



## Article

# Growth of New, Optically Active, Semi-Organic Single Crystals Glycine-Copper Sulphate Doped by Silver Nanoparticles

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**Abstract:** The purpose of this study is to modify all physicochemical properties of glycine–copper sulphate single crystals, such as crystal habits, molar mass, thermal stability, optical activity, and electrical properties. The novelty of this study is growth of glycine–copper sulphate single crystals doped by a low concentration of silver nanoparticles (SNPs) that improved both crystal habits and physicochemical properties. The originality of this work is that trace amounts of SNPs largely increased the crystal size. Crystals have molar stoichiometric formula  $[\text{glycine}]_{0.95}, [\text{CuSO}_4 \cdot 5\text{H}_2\text{O}]_{0.05}$  in the absence and presence of silver nanoparticles (SNPs) in different concentrations: 10 ppm, 20 ppm, and 30 ppm. The crystals' names and abbreviations are: glycine–copper sulphate (GCS), glycine–copper sulphate doped by 10 ppm SNPs (GCSN1), glycine–copper sulphate doped by 20 ppm SNPs (GCSN2), and glycine–copper sulphate doped by 30 ppm SNPs (GCSN3). Dopant silver nanoparticles increased: crystallinity reflecting purity, transparency to UV-Vis. electromagnetic radiation, thermal stability, and melting point of glycine–copper sulphate single crystal. GCSN3 is a super conductor. High thermal conductivity of crystals ranging from  $1.1 \text{ W} \cdot \text{min}^{-1} \cdot \text{K}^{-1}$  to  $1.6 \text{ W} \cdot \text{min}^{-1} \cdot \text{K}^{-1}$  enabled attenuation of electromagnetic radiation and rapid heat dissipation due to good dielectric and polar properties. On rising temperature, AC electrical conductivity and dielectric properties of perfect crystal GCSN3 increased confirmed attenuation of thermal infrared radiation.

**Keywords:** glycine; copper sulphate; silver nanoparticle; single crystals; doping; optical activity



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## 1. Introduction

Good optical, dielectric, and thermal properties of semi-organic single crystals enable application in modern technologies for design components in photonic devices, optical communication systems, optoelectronics, frequency convertors, and nonlinear optical (NLO) devices [1]. Single crystals of glycine amino acid containing copper sulphate ( $\text{CuSO}_4$ ) are used in optical high-resolution band pass filters for spectral devices [2]. Such crystals having a good optical quality are rarely reported. At room temperature, glycine amino acid in zwitterion form is crystalized to  $\alpha$ ,  $\beta$ , and  $\gamma$  polymorphs [3–5]. Glycine–copper sulphate single crystals possess NLO activity and thermal stability due to synergism of both organic and inorganic components [6,7]. Glycine has a chiral center that crystallizes in non-centrosymmetric space groups [8]. Inorganic copper sulphate enhanced mechanical and thermal stability of glycine crystals [9]. Crystals' growth in the presence of doping impurities modified crystal habit and properties [10,11]. No studies are reported doping glycine–copper sulphate crystals by silver nanoparticles that is widely used in food, medical, industrial, catalysis, and pharmaceutical applications [12,13]. This study aims to grow new single crystals of glycine–copper sulphate in the absence and presence of silver nanoparticles to add new unique properties for these blue-colored glycine–copper sulphate

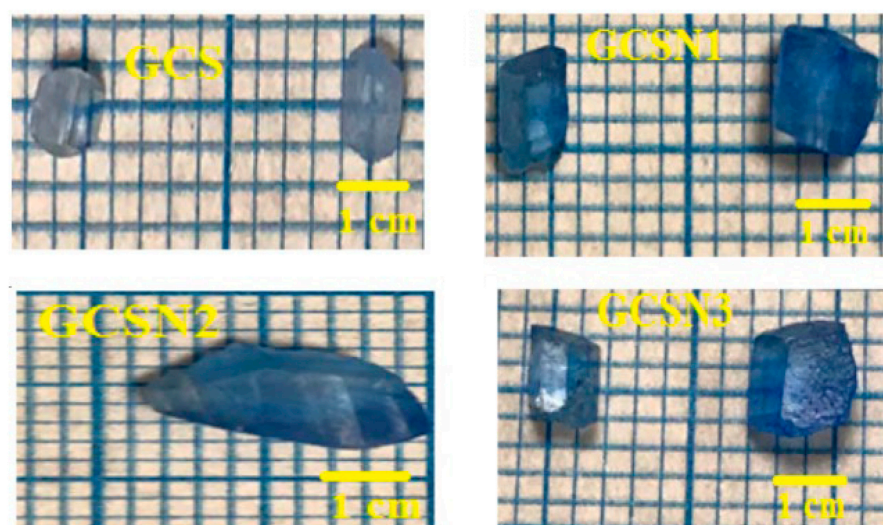
single crystals to improve absorption of weak infra-red photons reaching the earth from the sun causing global warming.

## 2. Experimental Procedure

### 2.1. Materials and Methods

All chemicals in this study are highly pure, of analytical grade, and used as received without further purification: glycine ( $\text{C}_2\text{H}_5\text{NO}_2$ , Oxford Co.) purity 98.5%,  $\text{CuSO}_4 \cdot 5\text{H}_2\text{O}$  (Sigma Aldrich Co., St. Louis, MO, USA), purity 98%. SNPs with polyvinylpyrrolidone, purity 99.9%, was purchased from Sigma Aldrich Co. with these physical characteristics: spherical shape nanoparticles: average diameter  $21.44 \pm 4.92$  nm and UV: electronic absorption bands at maximum wavelength ( $\lambda_{\text{max}}$ , 430 nm) due to delocalized electronic surface plasmon.

Slow solvent evaporation method is employed at  $25^\circ\text{C}$  for growth of GCS crystals doped by SNPs. The stoichiometric formula is  $(\text{glycine})_{0.95}$ ,  $[\text{CuSO}_4]_{0.05}$ . Salts are dissolved in double distilled water and agitated at 50 rpm using a magnetic stirrer for two h to obtain a homogeneous saturated solution. For SNP doping, a solution of glycine and  $\text{CuSO}_4$  are agitated at 50 rpm for two hours. Different concentrations of 10 ppm, 20 ppm, and 30 ppm SNPs are added to the filtrate that is further stirred for half an hour to complete homogeneity. The solution is covered by porous aluminum foil in a dust-free environment to allow slow solvent evaporation. High-quality, blue-color, pure crystals are harvested after one month, Figure 1.



**Figure 1.** Visual inspection of crystals.

GCS crystals increased in size and intensity of blue color as doping concentration of SNPs increased.

### 2.2. Characterization of Single Crystals

The grown crystals are characterized by:

Mass spectra (MS) by electron ionization technique at 70 eV using Thermo GCMS-ISQLT mass spectrometer; elemental analysis by energy dispersive X-ray analysis (EDX), and scanning electron microscope (SEM) using JSM-IT200 SEM.

Powder X-ray diffraction pXRD patterns at  $25^\circ\text{C}$  and diffraction, reflection angle ( $2\theta$ ) ranges from  $5^\circ$  to  $70^\circ$  at  $0.02^\circ$  step and scan rate  $1^\circ \text{ min}^{-1}$  using  $\text{Cu-K}\alpha$  radiation of wavelength  $1.5418 \text{ \AA}$  and acceleration voltage 40 kV using Bruker D8 advance diffractometer. Intensity of reflected X-rays in arbitrary units is plotted versus incidence and reflection angles  $2\theta^\circ$ , FTIR vibrational spectra at the frequency range  $400\text{--}4000 \text{ cm}^{-1}$  using IR Prestige-21, Borken, Germany.

UV-Vis. electronic absorption spectra using Helios alpha Unicom UV-Spectrophotometer at wavelength range 190–1200 nm; differential scanning calorimetry (DSC), thermal gravimetric analysis (TGA), and differential thermal gravimetric (DTG) analysis at temperature range: 25–800 °C using SDT Q600 V20.9 Build 20 instrument, 20 °C·min<sup>−1</sup> heating rate in de-aerated alumina cell to avoid sample oxidation by atmospheric oxygen; X-band electron spin resonance spectra ESR at room temperature, 9.43 GHz using reflection JES-RE1X ESR spectrometer in cylindrical resonance cavity with 100 kHz modulation, 5 mW power where applied magnetic field is controlled with LMR Gauss meter; electrical conductivity and dielectric characteristics of GCSN3 sample is measured using four probes Agilent 4294 A impedance bridge with sinusoidal voltage signal 10 mV amplitude. The sample is compressed as a pellet: 0.5 cm radius,  $0.23 \times 10^{-2}$  m thickness, and  $7.854 \times 10^{-5}$  m<sup>2</sup> geometrical area, coated on two opposite surfaces by silver paste for Ohmic contact with copper electrodes and annealed at 120 °C; thermal conductivity is measured at room temperature using hot disk TP 2500 [14].

### 3. Results and Discussion

MS is in Supplementary Information (SI); Figure S1 showed the relative abundance of the fragmented molecular ion versus  $\frac{\text{mass}(m)}{\text{charge}(z)}$  ratio peaked at  $m/z$  73.37, 75.2, 73.32, and 70.29, corresponding to molecular weight (M.wt) 75.07 g mol<sup>−1</sup> of glycine. The last peaks correspond to M.w. 614.32, 662.38, 758.57, 834.65 g mol<sup>−1</sup> for glycine–copper sulphate (GCS), glycine–copper sulphate doped 10 ppm SNPs (GCSN1), glycine–copper sulphate doped 20 ppm SNPs (GCSN2), and glycine–copper sulphate doped 30 ppm SNPs (GCSN3) crystals, respectively. An increasing concentration of SNPs increased chelation of glycine organic ligand to Cu(II) ion in CuSO<sub>4</sub>.

EDX spectra and SEM micrographs: Figure 2a–d show SEM-EDX analysis of GCS, GCSN1, GCSN2, and GCSN3 crystals, respectively. These spectra were produced as a focused electron beam on the sample ejected electrons from the inner-most electron atoms in the crystal leaving holes filled by ejected electrons from higher level emissions of X-ray [14,15].

EDX spectra confirmed that SNPs improved self-assembly of GCS from Figure 2a–d. Perfect crystallization is attained in Figure 2d.

The data in Table 1 indicated that oxygen, carbon, and nitrogen have maximum weight %, confirming that glycine is the base matrix of these single crystals. EDX spectrums of the doped crystals confirm the entry of both Cu(II) ion and Ag(I) ions into glycine crystal lattice.

