} - \frac{h(\Gamma^{-1}(z))}{h(\Gamma^{-1}(z)) - 2h(\Gamma^{-1}(z) - \frac{h(\Gamma^{-1}(z))}{h'(\Gamma^{-1}(z))})} \frac{h(\Gamma^{-1}(z) - \frac{h(\Gamma^{-1}(z))}{h'(\Gamma^{-1}(z))})}{h'(\Gamma^{-1}(z))}) \\ &= h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)} - \frac{f(z)}{f(z) - 2h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})}) \frac{h(\Gamma^{-1}(z) - \frac{f(z)}{af'(z)})}{af'(z)}) \\ &= h(\frac{z-b}{a} - \frac{f(z)}{af'(z)} - \frac{f(z)}{f(z) - 2f(y)} \frac{f(y)}{af'(y)}) \\ &= h(\frac{1}{a}(y - \frac{f(z)}{f(z) - 2f(y)} \frac{f(y)}{af'(z)} - b)) \\ &= f(1) + \frac{f(z)}{f(z) - 2f(y)} \frac{f(y)}{af'(z)} - b)) \\ &= f(y - \frac{f(z)}{f(z) - 2f(y)} \frac{f(y)}{af'(z)}) \\ &= f(w). \end{split}$$
(14)

From (14), we obtain $\Gamma \circ R_h \circ \Gamma^{-1}(z) = R_f(z)$; then this proof is completed. \Box

The above theorem shows that it allows us to conjugate the dynamical behavior of one operator with the behavior related to another, conjugated by an affine application.

Definition 1. Let $\mathcal{F} : X \subset \overline{C} \to X$ and $\mathcal{O} : Y \subset \overline{C} \to Y$ be two functions (representing two dynamical systems). We say that \mathcal{F} is conjugate to \mathcal{O} via Γ if there exists an isomorphism $\Gamma : Y \to X$ such that $\mathcal{F} \circ \Gamma = \Gamma \circ \mathcal{O}$. Such a map Γ is called a conjugacy [15].

According to Theorem 1 which was proven in [4], we find:

Theorem 2. Let \mathcal{F} and \mathcal{O} be defined by Definition 1 be of class C^1 and are conjugate to each other via the diffeomorphic conjugacy Γ . Moreover, let ξ be a fixed point of \mathcal{O} . Then, the following hold:

(a) The fixed points property remains invariant under a topological conjugacy Γ , that is:

 $\xi = \mathcal{O}(\xi)$ if and only if $\Gamma(\xi) = \mathcal{F}(\Gamma(\xi))$.

(b) The Poincaré characteristic multiplier [10] of ξ by \mathcal{O} , denoted by $m(\mathcal{O}, \xi)$, is invariant under a diffeomorphic conjugacy Γ , that is:

$$m(\Gamma \circ \mathcal{O} \circ \Gamma^{-1}, \Gamma(\xi)) = \mathcal{F}'(\Gamma(\xi)) = \mathcal{O}'(\xi) = m(\mathcal{O}, \xi).$$

Remark 1. Results of Theorem 2 state that a conjugacy Γ indeed preserves the dynamical behavior between the two dynamical systems; for example, if \mathcal{F} and \mathcal{O} are conjugate to each other via Γ , and ξ is a fixed point of \mathcal{O} , then $\Gamma(\xi)$ is a fixed point of \mathcal{F} . The converse is also true, that is, let ξ be a fixed point of \mathcal{F} , then $\Gamma^{-1}(\xi)$ is a fixed point of \mathcal{O} ; Furthermore, ξ is a critical point or free critical point of \mathcal{O} , then $\Gamma(\xi)$ is a critical point or free critical point of \mathcal{F} . The converse is also true. The above is based on the invariance properties of the fixed point and its multiplier under the conjugacy.

Furthermore, we discover $\mathcal{F} = \Gamma \circ \mathcal{O} \circ \Gamma^{-1}$ and $\mathcal{F}^n = (\Gamma \circ \mathcal{O} \circ \Gamma^{-1}) \circ (\Gamma \circ \mathcal{O} \circ \Gamma^{-1}) \cdots \circ (\Gamma \circ \mathcal{O} \circ \Gamma^{-1}) = \Gamma \circ \mathcal{O}^n \circ \Gamma^{-1}$. If \mathcal{F} and \mathcal{O} are extra invertible, we can also find $\mathcal{F}^{-1} = \Gamma \circ \mathcal{O}^{-1} \circ \Gamma^{-1}$ and $\mathcal{F}^{-n} = \Gamma \circ \mathcal{O}^{-n} \circ \Gamma^{-1}$, besides the topological conjugacy, Γ maps an orbit:

$$\cdots, \mathcal{O}^{-2}(y), \mathcal{O}^{-1}(y), y, \mathcal{O}(y), \mathcal{O}^{2}(y), \cdots$$

of \mathcal{O} onto an orbit:

$$\cdot, \mathcal{F}^{-2}(x), \mathcal{F}^{-1}(x), x, \mathcal{F}(x), \mathcal{F}^{2}(x), \cdots$$

of \mathcal{F} , where $x = \Gamma(y)$. From this, we find that the order of points is preserved. Thus, the orbits of the two maps behave similarly under homeomorphism Γ . According to the invariant properties of the fixed point, the multiplier as well as the scaling theorem, it is assuredly of value to study the dynamics of a conjugated map if simplified through conjugacy Γ .

3. Fixed Points and Stability

Due to topological invariance, we transform R_f in (3) into Ω by a linear fractional Möbius conjugacy map $M(z) = \frac{z-a}{z-b}$, $(a \neq b)$ fulfilling:

$$\Omega(z;a,b,t,k) = \frac{I(z;a,b,t,k)}{\Theta(z;a,b,t,k)},$$
(15)

when applied to an arbitrary prototype quadratic polynomial f(z) = (z - a)(z - b), where *I* and Θ are rational polynomials whose coefficients are dependent upon parameters *a*, *b*, *t*, *k*. One of our aims is to make coefficients of both *I* and Θ be minimally dependent on parameters. According to the previously mentioned conjugate, we find that all the coefficients of both *I* and Θ are only dependent upon *t* and *k*, and independent of *a* and *b*. Consequently, with the aid of the symbolic operation capability of Mathematica [16], (15) can be written as

$$\Omega(z;t,k) = -z^4 \frac{\sigma(z;t,k)}{\mu(z;t,k)},\tag{16}$$

where:

$$\sigma(z;t,k) = -k - 2t + (k+2t)z + 2tz^{2} + (k+2t)z^{3} + tz^{4}, \qquad (17)$$

$$\mu(z;t,k) = -t - (k+2t)z - 2tz^2 - (k+2t)z^3 + (k+2t)z^4,$$
(18)

We then apply a special treatment to t, the cases for values $t \in \left\{\frac{1}{2}, -\frac{1}{2}, 1, -1\right\}$. The aim is to simplify the rational function also to reduce the dependence on the parameters even further. In this study, the more interesting one is t = 1, which corresponds to the fixed point operator of the following form:

$$\Omega(z;k) = \Omega(z;1,k) = \frac{z^4 \left[-2 - k + (2+k)z + 2z^2 + (2+k)z^3 + z^4\right]}{1 + (2+k)z + 2z^2 + (2+k)z^3 - (2+k)z^4}.$$
(19)