The crystals have molar stoichiometric formula  $[\text{CH}_2\text{NH}_2\text{COOH}]_{0.95}$ ,  $[\text{CuSO}_4 \cdot 5\text{H}_2\text{O}]_{0.05}$  in the absence and presence of SNPs.

The vibrational band in FTIR spectra, as shown in Figure S2, is assigned to the function groups in the crystals, Table 2.

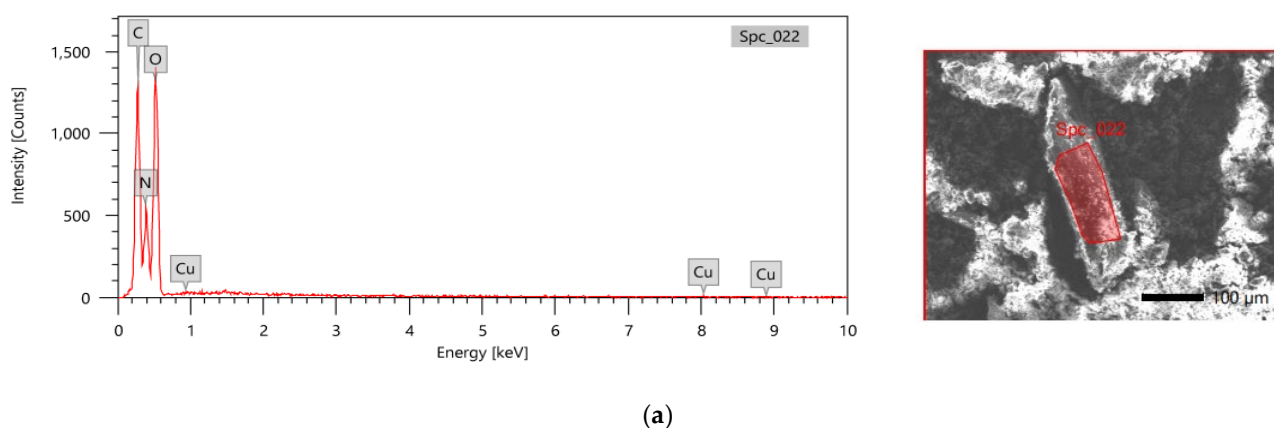
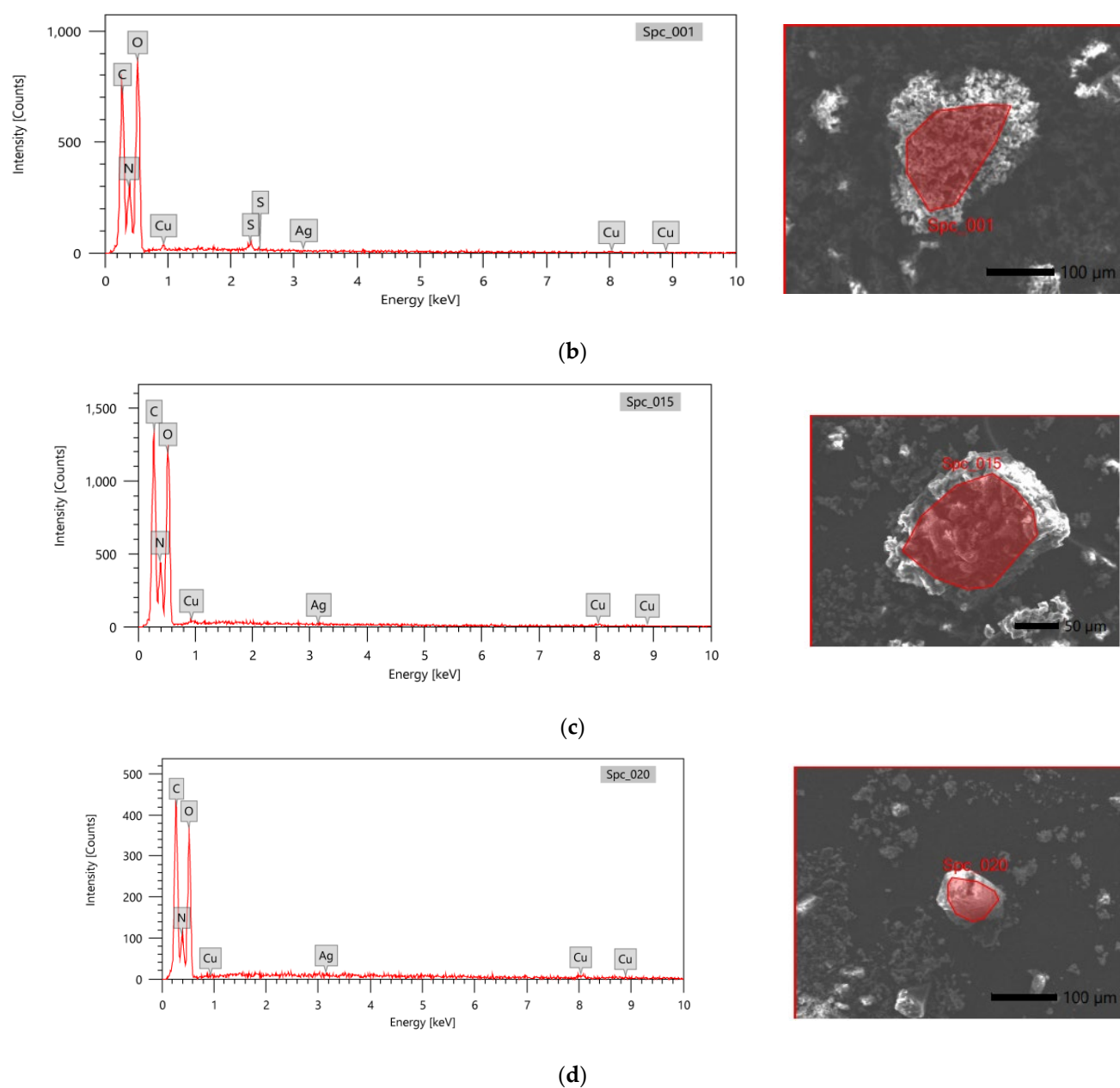


Figure 2. Cont.



**Figure 2.** (a–d): EDX spectra and SEM micrographs of GCS, GCSN1, GCSN2, and GCSN3, respectively.

**Table 1.** Weight percentage elements in crystals.

Crystal	Weight %					
	C	N	O	Cu	S	Ag
GCS	25.71	24.00	50.09	0.15	0.5	-
GCSN1	26.01	22.86	50.37	0.29	0.40	0.08
GCSN2	27.59	23.43	48.34	0.54	-	0.10
GCSN3	30.75	24.72	45.95	1.13	-	0.42

IR spectra of the GCS crystal showed a strong vibrational band at  $509.21\text{ cm}^{-1}$  due to Cu–N stretching [16], NH stretching band at  $3811.34\text{ cm}^{-1}$ , And medium peak at  $1111.00\text{ cm}^{-1}$  for  $\text{CH}_2$  rocking [17]. There is a strong peak at  $1334.79\text{ cm}^{-1}$  due to  $\text{CH}_2$  wagging. Intense peak C=O asymmetric stretching occurs at  $1604.77\text{ cm}^{-1}$  [18], symmetric stretching  $\text{COO}^-$  at  $1411.68\text{ cm}^{-1}$ , intense band asymmetric stretching  $\text{COO}^-$  at

1519.91  $\text{cm}^{-1}$  [19], and medium peak at 1033.85  $\text{cm}^{-1}$  for CCN asymmetric stretching deformation [20]. There is strong band  $\text{SO}_4^-$  stretching at 894.97  $\text{cm}^{-1}$ , an intense band due to bending  $\text{COO}^-$  at 694.37  $\text{cm}^{-1}$ , medium peak wagging  $\text{COO}^-$  at 609.51  $\text{cm}^{-1}$  [13], and  $\text{NH}_2$  asymmetric stretching at 2823.79  $\text{cm}^{-1}$  [21]. FTIR spectra of the samples GCSN1, GCSN2, and GCSN3 have small shift compared to that of GCS observed, which suggests the incorporation of SNPs into the crystals lattice.

**Table 2.** Assigned FTIR vibrational bands.

Vibration Mode	Wave Number ( $\text{cm}^{-1}$ )				
	Glycine [15]	GCS	GCSN1	GCSN2	GCSN3
Cu-N stretching	-	509.21	509.21	509.21	509.21
$\text{COO}^-$ wagging	606.30	609.51	609.51	609.51	609.51
$\text{COO}^-$ bending	696.9	694.37	694.37	694.37	694.37
CCN asymmetric stretching	1031.93	1033.85	1033.85	1033.85	1033.85
$\text{CH}_2$ rocking	1112.44	1111.00	1111.00	1126.43	1118.71
$\text{SO}_4^-$ stretching	-	894.97	894.97	894.97	902.51
$\text{COO}^-$ symmetric stretching	1492.39	1411.68	1411.89	1411.89	1404.18
$\text{CH}_2$ wagging	1311.89	1334.79	1334.74	1334.74	1319.31
$\text{COO}^-$ asymmetric stretching	1554.25	1519.91	1519.91	1504.48	1512.19
C=O asymmetric stretching	1643.53	1604.77	1604.77	1604.77	1604.77
$\text{NH}_2$ asymmetric stretching	2924.04	2823.79	2831.50	2831.50	2831.50
NH stretching	3921.39	3811.34	3996.21	3895.48	3903.92

Figure 3a–d showed indexed pXRD profile versus Rietveld refined pXRD patterns for GCS, GCSN1, GCSN2, and GCSN3 crystals. All pXRD patterns showed a prominent sharp diffraction peak at  $30^\circ$ . SNPs increased the peaks' intensity and modified crystal structure and lattice planes [22–24]. The crystals' structure and geometry agreed with Crystallography Opened Database, COD files. GCS and GCSN1 have monoclinic unit cell alpha glycine. Triclinic GCSN2 and GCSN3 have gamma glycine. pXRD patterns are refined using Full prof Suit software using CIF files containing crystal information. Peak patterns are refined following pseudo-Voigt profile analytical function [25]. Background and peak shapes are modeled with linear fitting by applying least-squares cycles and six background (polynomial 6th grade parameters) at the wavelength of Cu-detector and neglecting instrument contribution [26]. The crystallinity followed the trend:

$$\text{GCSN3} > \text{GCSN2} > \text{GCSN1} > \text{GCS}$$

Doping GCS with SNPs improved crystallinity, hence purity and crystal engineering.