We will investigate the fixed points of Ω and their stability. The fixed points of $\Omega(z;k)$ are given by the roots of:

$$\Omega(z;k) - z = \frac{z(z-1)\left[1 + (3+k)z + (5+k)z^2 + 3(3+k)z^3 + (5+k)z^4(3+k)z^5 + z^6\right]}{1 + (2+k)z + 2z^2 + (2+k)z^3 - (2+k)z^4} = \frac{z(z-1)\rho(z)}{\zeta(z)},$$
(20)

where we handily denote:

$$\rho(z) = 1 + (3+k)z + (5+k)z^2 + 3(3+k)z^3 + (5+k)z^4(3+k)z^5 + z^6, \tag{21}$$

$$\zeta(z) = 1 + (2+k)z + 2z^2 + (2+k)z^3 - (2+k)z^4.$$
⁽²²⁾

The conjugacy map we consider is a Möbius transformation $\Gamma(z) = M(z) = \frac{z-a}{z-b}$, with the following properties:

$$(i)\Gamma(a) = 0, (ii)\Gamma(b) = \infty, (iii)\Gamma(\infty) = 1.$$

This enable us to discover that z = 0 and $z = \infty$ are clearly two fixed points of $\Omega = \Gamma \circ R_f \circ \Gamma^{-1}$, with $\Gamma(z) = \frac{z-a}{z-b}$, respectively, corresponding to fixed points *a* and *b* of R_f or the roots *a* and *b* of polynomial f(z) = (z - a)(z - b). Furthermore, $0, \infty$ are free of the parameter *k*. In addition, they are super-attractive points and suggest the critical points of $\Omega(z;k)$.

Nonetheless, because $\Omega(0) = 0$, $\Omega(\infty) = \infty$, that is, their orbits draw near themselves, such fixed points would show a tiny impact on the dynamics. Fixed points not excluding $\{0, \infty\}$ are termed as strange fixed points which are different from the roots of f(z). In order to find further strange fixed points, we need to solve the equations $\Omega(z;k) - z = 0$ in (20) for z with given values of k.

3.1. Strange Fixed Points

We first check the existence of *k*-values for common factors (divisors) of $\rho(z)$ and $\zeta(z)$. Additionally, both $\rho(z)$ and $\zeta(z)$ will be checked if they have a factor (z - 1). The following theorem best describes the relevant properties of such existence as well as explicit strange fixed points.

Theorem 3. (a) If k = -1, then $\rho(z)$ and $\zeta(z)$ have a common factor (z + 1), the operator Ω

has the strange fixed points
$$z = 1, z = -1, z = -i\sqrt{\frac{1}{2}(3\pm\sqrt{5})}, i\sqrt{\frac{1}{2}(3\pm\sqrt{5})};$$

- (b) If k = -3, then $\rho(z)$ and $\zeta(z)$ have two common factors $1 + z^2$ and $1 z + z^2$, the operator Ω has the strange fixed points $z = 1, z = -0.5 \pm 0.866025i$;
- (c) If $= -\frac{27}{7}$, then $\rho(z)$ has a factor (z 1), the operator Ω has the strange fixed points z = 1(triple), $z = -0.332631 \pm 1.13238i$, $z = -0.238798 \pm 0.812947i$;
- (d) If k = -5, then $\zeta(z)$ has a factor (z 1), the operator Ω has the strange fixed points $z = 0.357785, z = 2.79497, z = -0.367445 \pm 1.27423i, z = -0.208934 \pm 0.724541i;$
- (e) For k satisfying $(k+1)(k+3)(k+\frac{27}{7})(k+5) \neq 0$, the strange fixed points are given by z = 1 and the six roots of $\rho(z) = 0$ in terms of parameter k.

Proof. (a), (b) Suppose that $\rho(z) = 0$ and $\zeta(z) = 0$ for some values of $z \in \overline{C}$. Observe that parameter k exists in a linear way in all coefficients of both polynomials. By eliminating k from the two polynomials, we obtain the relation: $(1 + z^2)^2(1 + z)(1 - z + z^2) = 0$. Hence, (1+z), $1+z^2$, $(1+z^2)^2$, $(1-z+z^2) = 0$ are candidates for common divisors of $\rho(z)$ and $\zeta(z)$. First, substituting z = -1, we find $\rho(-1) = 6 - 3(3+k) = 0$, $\zeta(-1) = -3 - 3k = 0$, from which k = -1 is obtained. Indeed, $\rho(z)$ and $\zeta(z)$ reduce to, respectively, $(1+z)^2$ $(1+3z^2+z^4)$ and $(1+z)(1+2z^2-z^3)$, as desired. Second, dividing both $\rho(z)$ and $\zeta(z)$ by $1 + z^2$ yields the identical remainder as -k - 3, from which the remainder becomes zero for k = -3. Indeed, $\rho(z)$ and $\zeta(z)$ reduce to, respectively, $(1+z^2)(1-z+z^2)(1+z+z^2)$ and $(1+z^2)(1-z+z^2)$. Accordingly, we find that $(1-z+z^2)$ is indeed a common factor of $\rho(z)$ and $\zeta(z)$. Thus, dividing both $\rho(z)$ and $\zeta(z)$ by $1 - z + z^2$ yields the remainder, respectively, as -3k - 9(2k + 6) - (k + 3) and the remainder becomes zero for k = -3. After that, it is the same as the previous common factor $(1 + z^2)$. Third, the direct division of $\rho(z)$ by $(1+z^2)^2$ is performed, then its remainder as $(k+3)z^3 - (k+2)z^{2-k-2}$, which cannot be zero for any k-values. The rest of the proof only needs to substitute the corresponding $\rho(z)$ and $\zeta(z)$ into (20) and solve its roots to be the strange fixed points.

- (c) We find that $\rho(z)$ has a factor for $k = -\frac{27}{7}$ via solving $\rho(1) = 27 + 7k = 0$. In fact, for $k = -\frac{27}{7}$, we also obtain $\rho'(1) = 0$. Hence, $\rho(z)$ has a factor $(z 1)^2$ for $k = -\frac{27}{7}$. Then, $\rho(z) = \frac{1}{7}(z 1)^2(7 + 8z + 17z^2 + 8z^3 + 7z^4)$, $\zeta(z) = \frac{1}{7}(7 13z + 14z^2 13z^3 + 13z^4)$. The rest of its proof is trivial.
- (d) If (z-1) is a divisor of $\zeta(z)$, then $\zeta(1) = 5 + k = 0$. For a value of k = -5, we discover $\zeta(z) = (z-1)(-1-2z+3z^3)$. As a result, we obtain the strange fixed points *z* satisfying $\rho(z) = 1 2z 6z^3 2z^5 + z^6 = 0$. The rest of the proof is similar to (a).
- (e) It is straightforward for us to claim the results. We numerically obtain the strange fixed points *z* satisfying $\rho(z) = 0$ for given values of $k \notin \{-1, -3, -\frac{27}{7}, -5\}$; detailed analysis will be shown later.

3.2. Stability of the Strange Fixed Points

To study the stability of fixed points, we need to compute the first derivative of Ω from (19):

$$\Omega'(z;k) = -\frac{z^3(1+z^2)\chi(z)}{\zeta(z)^2}.$$
(23)

where:

$$\chi(z) = 8 + 4k + (2 + 7k + 3k^2)z - 4(7 + 4k + k^2)z^2 - 2(22 + 13k + k^2)z^3 - 4(7 + 4k + k^2)z^4(2 + 7k + 3k^2)z^5 + (8 + 4k)z^6.$$
(24)

$$\zeta(z) = 1 + (2+k)z + 2z^2 + (2+k)z^3 - (2+k)z^4.$$
⁽²⁵⁾

We first check the existence of *k*-values for common factors (divisors) of $\chi(z)$ and $\zeta(z)$. Additionally, $\chi(z)$ and $\zeta(z)$ will be checked for whether they have divisors z^3 , $1 + z^2$. The following theorem best describes the relevant properties of such existence as well as explicit strange fixed points.

Theorem 4. (a) If k = -5, then $\Omega'(z;k) = \frac{6z^3(1+z^2)(2-3z-3z^3+2z^4)}{(-1-2z+3z^3)^2}$; (b) If k = -1, then $\Omega'(z;k) = -\frac{2z^3(1+z^2)(2-5z-5z^3+2z^4)}{(-1-2z+3z^3)^2}$; (c) If k = -3, then $\Omega'(z;k) = 4z^3$;

- (d) If k = -2, then $\Omega'(z;k) = -\frac{12z^5(1+z^2)^2}{(1+2z^2)^2};$
- (e) If $(k+5)(k+1)(k+3)(k+2) \neq 0$, then we can demonstrate the desired stability of z by varying parameter k with the aid of Theorem 6.

Proof. (a), (b), (c) Suppose that $\chi(z) = 0$ and $\zeta(z) = 0$ for some values of z. By eliminating k from the two polynomials, we obtain the relation: $(z-1)(z+1)(1+z^2)(1-z+z^2)(1+z^2-z^2$ $4z^3 + 2z^4 - 4z^5$ = 0. Hence, $(z-1), (z+1), (1+z^2), (1-z+z^2)$ are candidates for common divisors of $\chi(z)$ and $\zeta(z)$. After checking constraints $\chi(1) = \zeta(1) = 0, \chi(-1) = \zeta(-1) = 0$, we find that k = -5, -1 yielding common divisors. By checking zero remainders when dividing both $\chi(z)$ and $\zeta(z)$ by $(1+z^2)$ and $(1-z+z^2)$, respectively, we find k = -3yielding common divisors.

- (d) For $\chi(z)$ to have a factor *z*, we require k = -2 by solving $\chi(0) = 8 + 4k = 0$ for *k*. For $\chi(z)$ to have a factor $(1 + z^2)$, we require its remainder $(8k^2 + 40k + 48)z = 0$, from which k = -3, -2.
- If $\zeta(z)$ has a factor $(1 + z^2)$, then its remainder is -k 3, from which the remainder (c) becomes zero for k = -3. The rest of the above proof only needs to substitute the corresponding $\chi(z)$ and $\zeta(z)$ into (23).
- If k satisfies $(k+5)(k+1)(k+3)(k+2) \neq 0$, then the graphical stability of z yields (e) the desired results.

Regarding the stability of the strange fixed points, we conduct research in two cases. The first is the stability of the strange fixed points for special *k*-values; the second is the stability of the other strange fixed points corresponding to the root of $\rho(z)$ and z = 1.

Using the results of Theorem 4 to describe the stability of the fixed points in Theorem 3 in terms of parameter k. Table 1 presents the stability results of the first case.

k	ξ				No. of ξ
			$ \Omega'(\xi;k) $: m^w		
-1	1	-1	$-i\sqrt{\frac{1}{2}(3\pm\sqrt{5})}$	$i\sqrt{\frac{1}{2}(3\pm\sqrt{5})}$	6
	6:g	3.5:g	7.81025:g	7.81025:g	
-5	0.357785	2.79497	$-0.367445 \pm 1.27423i$	$-0.208934 \pm 0.724541i$	6
	11.787:g	11.7871:g	5.73:g	5.72996:g	
-27/7	1((triple)	$-0.332631 \pm 1.13238i$	$-0.238798 \pm 0.812947 \mathrm{i}$		7
	1:f	4.91055:g	4.91066:g		
-3	1	$-0.5 \pm 0.866025i$			3
	4:g	4:g			
-2	$-0.732786 \pm 0.68046 \mathrm{i}$	$0.0732949 \pm 0.673932 \mathrm{i}$	$0.159491 \pm 1.46648 \mathrm{i}$		6
	4.86704:g	10.8788:g	10.8786:g		

Table 1. Stability of strange fixed points ξ for special *k*-values.

 ${}^{w}|\Omega'(\xi;k)|$: *m* denotes that ξ is attractive, parabolic and repulsive, if $m = e(|\Omega'| < 1), m = f(|\Omega'| = 1),$ $m = g(|\Omega'| > 1)$, respectively.

Theorem 5. Let $k \notin \{-5, -1, -3, -2, -\frac{27}{7}\}$, then the relevant properties of the strange fixed point z = 1 are as follows:

- If k = -4, then z = 1 is a superattracting point; *(a)*
- (b)
- (c)
- If $|k + \frac{251}{63}| < \frac{8}{63}$, then z = 1 is an attractive point; If $|k + \frac{251}{63}| = \frac{8}{63}$, then z = 1 is a parabolic fixed point; At last, if $|k + \frac{251}{63}| > \frac{8}{63}$, then z = 1 is a repulsive point. (*d*)

Proof. Substituting z = 1 into Equation (23), we obtain:

$$|\Omega'(1;k)| = 8 \left| \frac{4+k}{5+k} \right|$$

It is easy to confirm that $|\Omega'(1;-4)| = 0$. Again, $8|\frac{4+k}{5+k}| \le 1$ is equivalent to $8|4+k| \le 1$ |5 + k|. Let k = x + iy be an arbitrary complex number. Then:

$$|4+k|^2 = (4+x)^2 + y^2$$
 and $|5+k|^2 = (5+x)^2 + y^2$.

So:

$$64\Big[(4+x)^2+y^2\Big] \le (5+x)^2+y^2.$$

By simplifying:

$$63x^2 + 63y^2 + 502x + 999 \le 0.$$

That is:

$$\left(x+\frac{251}{63}\right)^2+y^2=\left(\frac{8}{63}\right).$$

Therefore:

 $|\Omega'(1;k)| \le 1$ if and only if $|k + \frac{251}{63}| \le \frac{8}{63}$.