During refinement, the number and order of crystalline planes and diffraction peaks increased in the same order. Many iteration cycles and all noise data are neglected; too long iteration time is consumed for GCSN2 and GCSN3 due to extra high crystallinity, long cartesian coordinates, and different angles in the triclinic unit cell. Intense peaks shifted to lower 2-theta, indicating a pillared crystal structure. Small peak absence and no polycrystallinity regions are observed in the perfect GCSN3 crystal, which confirmed good surface. There is good fitting of pXRD spectra (calculated intensity of the diffraction peaks are close to each other, resulting in a very negligible difference between the observed and calculated intensities ( $Y_{\text{obs}} - Y_{\text{cal}}$ )) with less than zero in arbitrary unit.

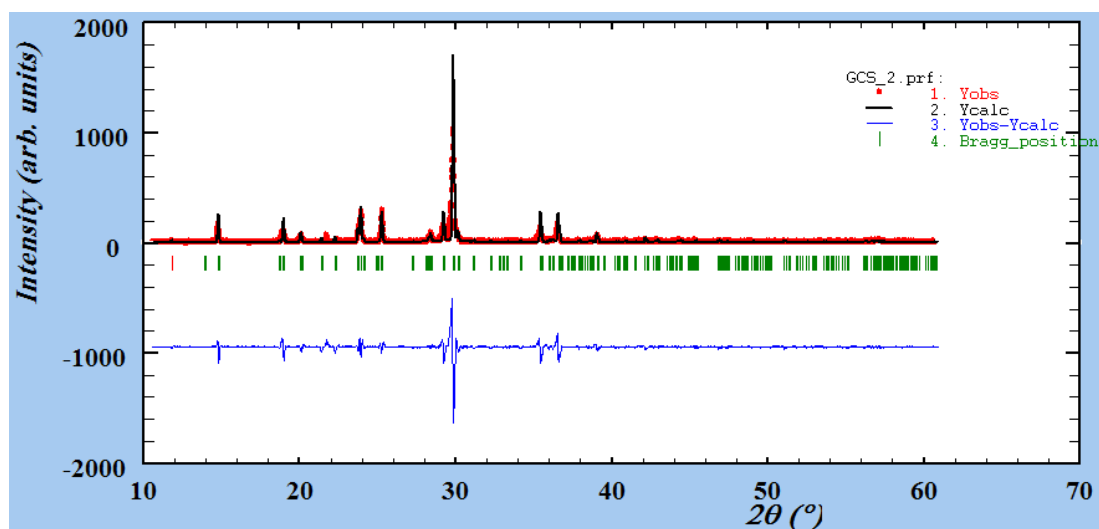
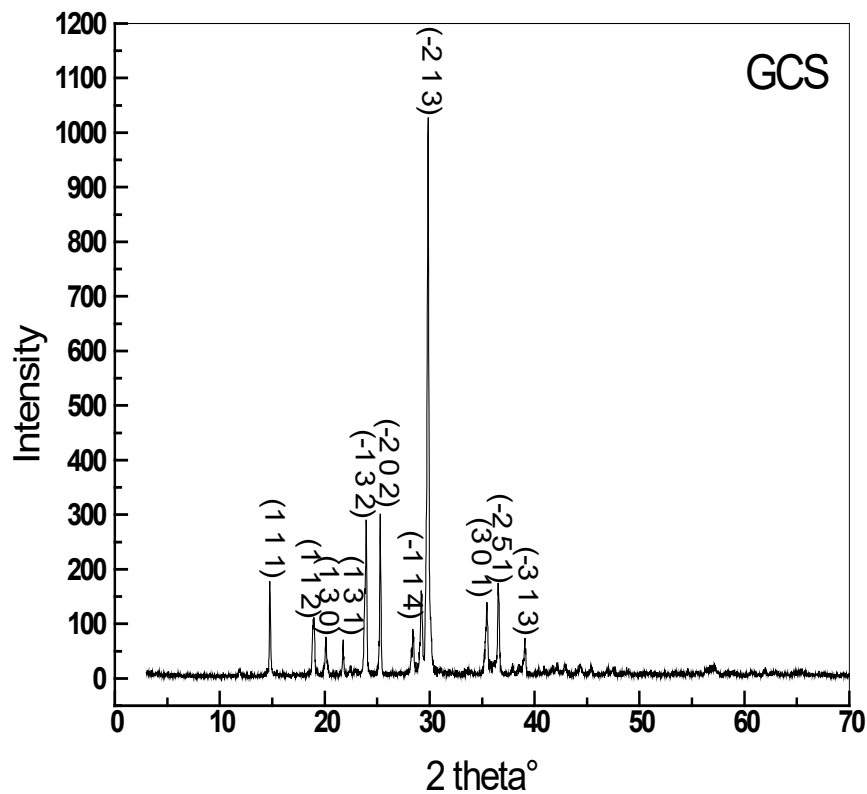
Both observed and calculated profiles are closely coincided to each other in a nonlinearly fit.

Sharp intense pXRD patterns with a dominating diffraction peak in crystals confirmed good crystallinity. Intensity changes and a slight shift in peak positions of GCS by SNPs



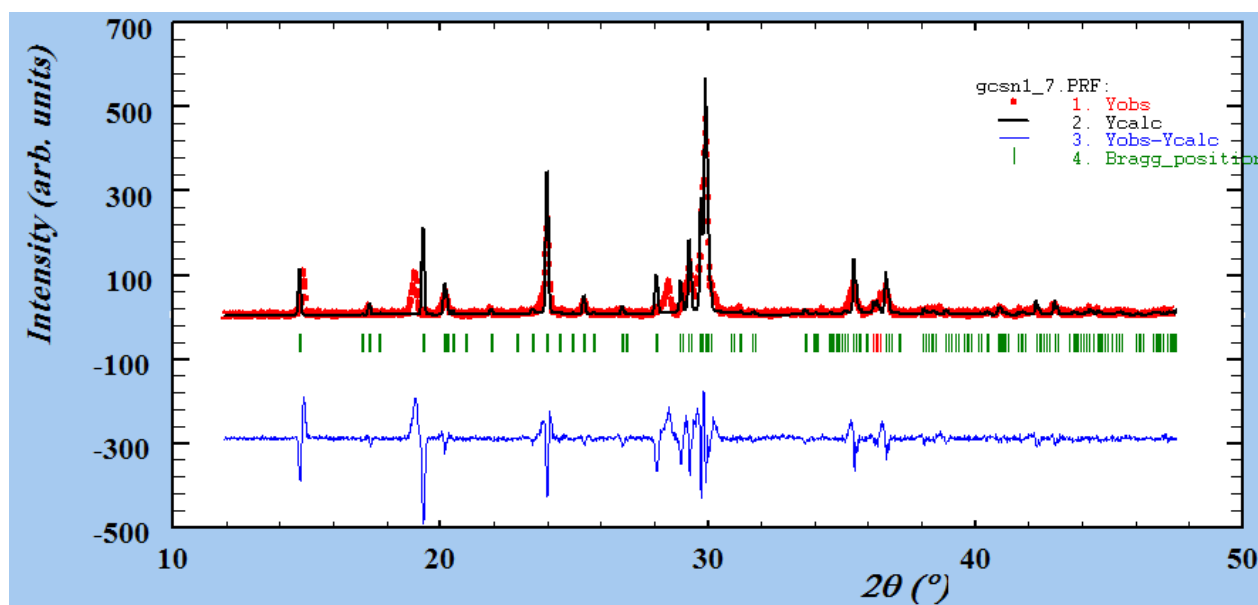
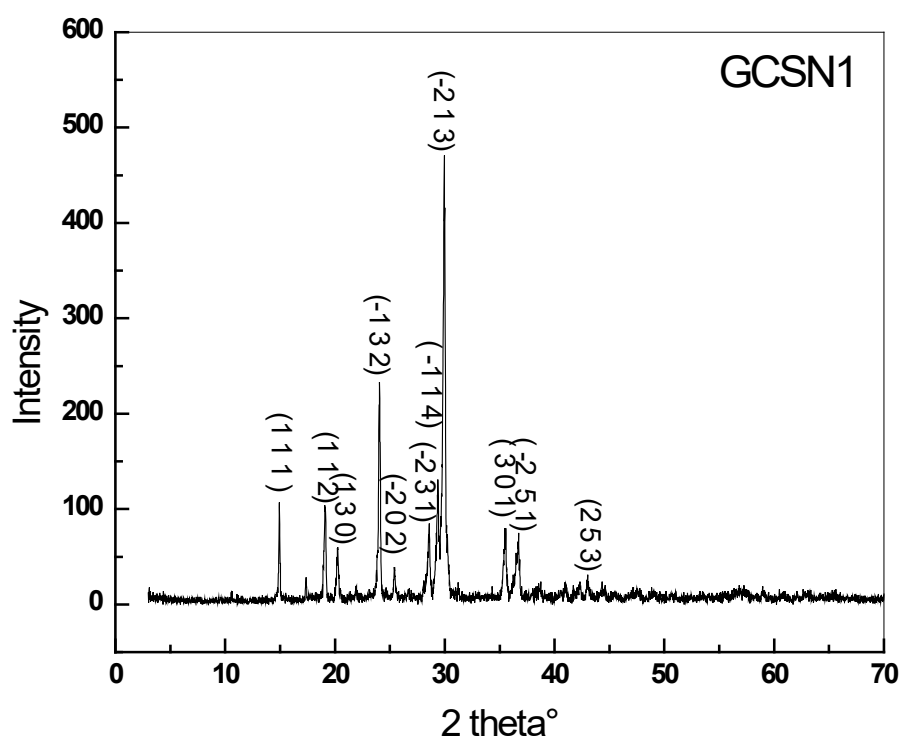
reflected modified crystalline planes [23]. Diffraction peaks at  $44.4^\circ$  and  $64.7^\circ$  in GCSN2 and GCSN3 confirmed doping impurities [24].

Miller indices  $h\ k\ \ell$  of the crystal planes, full width at half maximum (FWHM) of XRD patterns, peak position, and inter planar distance ( $d_{\text{cal.}}$ ), are collected in Table S1. The crystallinity and geometry of single crystals are deduced from pXRD by matching these diffraction patterns to pdf cards of similar crystals in Crystallography Opened Database using FWHM that characterized different material properties and surface integrity features [25,26]. The unit cell parameters are collected in Table 3.