 \square

Remark 2. For later use, we conveniently let $A = \{k \in C : |\Omega'(1;k)| < 1\}, S = \{k \in C : |\Omega'(1;k)| < 1\}$ $|\Omega'(1;k)| = 1$ and $S = \{k \in C : |\Omega'(1;k)| > 1\}$. Then, A, S, B indicate the sets where the strange fixed point z = 1 becomes attractive, repulsive, and indifferent, respectively. As seen in *Figure 1, S is the boundary between A as a circle of radius* $\frac{8}{63}$ *centered at* $\left(-\frac{251}{63},0\right)$ *. Furthermore, Figure 2 shows a stability surface with* z = 1 (*independent of parameter k*).



Figure 1. Stability circle *S* for the strange fixed point z = 1.



Figure 2. Stability surfaces of strange fixed point z = 1.

Lemma 1. If $\xi \in \overline{C}$ is the root of $\rho(z)$ defined in (21), then $\frac{1}{\xi}$ is also a root of $\rho(z)$ for any $k \in C$ and any $z \in \overline{C}$.

Proof. Let $\rho(z) = 0$. Substituting $\frac{1}{\overline{c}}$ into $\rho(z)$ for *z*, we obtain:

$$\rho\left(\frac{1}{\xi}\right) = 1 + (3+k)\frac{1}{\xi} + (5+k)\frac{1}{\xi^2} + 3(3+k)\frac{1}{\xi^3} + (5+k)\frac{1}{\xi^4} + (3+k)\frac{1}{\xi^5} + \frac{1}{\xi^6} \\
= \frac{\xi^6 + (3+k)\xi^5 + (5+k)\xi^4 + 3(3+k)\xi^3 + (5+k)\xi^2 + (3+k)\xi + 1}{\xi^6} \\
= \frac{\rho(\xi)}{\xi^6} \\
= 0.$$
(26)

This completes the proof. \Box

Corollary 1. If $\xi \in \overline{C}$ is a fixed point of Ω defined in (19), then $\frac{1}{\xi}$ is also a fixed point of Ω .

Proof. Based on the fact that Ω is conjugate to Ω via conjugacy $\Gamma(z) = \frac{1}{z}$. yThus, let $X = Y = \overline{C}, \mathcal{O} = \mathcal{F} = \Omega$ in Theorem 2, then Theorem 2 (b) completes this proof. \Box

Lemma 2. Let $(k+5)(k+1)(k+3)(k+2)(k+\frac{27}{7}) \neq 0$ and $\xi_1, \xi_2, \xi_3 \in C \setminus \{0\}$ be three roots of $\rho(z)$. Then, $\rho(z)$ can be expressed as the product of three second-degree polynomials as follows:

$$\rho(z) = \prod_{i=1}^{3} (1 + d_i z + z^2) = \prod_{i=1}^{3} (z - \xi_i) \left(z - \frac{1}{\xi_i} \right).$$

where $d_i = -\left(\xi_i + \frac{1}{\xi_i}\right)$ or $\xi_i = \frac{-d_i - \sqrt{d_i^2 - 4}}{2}$ for $1 \le i \le 3$ in terms of k.

Proof. In view of Lemma 1 and Corollary 1, we find that if $\xi \neq 0$ is a strange fixed point of $\Omega(z;k)$ found from the roots of $\rho(z) = 0$ for $k \notin \{-5, -1, -3, -2, -\frac{27}{7}\}$, then so is $\frac{1}{\xi}$.

Hence, the above factorization is valid. Then $\prod_{i=1}^{3} (z - \xi_i) \left(z - \frac{1}{\xi_i}\right)$ is easily acquired under the six-degree polynomials $\rho(z)$. $(1 + d_i(z) + z^2)$ can be obtained by expanding $\left(z - \frac{1}{\xi_i}\right)$, it is clear that $d_i = -\left(\xi_i + \frac{1}{\xi_i}\right)$ holds.

In order to continue to analyze the stability of the other fixed points from the root of $\rho(z)$, we need to find the root of $\rho(z)$ first by solving $\rho(z) = 1 + (3+k)z + (5+k)z^2 + 3(3+k)z^3 + (5+k)z^4 + (3+k)z^5 + z^6 = 0$. By comparing the expansion of $\prod_{i=1}^3 (1+d_iz+z^2)$ with the coefficients of the same order term of $\rho(z)$, we find that the following three equations are related to d_1, d_2, d_3 :

$$\begin{cases} d_1 + d_2 + d_3 = 3 + k, \\ 3 + d_1 d_2 + d_1 d_3 + d_2 d_3 = 5 + k, \\ 2d_1 + 2d_2 + 2d_3 + d_1 d_2 d_3 = 3(3 + k). \end{cases}$$
(27)

We eliminate variables d_2 , d_3 in (27) using the strong symbolic operation ability of Mathematica to ultimately access a decic equation in d_1 beneath:

$$d_1^{3} - (3+k)d_1^{2} + (2+k)d_1 - (3+k) = 0.$$

For convenience, we let $\Phi(d_1; k)$ denote $d_1^3 - (3+k)d_1^2 + (2+k)d_1 - (3+k)$. Eliminating d_1, d_3 and d_1, d_2 , respectively, we then obtain a unique decic equation in $c \in \{d_1, d_2, d_3\}$ as follows:

$$\Phi(c;k) = 0.$$

Now, we calculate the roots of $\Phi(c;k)$ and attain:

$$\begin{aligned} d_1(k) &= \frac{3+k}{3} - \frac{\sqrt[3]{2}(-3-3k-k^2)}{3\cdot n(k)} + \frac{n(k)}{3\cdot \sqrt[3]{2}}, \\ d_2(k) &= \frac{3+k}{3} - \frac{(1+i\sqrt{3})(-3-3k-k^2)}{3\cdot \sqrt[3]{4}\cdot n(k)} - \frac{(1-i\sqrt{3}\cdot n(k))}{6\cdot \sqrt[3]{2}}, \\ d_3(k) &= \frac{3+k}{3} - \frac{(1-i\sqrt{3})(-3-3k-k^2)}{3\cdot \sqrt[3]{4}\cdot n(k)} - \frac{(1+i\sqrt{3}\cdot n(k))}{6\cdot \sqrt[3]{2}}. \end{aligned}$$

where:

$$n(k) = \sqrt[3]{81 + 36k + 9k^2 + 2k^3 + 3\sqrt{3} \cdot \sqrt{239 + 204k + 86k^2 + 24k^3 + 3k^4}}.$$

Then, substituting $d_i(k)$ for $\xi_i = \frac{-d_i - \sqrt{d_i^2 - 4}}{2}$, i = 1, 2, 3. The result is a *k*-dependent strange fixed point of Ω . \Box

The stability surfaces of the strange fixed points dependent on parameter k are illustrated in Figure 3.



Figure 3. Stability surfaces of the strange fixed points ξ_i , i = 1, 2, 3.

Lemma 3. If $\xi \in \overline{C}$ is the root of $\chi(z)$ defined in (24), then $\frac{1}{\xi}$ is also a root of $\chi(z)$ for any $k \in C$ and any $z \in \overline{C}$.