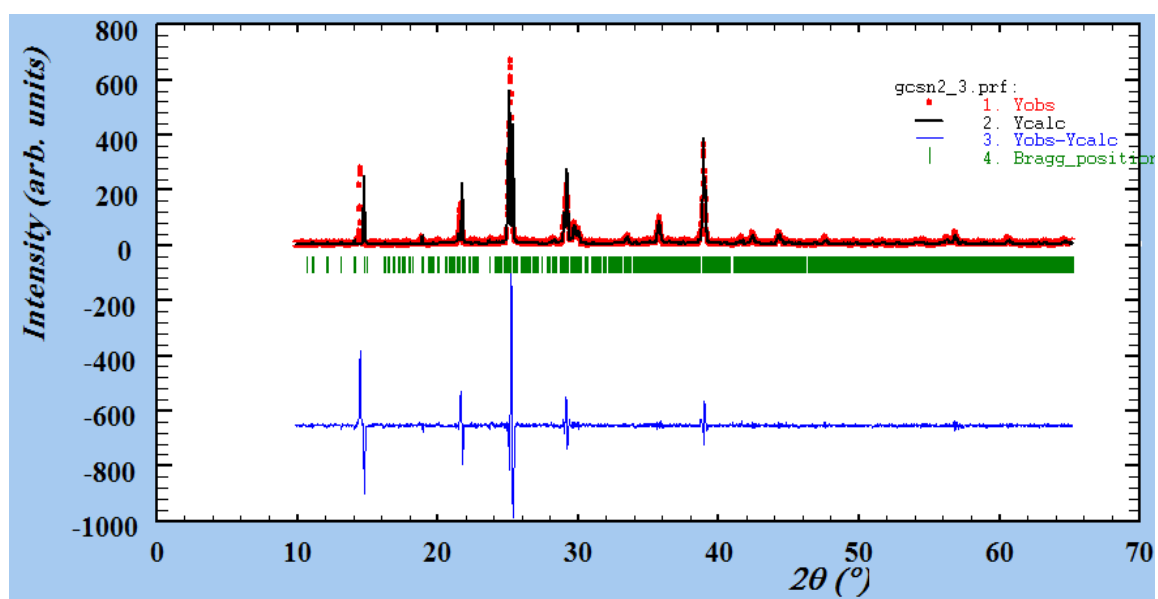
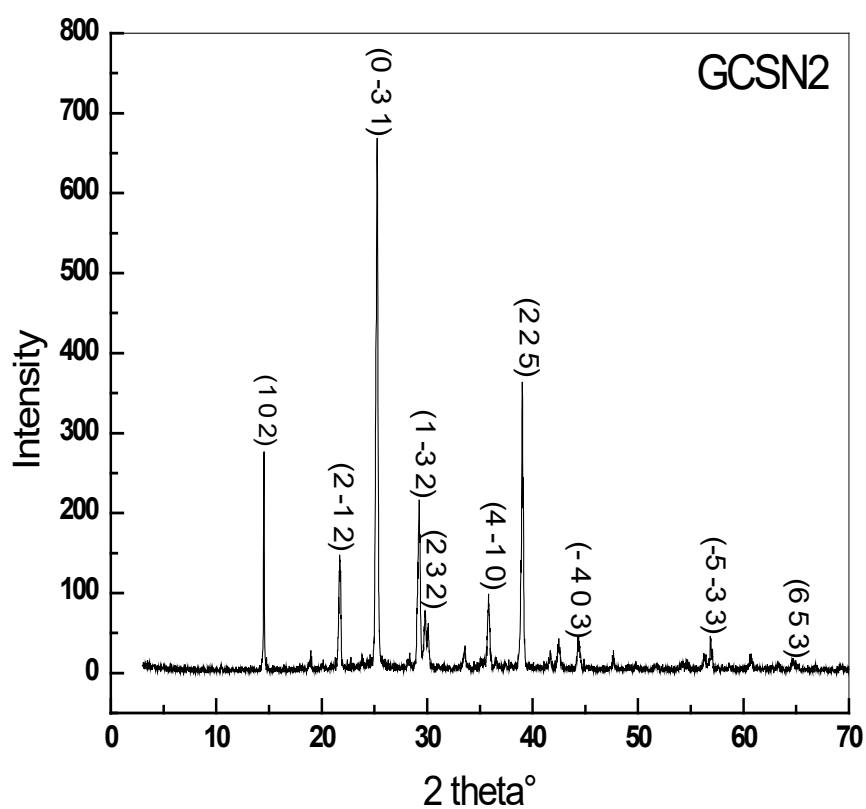
(a)

Figure 3. Cont.



(b)

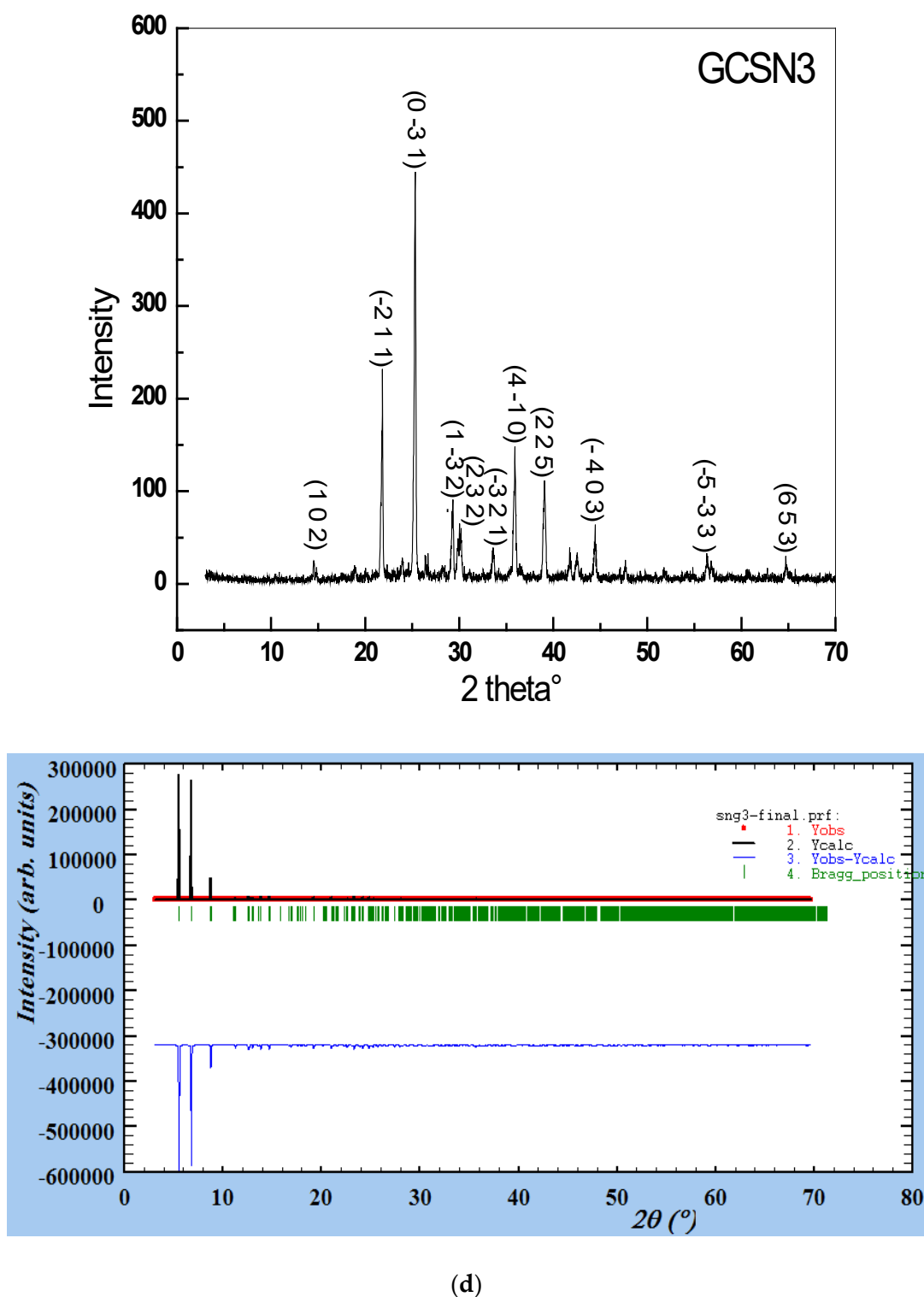
Figure 3. Cont.



(c)

Figure 3. Cont.



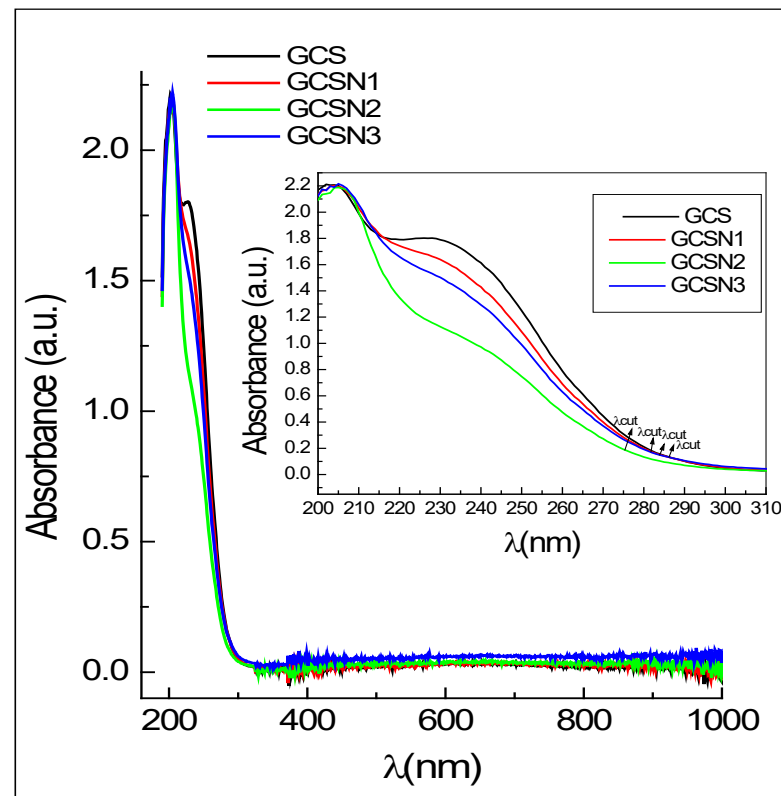


**Figure 3.** Observed and refined pXRD patterns of crystals respectively: (a): GCS, (b): GCSN1, (c): GCSN2 and (d): GCSN3.

UV-Vis. absorbance spectrum of crystals at wavelength range 190–1100 nm are shown in Figure 4. UV-Vis. absorbance curve showed cut-off of wavelength  $\lambda_{\text{cut off}}$  is lower than that of glycine [27].