Proof. As in the proof of Lemma 1, we obtain by direct calculation that $\chi(\frac{1}{\xi}) = \frac{\chi(\xi)}{\xi^6}$, because $\chi(\xi) = 0$, proof completed. \Box

Theorem 6. Let z be a strange fixed point of $\Omega(z; k)$ satisfying $\rho(z) = 0$. Then, the following holds:

$$\Omega'(z;k) = \Omega'\left(\frac{1}{z};k\right).$$
(28)

Proof. The derivative of Ω is $\Omega'(z;k) = -\frac{z^3(1+z^2)\chi(z)}{\zeta(z)^2}$. Substituting $\frac{1}{z}$ into $\Omega'(z;k)$ for z and knowing from the proof of Lemma 3 that $\chi\left(\frac{1}{z}\right) = \frac{\chi(z)}{z^6}$, eventually, we obtain $\Omega'\left(\frac{1}{z};k\right) = -\frac{(1+z^2)\chi(z)}{z^{11}\zeta\left(\frac{1}{z}\right)^2}$. As a result, it continues to check the relation:

$$\frac{z^3}{\zeta(z)^2} = \frac{1}{z^{11}\zeta\left(\frac{1}{z}\right)^2}$$

By direct computation:

$$\frac{z^3}{\zeta(z)^2} - \frac{1}{z^{11}\zeta(\frac{1}{z})^2} = \frac{z^{14}\zeta(\frac{1}{z})^2 - \zeta(z)^2}{z^{11}\zeta(z)^2\zeta(\frac{1}{z})^2} = \frac{-1 - (k+1)z + kz^2 + 2z^3 - 2z^5 - kz^6 + (1+k)z^7 + z^8}{z^{11}\zeta(z)^2\zeta(\frac{1}{z})^2} \cdot \rho(z) = 0.$$

is found. This completes the proof. \Box

The following lemma depicts how to find the super-attraction point under the strange fixed point of conjugated maps $\Omega(z;k)$ for some values of k. We need to concurrently solve $\rho(z) = \chi(z) = 0$ for some values of k; in this way, the super-attractors of $\Omega(z;k)$ can be found. We first eliminate the parameter k from $\rho(z) = \chi(z) = 0$, resulting in a polynomial over z with degree 16:

$$(1+z)^2(1+z^2)(1-z+z^2)(1+6z^2-9z^3+12z^4-16z^5+12z^6-9z^7+6z^8+z^{10}) = 0.$$

In view of Theorem 6, there are only eight pairs super-attractors denoted by $(z_i, \frac{1}{z_i})$ for $1 \le i \le 8$, the stability of z_i and $\frac{1}{z_i}$ are the same. Then, after solving $\Omega'(z;k) = 0$ for k numerically, we achieve the following Lemma deeming the needed super-attractors:

Lemma 4. Let $z_1 = -1, z_2 = i, z_3 = 0.5 + 0.866025i, z_4 = -0.791164 + 2.35788i, z_5 = -0.255105 - 0.966913i, z_6 = -0.127905 + 0.381192i, z_7 = 0.215007 + 0.976613i, z_8 = 0.959167 + 0.282841i$. Then $(z_i, \frac{1}{z_i}), (i = 1, 2, \dots, 8)$ are super-attractors of $\Omega(z; k)$, respectively, for eight values of $k \in \{0, -3, -3.33333 + 9.87028 \times 10^{-7}i, -0.870604 - 2.7565i, -1.2134 - 1.08719 \times 10^{-6}i, -0.87065 + 2.7565i, -2.78296 + 2.96051 \times 10^{-6}i, -3.77916 - 2.48912 \times 10^{-7}i\}$ or $\{-1, -2, -3 + 5.64784 \times 10^{-13}i, -2.23529 + 0.297382i, -3.50312 + 1.25267 \times 10^{-6}i, -2.23529 - 0.297381i, -2.71908 - 4.18506 \times 10^{-6}i, -4.79253 - 2.27167 \times 10^{-7}i\}$.

4. Critical Points and Free Critical Points

The critical points of map $\Omega(z;k)$ are defined as the roots of $\Omega'(z;k) = 0$. From this known fact, each basin of attraction contains at least one free critical point. Operator $R_f(z)$ has as critical point $0, \infty$ along with the roots of this sixth degree polynomial $\chi(z)$ where 0 and ∞ are related to *a* and *b*. Critical points that are different from 0 and ∞ are defined as free critical points. We are interested due to the orbital behavior of the free critical points, and we will deal with the free critical points which depend on parameter *k*.

Regarding the free critical points, we also divide the study into two parts such as the fixed points. One is to find the corresponding critical points for some special parameter valve k; the other is to find the critical points corresponding to the root of $\chi(z)$.

In view of Theorem 4, the free critical points for special *k*-values can easily be found and shown in Table 2.

k	η	No. of η
-5	$\pm i$, 0.557699, 1.79308, $-0.425391 \pm 0.90501i$	6
-1	$\pm \mathrm{i}, 0.360048, 2.77741, -0.318729 \pm 0.947846\mathrm{i}$	6
-3		0
-2	i (double), —i (double)	4
-27/7	$\pm {\rm i}, 0.326533 \pm 0.945186 {\rm i}, 0.67332 \pm 0.739352 {\rm i}, 0.974642 \pm 0.223771 {\rm i}$	8

Table 2. Free critical points ξ for special *k*-values.

Lemma 5. If $\eta \in \overline{C}$ is the root of $\chi(z)$ defined in (24), then $\frac{1}{\eta}$ is also a root of $\chi(z)$ for any $k \in C$ and any $z \in \overline{C}$, that is, $\chi\left(\frac{1}{z}\right) = \frac{\chi(z)}{z^6}$ hold.

Corollary 2. If η is any critical point, then so is $\frac{1}{\eta}$.

Proof. According to Theorem 2 (b), let $X = Y = \overline{C}$, $\mathcal{O} = \mathcal{F} = \Omega$ and $\Gamma(z) = \frac{1}{z}$ under the fact that Ω is conjugated to itself.