**Table 3.** Crystal parameters.

Crystal	GCS	GCSN1	GCSN2	GCSN3
Identifier ID	1,505,763	1,505,763	4,342,413	4,342,413
Crystal System	Monoclinic		Triclinic	
Polymorph	$\alpha$		$\gamma$	
Space Group	P 2 <sub>1</sub>		P $\bar{1}$	
Lattice Parameter	a = 7.890 Å, b = 15.94 Å, c = 13.063 Å $\alpha = \gamma = 90^\circ$ , $\beta = 94.58^\circ$		a = 10.8222 Å, b = 11.0221 Å, c = 12.8735 Å $\alpha = 88.85^\circ$ , $\beta = 77.52^\circ$ , $\gamma = 84.55^\circ$	
Volume (Å <sup>3</sup> )	1638.14		1492.54	

**Figure 4.** UV-Vis. absorbance curve.

Doping by SNPs decreased  $\lambda_{\text{cut off}}$ , i.e., increased band gap for UV–electronic transition. SNPs enhanced transparency of crystals to UV radiation-enabled deposition as thin film on glass for protection against UV radiation [28,29]:

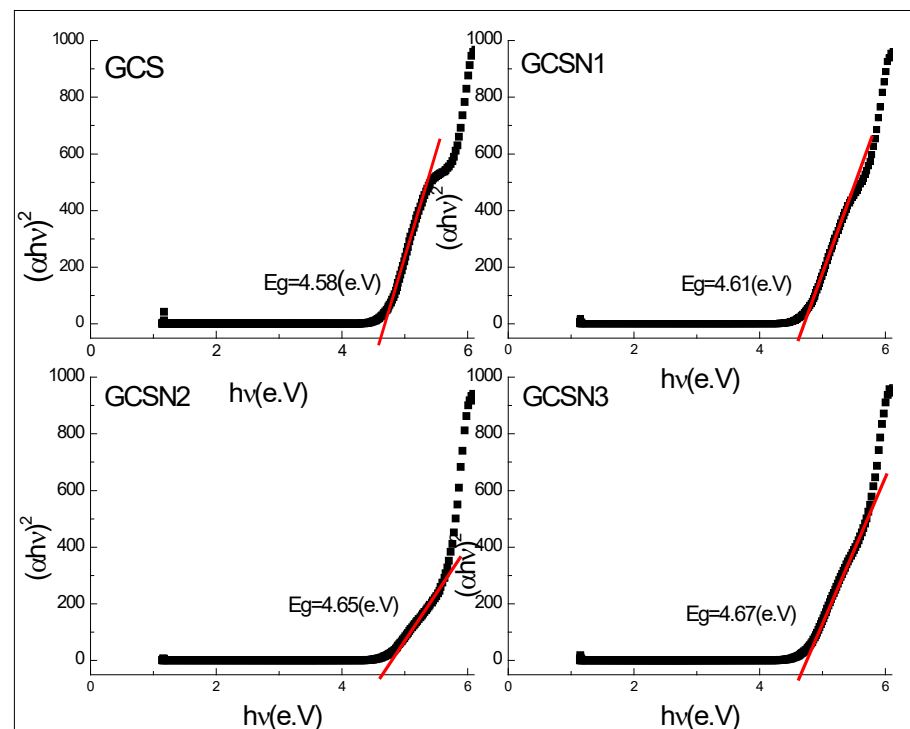
$$\text{Absorption fraction } \alpha = \frac{2.3042 \times \text{absorbance } A}{\text{sample thickness } t} \quad (1)$$

Band gap-controlled UV absorption coefficient depends on the energy of the incident photon and is estimated using Equation (2) [30,31].

$$\alpha h\nu = A(h\nu - E_g)^r \quad (2)$$

where  $\nu$  frequency of the incident radiation is inversely proportional to the wavelength of absorbance ( $\lambda$ ),  $A$  is constant, and the exponent  $r$  depends on the nature of electronic transition.  $r = 2$  for indirect transition, and  $r = \frac{1}{2}$  for allowed direct transition,  $r = \frac{1}{2}$ . Since all crystals are blue colored, allowed direct transition is considered [15,30]. The optical gaps,  $E_g$ , are calculated from plots  $(\alpha h\nu)^2$  as a function of photon energy ( $h\nu$ ), as shown

in Figure 5, and the observed values  $E_g$  are given in Table 4 along with the values of all reported glycine single crystals.



**Figure 5.** Tauc's plots for calculation of band gap.

**Table 4.** Comparison between reported and present study: optical band gap ( $E_g$ ) and  $\lambda_{\text{cut off}}$ .

Single Crystal	Band Gap		$\lambda_{\text{cut off}}$ (nm)	
	Value	Ref.	Value	Ref.
Glycine	3.13	[15]	346	[15]
GCS	4.58	Present study	287	Present study
GCSN1	4.61		283	
GCSN2	4.65		276	
GCSN3	4.67		280	
GCS	3.8	[23]	250	[23]
GCN	5.77	[15]	231	[15]
GLS	5.428	[32]	330	[32]
GZS	6.2	[33]	200	[33]

$\lambda_{\text{cut off}}$  287 nm for GCS decreased to 283, 276, and 280 nm for GCSN1, GCSN2, and GCSN3, respectively, which indicated an increasing band gap. Blue shift of  $\lambda_{\text{cut off}}$  to lower values indicated SNPs' improved polarizability of electron density on the single crystals. This finding suggested the suitability of GCS-doped SNPs single crystals for applications in optoelectronic devices such as frequency multiplier, sum-, difference-, and blue laser frequency generators, etc. [34]. Optical band gaps' eV are 4.58, 4.61, 4.65, and 4.67 for GCS, GCSN1, GCSN2, and GCSN3 crystals, respectively. SNPs increased  $E_g$  and enhanced optical properties of GCS. High  $E_g$  indicated the decrease in the localized energy states on doping by SNPs due to extrinsic defects or disorders in GCS caused by interstitial doped SNPs [35].

Good optical properties of the GCS crystals doped by SNPs are confirmed by extinction coefficient  $K$  calculated using Equation (3);  $\alpha$  and refractive index ( $n$ ) reflect dissipated incident radiation by absorption and scattering [36,37].

$$K = \frac{\lambda\alpha}{4\pi} \quad (3)$$

Reflectance  $R$  and  $K$  depend on photon energy, Figure S3: The reflectance is calculated using Equation (4): [36]

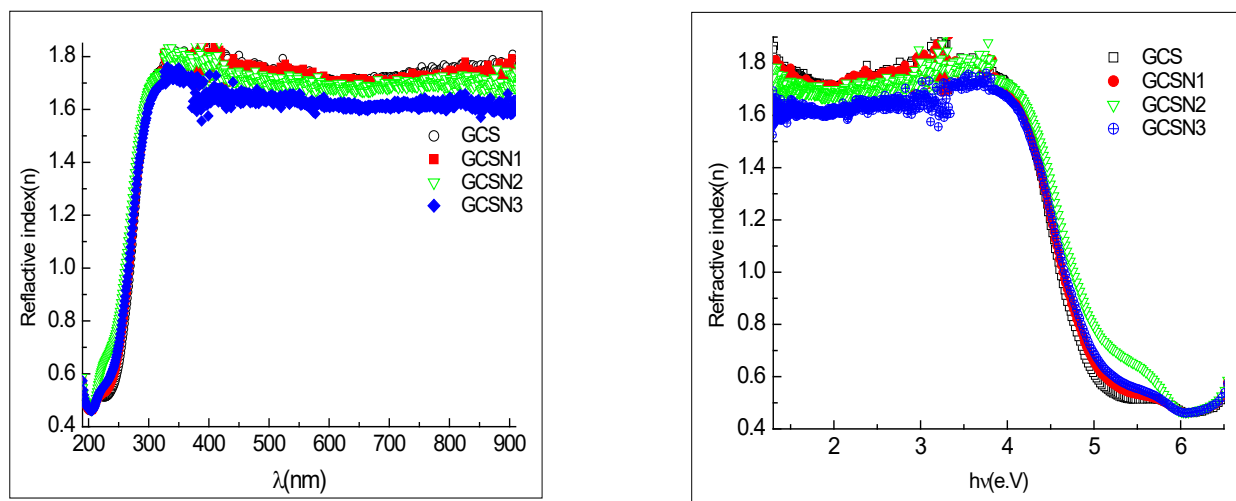
$$R = 1 \pm \frac{\sqrt{1 - \exp(-\alpha t) + \exp(\alpha t)}}{1 + \exp(-\alpha t)} \quad (4)$$

where  $t$  is the sample thickness.

Refractive index ( $n$ ) depends on wavelength, and photon energy is calculated using mathematical Equation (5): [36–38], Figure 6.

$$n = -\frac{(R + 1) \pm \sqrt{3R^2 + 10R - 3}}{2(R - 1)} \quad (5)$$

Positive refractive index ( $n$ ) 1.6–1.8 indicated dispersion of incident radiation on the crystals is inversely proportional to the photon energy. Refractive index ( $n$ ) is decreased by increasing photon energy and concentration of SNPs. High photon energy enables passing through the crystal lattice with low dispersion. SNPs decreased dispersion of incident radiation by improving transparency to UV radiation. High transmission and low absorbance of UV radiation and low refractive index suggest the single crystals are suitable for antireflection coating in solar thermal devices and NLO applications [38].



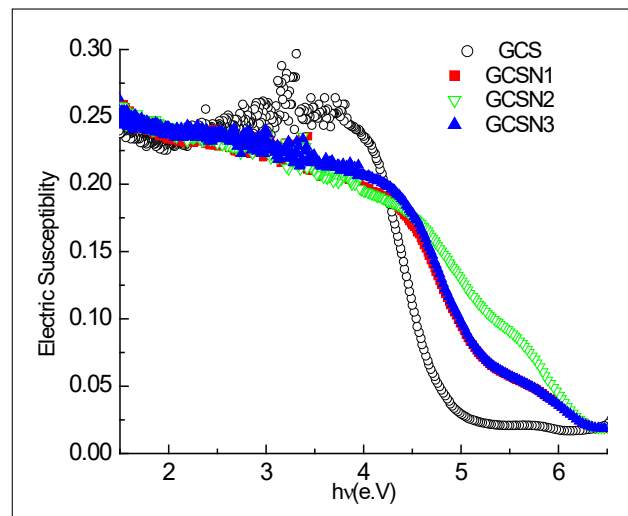
**Figure 6.** Refractive index with wavelength and photon energy.

Figure 7 shows the nonlinear variation of the electrical susceptibility ( $\chi$ ) with the photon energy.

GSC crystals showed large noise scattering in electrical susceptibility. This scattering disappeared on doping by SNPs that elevated electrical susceptibility ( $\chi$ ) at high photon energy following the order:

$$\text{GCSN2} > \text{GCSN3} > \text{GCSN1} >>> \text{GCS}$$

GCSN3 showed lower  $\chi$  than GCSN2 due to extra high crystallinity; GCSN3 decreased the mean free path for electron charge transfer.



**Figure 7.**  $\chi$ -photon energy plot.

The calculated electric susceptibility is plotted with photon energy; Figure 7 shows electric susceptibility ( $\chi$ ) is about 0.25 for all crystals at low photon energy.

Melting point, crystallization, phase transition, enthalpy, and entropy changes are clarified using DSC thermograms [39], Figure 8a–d. A sharp endothermic peak at decomposition temperature reflects good crystallinity. A weak broad peak at 195.80 and 206.29 °C in GCSN2 and GCSN3 characterized phase change from  $\gamma$ -glycine to  $\alpha$ -glycine [40]. Melting point (m.p.) is 253.83, 262.37, 265.64, 268.25 °C for GCS, GCSN1, GCSN2, and GCSN3, respectively. Increased m.p. on doping by SNPs is attributed to high thermal stability.

Heat capacity “Cp” is the ratio between heat flow to heating rate and depends on temperature: [41–44].

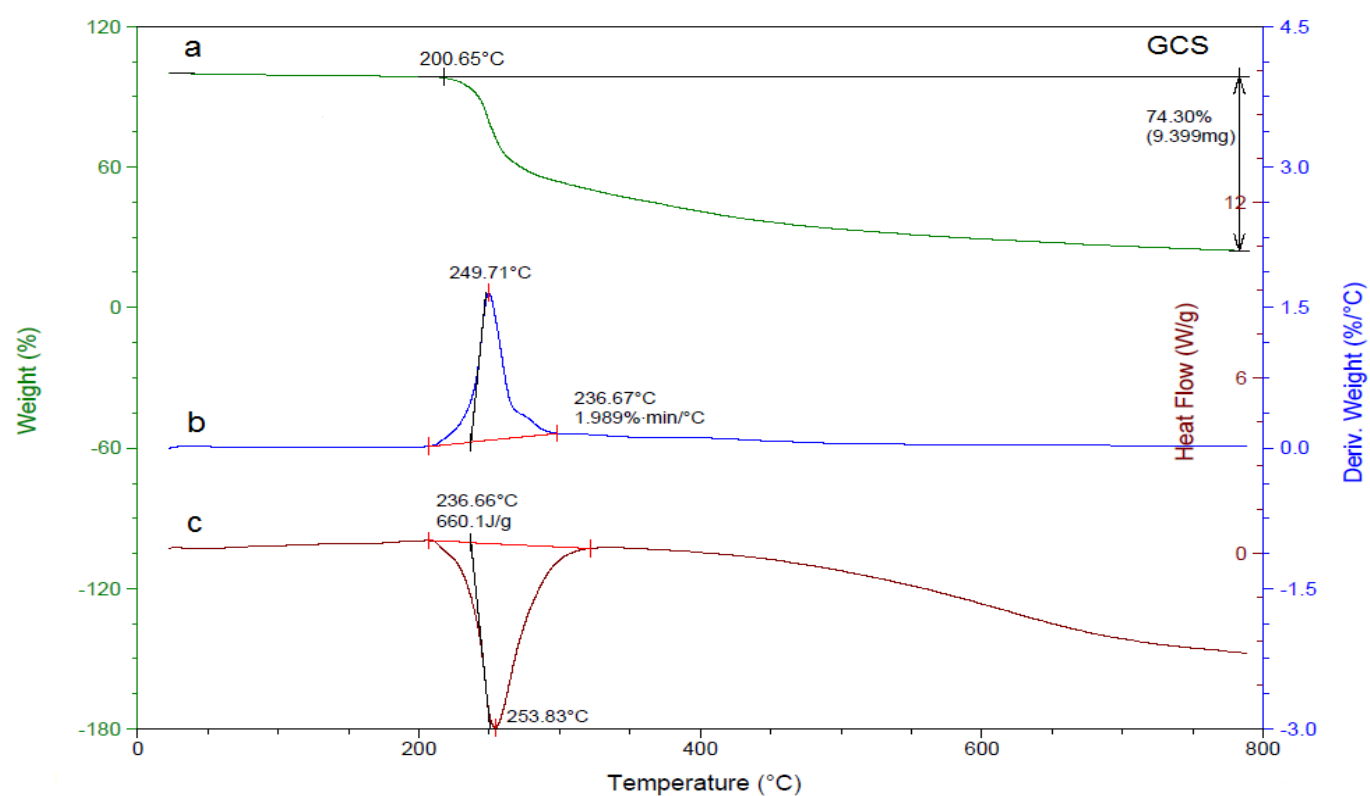
$$C_p = aT + b, \frac{C_p}{T} = \alpha T^2 + \gamma \quad (6)$$

where parameters a and b are calculated from intercept (a) and slope b of straight-line Cp-T plot. Cp of all prepared single crystals showed nonlinear variation with the absolute temperature, Figures S4–S7.

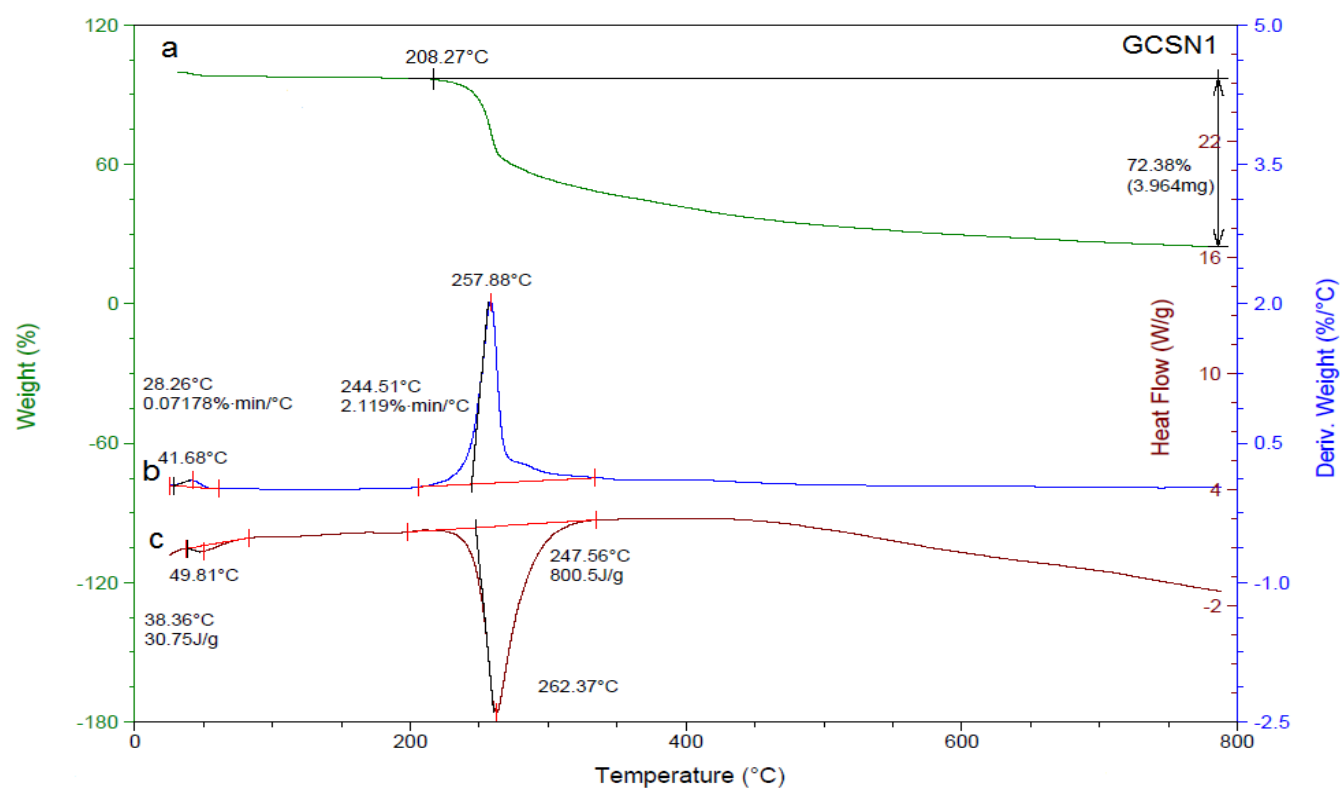
Thermal lattice coefficient  $\alpha$  and electronic heat capacities  $\gamma$  are obtained from linear  $\frac{C_p}{T}$  versus  $T^2$  plot, (coefficient,  $R^2$  above 0.99) Figures S4a,b and S7a,b. Table 5 shows linear fits parameters of DCS.

Heat capacity at constant pressure, Cp, is the heat required to raise the temperature of the crystal sample by 1 °C and represents the variation of the heat content of the crystal sample on heating. The variation of  $\alpha$ ,  $\gamma$  coefficients approved Cp variation with temperature. This finding indicated that the Cp amount of thermal heat absorbed by the crystals increases on heating, enabling the application of a heat shielding coating on thin film glasses.

Thermograms are shown in Figure 8a–d. DTG showed one peak confirmed one thermal decomposition step. GCS showed weight loss (wt. loss) 74.3 wt. % at temperature range 200.65–787.67 °C, DTG peak at 249.7 °C. Glycine decomposed into CO<sub>2</sub>, NH<sub>3</sub> [45] leaving 25.7% residue, 72.31%. GCSN1 decomposed at 208.27–788.29 °C and 27.69% residue. DTG peak was at 257.88 °C. In addition, 72.93 wt. % of GCSN2 decomposed at 217.41–788.84 °C, 27.07 wt. % residue, DTG peaked at 262.28 °C. Additionally, 72.86 wt. % of GCSN3 decomposed at 221.54–787.96 °C, gave 27.14 wt. % residue, and peak at 264.07 °C. SNPs decreased wt. loss on thermal decomposition of the crystals. Higher residue left after thermal decomposition confirmed improved crystallinity on doping GCS by SNPs.