To seek *k*-dependent free critical points of $\Omega(z;k)$, we need to solve $\chi(z) = 0$ as defined in (24) for $k \notin \{-5, -1, -3, -2, -\frac{27}{7}\}$. Based on Lemma 5, $\chi(z)$ can be written as second-degree polynomials. As a result, $\chi(z) = (8 + 4k) \prod_{j=1}^{3} (1 + e_j z + z^2)$. We define a function $\psi(r) = \frac{-r - \sqrt{r^2 - 4}}{2}$ for the convenience of the following. Then, six roots η of $\chi(\eta) = 0$ are the *k*-dependent free critical points η_i of map $\Omega(z;k)$ given as follows:

$$\eta_j = \psi(e_j), \eta_{j+3} = \frac{1}{\eta_j}, (1 \le j \le 3)$$

where:

$$e_{1} = \frac{(2+7k+3k^{2})\delta_{1} - \delta_{0} + \delta_{1}^{2}}{12(2+k)\delta_{1}},$$

$$e_{2} = \frac{(4+14k+6k^{2})\delta_{1} + (1+i\sqrt{3})\delta_{0} - (1-i\sqrt{3})\delta_{1}^{2}}{24(2+k)\delta_{1}},$$

$$e_{3} = \frac{(4+14k+6k^{2})\delta_{1} + (1-i\sqrt{3})\delta_{0} - (1+i\sqrt{3})\delta_{1}^{2}}{24(2+k)\delta_{1}},$$

$$\delta_{0} = -1252 - 1324k - 493k^{2} - 90k^{3} - 9k^{4},$$

$$\delta_{1} = \sqrt[3]{-37720} - 65868k - 30990k^{2} + 739k^{3} + 2943k^{4} + 405k^{5} + 27k^{6} + 24\sqrt{3}\sqrt{\gamma_{0}},$$

 $\gamma_0 = -312336 - 727456k - 1288216k^2 - 2095336k^3 - 2200225k^4 - 1345816k^5 - 474904k^6 - 94582k^7 - 10240k^8 - 592k^9 - 9k^{10}.$

5. Parameter Spaces and Dynamical Planes

Lemma 6. $\Omega\left(\frac{1}{z};k\right) = \frac{1}{\Omega(z;k)}$ and $\Omega^q\left(\frac{1}{z};k\right) = \frac{1}{\Omega^q(z;k)}$ hold for any $k \in C, z \in \overline{C}$ and $q \in N$ are given.

Proof. The proof is simple. First, $\Omega(z;k) \circ \Gamma(z) = \Gamma(z) \circ \Omega(z;k)$, we chose a topological conjugacy $\Gamma(z) = \frac{1}{z}$, that is to say, $\Omega(z;k) \circ \frac{1}{z} = \Omega(\frac{1}{z};k) = \frac{1}{z} \circ \Omega(z;k) = \frac{1}{\Omega(z;k)}$. Then, $\Omega^q(z;k) = \Gamma(z) \circ \Omega^q(z;k) \circ \Gamma^{-1}(z)$ for a given $q \in N$. Similarly, $\Gamma(z) = \frac{1}{z}$ such that $\Omega^q(z;k) \circ \Gamma(z) = \Omega^q(\frac{1}{z};k) = \Gamma(z) \circ \Omega^q(z;k) = \frac{1}{\Omega^q(z;k)}$. \Box

Corollary 3. If $z \in \overline{C}$ is a q-periodic point of Ω , then so is $\frac{1}{z}$, let $q \in N$ be given.

Proof. With the help of Lemma 6, we have $\Omega^q(\frac{1}{z};k) = \frac{1}{\Omega^q(z;k)}$. As *z* is a *q*-periodic point of Ω , clearly, $\Omega^q(z;k) = z$. Hence, $\Omega^q(\frac{1}{z};k) = \frac{1}{z}$, this shows $\frac{1}{z}$ is also a *q*-periodic point. \Box

The above corollary explains whether the orbit of free critical point η_j , (j = 1, 2, 3) approaches a q-periodic point ξ_j of Ω , then the orbit of $\eta_{j+4} = \frac{1}{\xi_j}$ approaches a q-periodic point $\frac{1}{\xi_j}$ for any integer $q \ge 1$ and $k \notin \{-5, -1, -3, -2, -\frac{27}{7}\}$. This means that only one branch of η_j needs to be considered for its orbit behavior. We naturally define two concepts called the parameter space \mathcal{P} and dynamical plane \mathcal{D} for the iterative map $\Omega(z;k)$ as follows:

 $\mathcal{P} = \{k \in \overline{C}: \text{ an orbit of a free critical point } \eta \text{ tends towards a number } v \text{ under the action of } \Omega(z;k)\}.$

 $\mathcal{D} = \{z \in \overline{C}: \text{ an orbit of } z(k) \text{ for a given } k \in \mathcal{P} \text{ tends towards a number } u \text{ under the action of } \Omega(z;k)\}.$

We are going to seek the best members of the family by virtue of the parameter space connected with the free critical point η_i (see [6]) and dynamic plane.

Now, we analyze the asymptotic behavior of the free critical points of our proposed method family (2) with the technique [13]. These free critical points that depend on *k* are the roots of polynomial $\chi(z)$. For this, we draw the parameter space related to points η_j . In the previous analysis, we already know that the free critical points are conjugated, so only three planes need to be drawn. We create a mesh of 500×500 points in the complex plane and each point corresponds a different complex parameter value *k*. We select $z_0 = \eta_j(k)$ as an initial estimate for iterative schemes (2) where *k* goes through every point in the mesh and set the maximum number of iterations to 50. If the method converges to z = 1, the points are painted in blue; red denotes the convergence of the method to 0 (associated with *a*); yellow denotes those points that eventually converge to ∞ (associated with the *b*) and they are black in other cases, i.e., they diverge.

The parameter space related to the free critical points η_j are obtained in Figures 4–6 about different intervals. $k \in [a, b] \times [c, d]$ where [a, b] and [c, d] are the range of the real and imaginary parts of k, respectively; (b) is a detail on (a). In Figure 4b, we observe a small disk (the blue is denoted by D): D corresponds to values of k for which z = 1 is attractive or superattractive. In addition, it is very clear that, except for the red and yellow areas, the other color domains are not the best for the selection of the parameter k-value in terms of stability. From Figures 4–6, we see broad regions for red and yellow, suggesting that some members of family (2) are numerically stable. We will analyze the dynamic plane of family (2) with parameter k in the red and yellow areas of Figures 4–6.



Figure 4. Parameter space associated with the free critical point η_1 .



Figure 5. Parameter space associated with the free critical point η_2 .





The dynamic plane of the new family (2) for a given value of k is presented. Each basin of attraction is painted in a different color: convergence to 0 and ∞ are in orange and blue, respectively; if it converges to the fixed point z = 1 then it is in green; black indicates that the point does not converge to any root. We plotted the orbit of a point in red. Similarly to the parametric space, the maximum number of iterations is 50, and these dynamic planes are generated by a mesh of 500×500 points.

First of all, we provide dynamical planes for k = -1, k = -2, k = -3, k = -5 and k = -27/7 in Figure 7, respectively. It can be seen that only the black color appears in Figure 7e. In addition to Figure 7e, the attraction basins in other planes are only related to *a* or *b*, which means that *k*-values have good convergence properties. In other words, the corresponding members of the Ostrowski-type iterative family (2) are numerically stable.

Figure 8 presents the dynamic planes converging to the strange fixed point for k = -4, k = -3.75 and k = -0.870604 + 2.7565i. Furthermore, three different basins of attraction appear in each figure.

In Figure 9, different kinds of unstable behavior can be found. Figure 9 shows that z = 1 is a parabolic fixed point which is consistent with Theorem 5(c) and then it exists in the Julia set but it has its own basin of attraction (black in the figure).