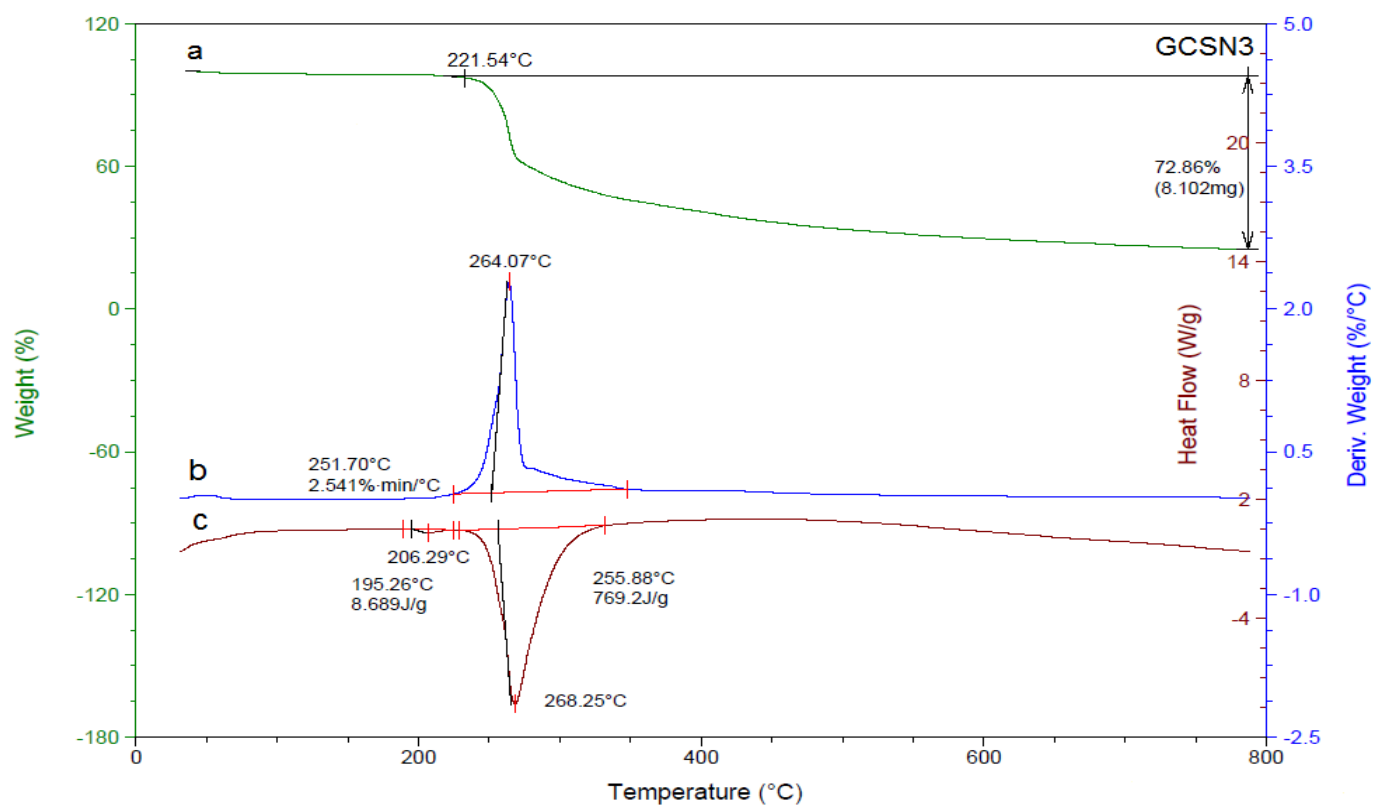
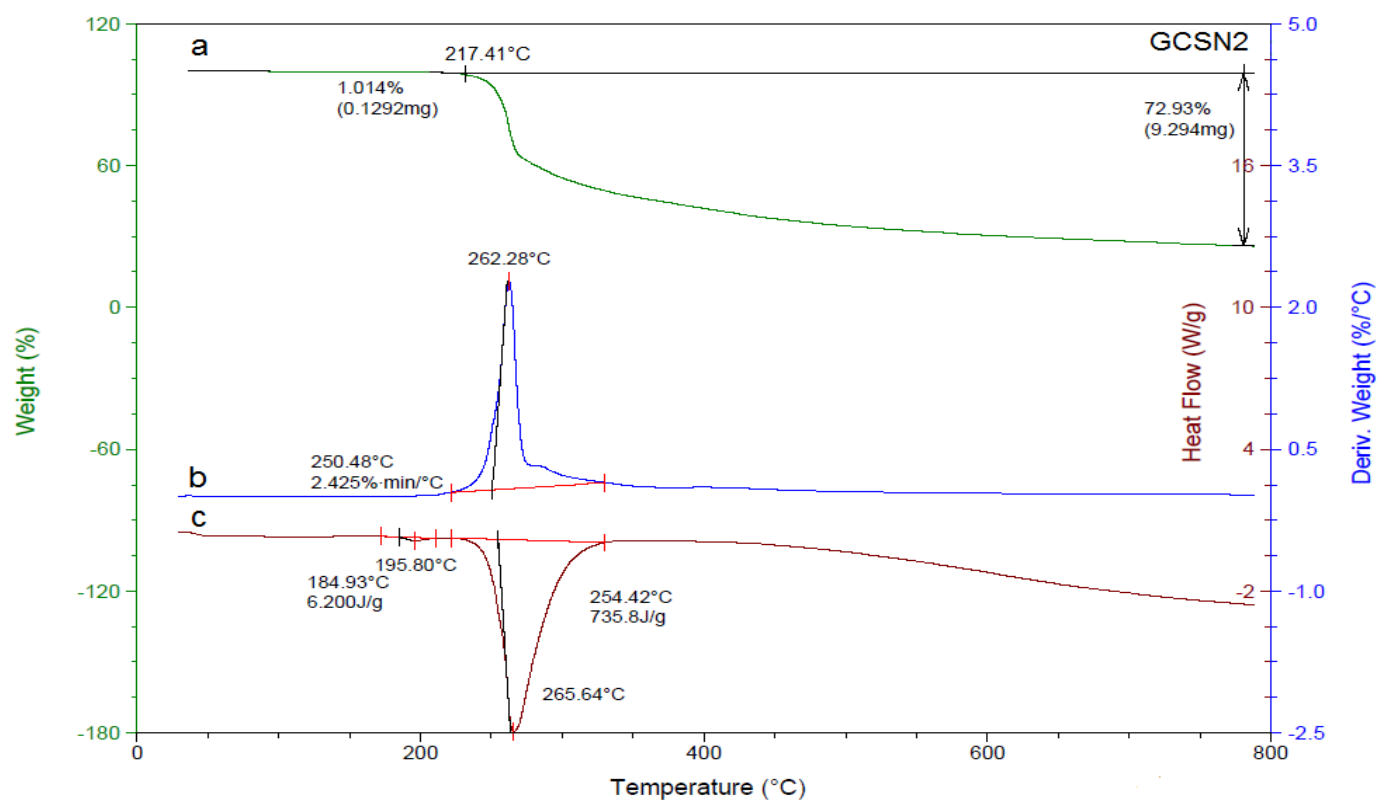
(a)



(b)

Figure 8. Cont.





**Figure 8.** (a–d): Thermograms: a. TGA, b. DTG, and c. DSC of GCS, GCSN1, GCSN2, and GCSN3.

**Table 5.** Linear fits parameters of DSC curves.

Crystal	Temperature Range (°C)	$C_p = a T + b$		$C_p/T = \alpha T^2 + \gamma$	
		a	b	$\alpha \times 10^{-6}$	$\gamma$
GCS	22–206	0.003	−0.441	0.00707	0.001
	206–254	−0.379	185.6	−0.7	0.181
	253.8–321	0.231	−134.9	0.5	−0.166
	321–788	−0.031	21.91	−0.02	0.007
GCSN1	23.8–86	0.0429	−10.95	−0.15	0.006
	86–223	0.007	1.851	−0.01	0.139
	223–262	−1.266	657.3	−2.10	0.575
	262–325	0.410	−233.6	0.89	−0.284
	325–788	−0.031	29.75	−0.023	0.0212
GCSN2	28.8–227	−0.002	2.029	−0.016	0.005
	226.8–265	−0.986	510.2	−1.4	0.378
	265–331	0.393	−231.7	0.62	−0.213
	331–788	−0.021	14.65	−0.014	0.006
GCSN3	30–87	0.045	−15.04	0.2	−0.024
	87–228	$-1.29 \times 10^{-5}$	1.34	−0.0077	0.004
	228–268	−0.981	509.01	−1.9	0.535
	268–334	0.349	−206.01	0.59	−0.208
	334–787	−0.011	10.78	−0.009	0.008

Crystals contain no water of crystallization as wt. loss at 100 °C is absent [46]. TGA thermograms confirmed thermal stability up to 200.65, 208.27, 217.41, and 221.54 °C for GCS, GCSN1, GCSN2, and GCSN3, which enabled laser applications [47]. Peak temperature  $T_m$  increased from 249.71 to 257.88, 262.28, 264.07 °C on increasing SNPs' concentration.

ESR spectra of powder sample crystals are shown in Figure 9. Anisotropy g-factor for crystals confirmed low symmetry. Spin Hamiltonian parameters g and A tensors revealed rhombic symmetry crystal field around Cu(II) ion split ground state. Degeneracy of ground state energy level is lifted giving static Jahn–Teller distortion [48,49].

Value  $g_{\perp}$  and unpaired electronic state R are given by Equations (7) and (8): [50,51].

$$g_{\perp} = g_x + g_y / 2 \quad (7)$$

$$R = (g_x - g_y) / (g_z - g_x) \quad (8)$$

Table 6 includes g factor, hyper fine constants A, and R for the crystals.

Values A and g factor have no axial symmetry in the crystal lattice (no dynamic Jahn–Teller) [52].  $R = 0.1805, 0.1224, 0.1673, 0.1418$  for GCS, GCSN1, GCSN2, and GCSN3, respectively; less than unity indicated  $d_{x^2-y^2}$  ground state for unpaired electron [53].

$A_x, A_y$  equals  $A_{\parallel}$  are lower than  $A_z$ ; g-parallel is greater than perpendicular  $g_{\perp}$  and confirmed  $d_{x^2-y^2}$  ground state [52].  $g_{\parallel}$  value is less than 2.3, indicating strong covalent copper–glycine bond [54].