Finally, in Figure 10a, the dynamical plane of the iterative scheme corresponds to k = -251/63 (the center of circle *S*) is shown. Figure 10b–d are the dynamic planes with respect to the neighborhood of -251/63 but still within the circle *S*. We find that with -0.01198 - 1.262i, -0.01896 - 1.283i, -0.01198 - 1.269i and -0.06088 - 1.283i as the

initial points, respectively, their orbits eventually converge to z = 1, which corresponds to Theorem 5 (b) and Remark 2.

Through the above analysis, we can see that these methods corresponding to k = -1, k = -2, k = -3 and k = -5 outperform the other elements of the Ostrowski-type iterative family (2) in most numerical applications. Furthermore, t = 1, k = -3 is the most stable member due to Cayley's test [17].



Figure 7. Dynamical planes for special *k*-values.



Figure 8. Dynamical planes converging to the strange fixed points.



Figure 9. Dynamical planes for |k + 251/63| = 8/63.



Figure 10. Dynamical planes of the neighborhood of -251/63.

6. Numerical Experiments

In this section, different test functions and previous methods are used to verify the effectiveness of the proposed method OM. The iterative methods used are as follows. Gupta et al. presented the method (GM, see [18]):

$$\begin{cases} y_n = x_n - \frac{f(x_n)}{f'(x_n)}, \\ w_n = y_n - \frac{f(y_n)}{f'(x_n)} \frac{f(x_n)}{f(x_n) - 2f(y_n)}, \\ x_{n+1} = w_n - \frac{f(w_n)(w_n - y_n)}{f(w_n) - f(y_n)}. \end{cases}$$
(29)

Maroju et al. showed the method (MM, see [11]):

$$\begin{cases} y_n = x_n - \frac{f(x_n)}{f'(x_n)}, \\ z_n = x_n - \left(1 + \frac{f(y_n)}{f(x_n)} \frac{f(x_n) + f(y_n)}{f(x_n) - f(y_n)}\right) \cdot \frac{f(x_n)}{f'(x_n)}, \\ x_{n+1} = z_n - \frac{f(z_n)}{f'(x_n) + g(x_n)(z_n - x_n)}. \end{cases}$$
(30)

where:

$$g(x_n) = \frac{2f(y_n)f'(x_n)^2}{f(x_n)^2}$$

Table 3 shows the test functions used in the subsequent comparison. As seen in Table 4, we clearly know that the error of k = -3 is the smallest under the same number of iterations, which is the same as the result in the dynamic plane. As Table 5 suggests, our proposed method shows favorable performance compared to Maroju and Gupta methods.

i	$f_i(x)$	<i>x</i> ₀	α
1	$\sin^2(x) - x^2 + 1$	1.3	1.4044916482153412
2	$x^3 - 10$	2	2.1544346900318837
3	$\cos\left(\frac{\Pi}{2}x\right) + x^2 - \Pi$	2	2.0347248962791266
4	$\exp^{\hat{x}}\sin x + \log(x^2 + 1)$	1	0
5	$x^4 - \log x - 5$	1.6	1.5259939537536892

Table 3. The function $f_i(x)$, initial guesses x_0 and zeros α .

Table 4. Comparison for special k-values.

k	$f_i(x)$	$ x_1 - x_0 $	$ x_2 - x_1 $	$ x_3 - x_2 $	$ x_4 - x_3 $
-5	$f_1(x)$	0.10449	1.7367×10^{-6}	2.8778×10^{-35}	$5.9578 imes 10^{-208}$
	$f_2(x)$	0.15443	$9.537 imes10^{-7}$	$4.6831 imes 10^{-38}$	$6.5657 imes 10^{-226}$
-1	$f_1(x)$	0.10449	$2.0417 imes10^{-6}$	$7.074 imes10^{-35}$	$1.2239 imes 10^{-205}$
	$f_2(x)$	0.15444	1.0177×10^{-6}	5.8526×10^{-38}	$2.1164 imes 10^{-225}$
-3	$f_1(x)$	0.10449	$1.4665 imes10^{-7}$	$3.5972 imes 10^{-43}$	$7.8344 imes 10^{-257}$
	$f_2(x)$	0.15443	$1.0231 imes10^{-7}$	$5.4912 imes 10^{-45}$	$1.3126 imes 10^{-268}$
-2	$f_1(x)$	0.10449	$8.5343 imes10^{-7}$	$1.8169 imes 10^{-37}$	$1.6918 imes 10^{-221}$
	$f_2(x)$	0.15444	$4.1677 imes 10^{-7}$	$1.2544 imes 10^{-40}$	$9.3276 imes 10^{-242}$
-27/7	$f_1(x)$	0.10449	$8.8583 imes 10^{-7}$	2.272×10^{-37}	$6.467 imes 10^{-221}$
	$f_2(x)$	0.15443	4.9404×10^{-7}	4.2762×10^{-40}	1.7983×10^{-238}

Table 5. Comparison of iterative methods for the test functions.

$f_{i}(x)$	n		ОМ		MM		OGM	
$J_i(x)$		$ x_n-x_{n-1} $	$ f(x_n) $	$ x_n-x_{n-1} $	$ f(x_n) $	$ x_n-x_{n-1} $	$ f(x_n) $	
$f_1(x)$	1	1.0449×10^{-1}	3.6406×10^{-7}	1.0449×10^{-1}	$5.16692 imes 10^{-7}$	1.0449×10^{-1}	$1.17938 imes 10^{-6}$	
	2	$1.4665 imes 10^{-7}$	$8.9299 imes 10^{-43}$	$2.0814 imes 10^{-7}$	$2.43255 imes 10^{-41}$	$4.7508 imes 10^{-7}$	$7.22718 imes 10^{-39}$	
	3	$3.5972 imes 10^{-43}$	$1.94486 imes 10^{-256}$	$9.7989 imes 10^{-42}$	$2.64877 imes 10^{-247}$	$2.9113 imes 10^{-39}$	3.82701×10^{-232}	
	4	7.8344×10^{-257}	$2.07559 imes 10^{-1538}$	$1.067 imes 10^{-247}$	$4.41503 imes 10^{-1483}$	1.5416×10^{-232}	$8.43734 imes 10^{-1392}$	
$f_2(x)$	1	$1.5443 imes 10^{-1}$	$1.42466 imes 10^{-6}$	$1.5443 imes 10^{-1}$	$3.95386 imes 10^{-6}$	$1.5443 imes10^{-1}$	$3.53539 imes 10^{-6}$	
	2	1.0231×10^{-7}	$7.6464 imes 10^{-44}$	$2.8394 imes 10^{-7}$	$1.39756 imes 10^{-40}$	$2.5389 imes 10^{-7}$	$5.35701 imes 10^{-41}$	
	3	$5.4912 imes 10^{-45}$	$1.82778 imes 10^{-267}$	$1.0037 imes 10^{-41}$	$2.72565 imes 10^{-247}$	$3.8471 imes 10^{-42}$	$6.48392 imes 10^{-250}$	
	4	1.3126×10^{-268}	$3.4098 imes 10^{-1609}$	$1.9574 imes 10^{-248}$	$1.49989 imes 10^{-1487}$	$4.6564 imes 10^{-251}$	$2.03858 imes 10^{-1503}$	
$f_3(x)$	1	3.4725×10^{-2}	$9.21225 imes 10^{-12}$	$3.4725 imes 10^{-2}$	$1.68613 imes 10^{-11}$	3.4725×10^{-2}	$3.70924 imes 10^{-10}$	
	2	2.2171×10^{-12}	$6.67277 imes 10^{-73}$	$4.058 imes 10^{-12}$	$7.38308 imes 10^{-71}$	$8.927 imes10^{-11}$	$9.67886 imes 10^{-62}$	
	3	1.6059×10^{-73}	$9.63714 imes 10^{-440}$	$1.7769 imes 10^{-71}$	$5.20373 imes 10^{-427}$	$2.3294 imes 10^{-62}$	$3.05532 imes 10^{-371}$	
	4	$2.3194 imes 10^{-440}$	$8.74583 imes 10^{-2641}$	1.2524×10^{-427}	$6.3794 imes 10^{-2564}$	7.3532×10^{-372}	$3.02312 \times 10^{-2228}$	
$f_4(x)$	1	0.95175	$5.29461 imes 10^{-2}$	0.95137	$5.33921 imes 10^{-2}$	0.95264	$5.18821 imes 10^{-2}$	
	2	$4.8254 imes 10^{-2}$	$1.67183 imes 10^{-8}$	$4.8627 imes 10^{-2}$	$4.89593 imes 10^{-8}$	$4.7363 imes 10^{-2}$	$2.02955 imes 10^{-7}$	
	3	$1.6718 imes 10^{-8}$	$5.33748 imes 10^{-47}$	$4.8959 imes10^{-8}$	$1.0712 imes 10^{-43}$	$2.0296 imes 10^{-7}$	$2.05003 imes 10^{-39}$	
	4	$5.3375 imes 10^{-47}$	$5.65194 imes 10^{-278}$	$1.0712 imes 10^{-43}$	1.17509×10^{-257}	$2.05 imes10^{-39}$	2.17731×10^{-231}	
$f_5(x)$	1	$0.74006 imes 10^{-1}$	$4.74263 imes 10^{-7}$	$0.74006 imes 10^{-1}$	$2.48963 imes 10^{-6}$	0.074006	$1.29563 imes 10^{-6}$	
	2	$3.4978 imes 10^{-8}$	$7.50338 imes 10^{-45}$	$1.8362 imes 10^{-7}$	$6.84712 imes 10^{-40}$	$9.5557 imes10^{-8}$	$7.70349 imes 10^{-42}$	
	3	$5.534 imes10^{-46}$	$1.17674 imes 10^{-271}$	$5.05 imes 10^{-41}$	$2.96314 imes 10^{-241}$	$5.6816 imes 10^{-43}$	$3.40343 imes 10^{-253}$	
	4	8.6788×10^{-273}	$1.7507 imes 10^{-1632}$	$2.1854 imes 10^{-242}$	$1.94632 imes 10^{-1449}$	$2.5101 imes 10^{-254}$	$2.53102 imes 10^{-1521}$	

7. Conclusions

We analyze the stability of the Ostrowski-type iterative method with six orders. The dynamic analysis of any quadratic polynomial was performed to select the family members with better stability characteristics. The stability of the strange fixed points is studied. There are many possible values for which members of this family have been proven to have good stability from the parameter space. The dynamic plane also shows that some members of this family are numerically stable. The results in Figure 7 and Table 4 are consistent, that is, t = 1, k = -3 has the best numerical behavior. Some numerical experiments support that the new method for t = 1, k = -3 has better performance than existing methods.

Author Contributions: Methodology, X.W.; writing—original draft preparation, X.C. All authors have read and agreed to the published version of the manuscript.

Funding: This research was supported by the National Natural Science Foundation of China (No. 61976027).

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: Not applicable.

Conflicts of Interest: The authors declare no conflict of interest.

References

- 1. Ortega, J.M.; Rheinbolt, W.C. Iterative Solution of Nonlinear Equations in Several Variables; Academic Press: New York, NY, USA, 1970.
- Wang, X. A family of Newton-type iterative methods using some special self-accelerating parameters. Int. J. Comput. Math. 2017, 95, 2112–2127. [CrossRef]
- 3. Wang, X. An Ostrowski-type method with memory using a novel self-accelerating parameter. *J. Comput. Appl. Math.* **2017**, *330*, 710–720. [CrossRef]
- Geaum, Y.H.; Kim, Y.I. Long-term orbit dynamics viewed through the yellow main component in the parameter space of a family of optimal fourth-order multiple-root finders. *Discret. Contin. Dyn. Syst. Ser. B* 2020, 25, 3087.
- 5. Ostrowski, A.M. Solution of Equations in Euclidean and Banach Space; Academic Press: New York, NY, USA, 1973.
- 6. Cordero, A.; Guasp, L.; Torregrosa, J.R. Choosing the most stable members of Kou's family of iterative methods. *J. Comput. Appl. Math.* **2018**, 330, 759–769. [CrossRef]
- Chun, C.; Ham, Y.M. Some sixth-order variants of Ostrowski root-finding methods. *Appl. Math. Comput.* 2007, 193, 389–394. [CrossRef]
- 8. Blanchard, P. Complex analytic dynamics on the Riemann sphere. Bull. Amer. Math. Soc. 1984, 11, 85–141. [CrossRef]
- 9. Carleson, L.; Gamelin, T.W. Complex Dynamics; Springer: New York, NY, USA, 1993.
- 10. Thompson, J.M.T.; Stewart, H.B. Nonlinear Dynamics and Chaos; John Wieley and Sons Ltd.: New York, NY, USA, 1986.
- 11. Maroju, P.; Magrenan, A.A.; Mosta, S.S.; Sarria, I. Second derivative free sixth order continuation method for solving nonlinear equations with applications, *J. Math. Chem.* **2018**, 56, 2099–2116. [CrossRef]
- 12. Magreñán, Á.A. Different anomalies in a Jarratt family of iterative root-finding methods. Appl. Math. Comput. 2014, 233, 29–38.
- Chicharro, F.I.; Cordero, A.; Torregrosa, J.R. Drawing dynamical and parameters planes of iterative families and methods. *Sci.* World J. 2013, 2013, 780153. [CrossRef] [PubMed]
- 14. Ahlfors, L.V. Complex Analysis; McGraw-Hill Book, Inc.: New York, NY, USA, 1979.
- 15. Beardon, A.F. Iteration of Rational Function; Springer: New York, NY, USA, 1991.
- 16. Wolfram, S. The Mathematica Book, 5th ed.; Wolfram Media: Champaign, IL, USA 2003.
- Babajee, D.K.R.; Cordero, A.; Torregrosa, J.R. Torregrosa, Study of multipoint iterative methods through the Cayley Quadratic Test. J. Comput. Appl. Math. 2014, 291, 358–369. [CrossRef]
- 18. Parhi, S.K.; Gupta, D.K. A sixth order method for nonlinear equations. Appl. Math. Comput. 2012, 218, 10548–10556. [CrossRef]