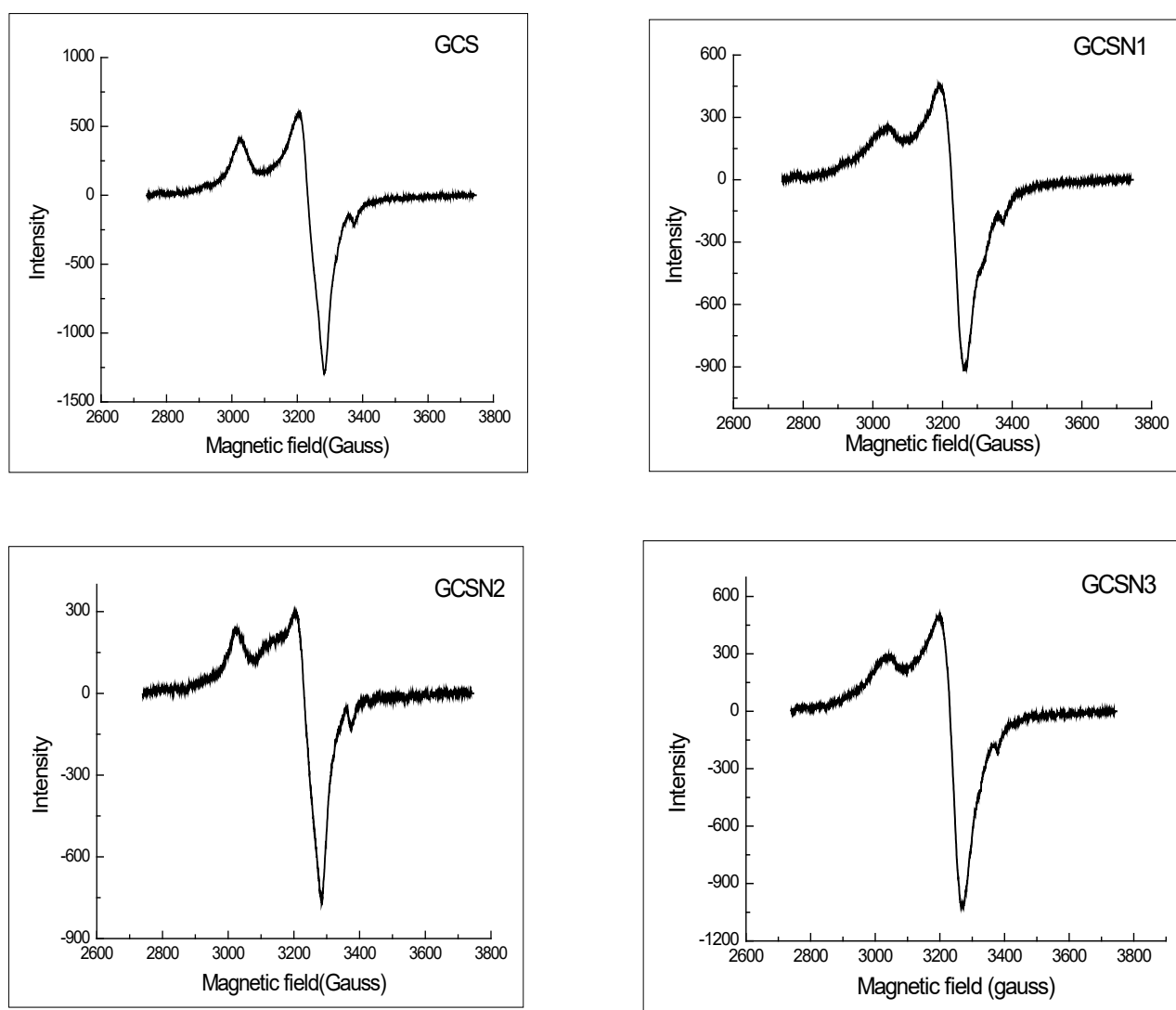
$\sigma$  bond coefficient or covalence parameter  $\alpha^2$  for unpaired electron density on Cu(II) ion is estimated using Equation (9) [54,55], Table 7.

$$\alpha^2 = A_{\parallel} / 0.036 + (g_{\parallel} - 2.0023) + 3/7(g_{\perp} - 2.0023) + 0.04 \quad (9)$$

where

$$A_{\parallel} = A_z, A_{\perp} = (A_x + A_y) / 2 \quad (10)$$

The parameter  $\alpha^2$  is less than unity, indicating covalent bonding between Cu(II) and glycine ligand [56]. SNPs decreased covalence parameter ( $\alpha^2$ ) of glycine–Cu(II) bond, except that GCSN2 showed abnormally high  $\alpha^2$ , which confirmed its highest electrical susceptibility.



**Figure 9.** Powder ESR spectra of Cu(II) ions in crystals.

**Table 6.** g values and A (mT) matrices for Cu(II) ion in crystals.

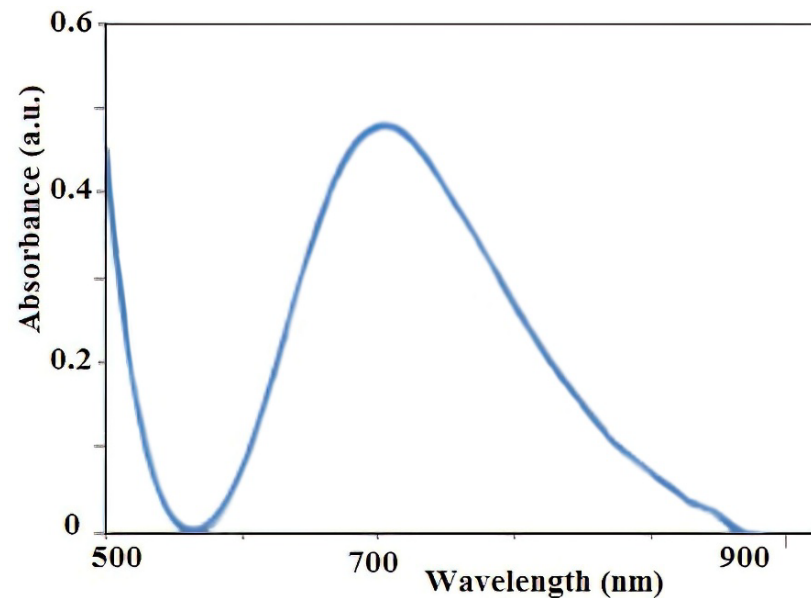
Crystal	$g_x$	$g_y$	$g_z$	$A_x$	$A_y$	$A_z$	$g_{  }$	$g_{\perp}$	R
GCS	2.0859	2.0531	2.2676	13.309	10.872	18.183	2.2676	2.0691	0.1805
GCSN1	2.0921	2.0705	2.2685	12.548	13.949	14.229	2.2685	2.0813	0.1224
GCSN2	2.0883	2.0551	2.2868	11.154	16.183	17.871	2.2868	2.0717	0.1673
GCSN3	2.0887	2.0637	2.265	13.27	12.465	15.965	2.265	2.0762	0.1418

**Table 7.** EPR parameters.

Crystals	$A_{\perp}$ (G)	$A_{  }$ (G)	$\alpha^2$	Crystals	$A_{\perp}$ (G)	$A_{  }$ (G)	$\alpha^2$
GCS	120.908	181.831	0.8392	GCSN2	136.685	178.707	0.8507
GCSN1	132.491	142.299	0.7353	GCSN3	128.719	159.649	0.7778

Figure 10 showed UV of thin film coating of GCSN3 on the aluminum (Al) foil sample. An aqueous solution of GCSN3 was evaporated under ultra-high vacuum conditions onto the Al foil where carboxylate COOH of glycine zwitterion amino acid is chemically adsorbed on the aluminum surface. Absorption at long  $\lambda$  900–1100 nm for this crystal near

IR region indicates absorption of thermal energy of IR radiation. Phonon bands at 900 nm originate from vibrational modes of harmonic and unharmonic oscillators in the crystal lattice. Absorbed IR radiation causes thermal vibrations of atoms or molecules and creates thermal phonon waves that propagate in the crystal lattice, dissipating thermal IR energy.



**Figure 10.** UV-Vis. absorbance spectra of GCSN3 thin film coated on Al foil.

Absorptivity of GCSN3 near the IR region of electromagnetic radiation indicated that crystals can shield thermal heat of IR radiation on the coating as dispersed thin film on alumetal.

The sun provides thousands  $\text{W}\cdot\text{m}^{-2}$  energy on the earth's surface daily. Total solar energy in the upper atmosphere contains 50% IR radiation, 40% Vis. Light, and 10% UV radiation. IR radiation causes vibrations that heats earth's surface [57]. Attenuation of thermal energy can be achieved by painting glass windows with these blue color crystals transparent to UV radiation, filtering, and that dissipates IR radiation.

High thermal conductivity of crystals equals: 1.10, 1.21, 1.54, and  $1.6 \text{ W}\cdot\text{m}^{-1}\text{K}^{-1}$  for GCS, GCSN1, GCSN2, and GCSN3 confirmed rapid attenuation of many incident EM waves by dielectric components and rapidly dissipated as heat.

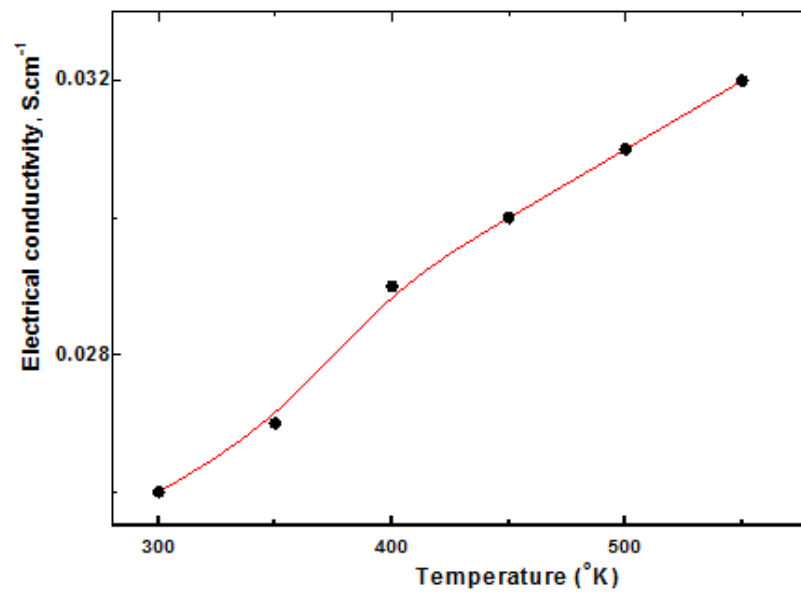
Figure 11 showed electrical conductivity of GCSN3 increased on heating as a typical semiconductor behavior due to thermally activated charge carriers' mobility [58].

Impedance plots confirmed super conductivity. A plateau region at low 0.1Hz frequency region represents total conductivity of grain boundary. A high-frequency region at 100 kHz represents the contribution of grains to total conductivity. An intermediate frequency region at 1 kHz is due to charges trapped between grain boundaries and grains [15]. AC conductivity confirmed the dielectric nature of a single crystal can dissipate heat rapidly. The high-frequency dielectric constant is 4.49.

Dielectric study of GCSN3 crystal response of charges to applied electric field showed dielectric constants at 100 Hz, 1 kHz, 10 kHz, and 100 kHz and a temperature range of 200–550 K. Dielectric constant was calculated using equation [59]:

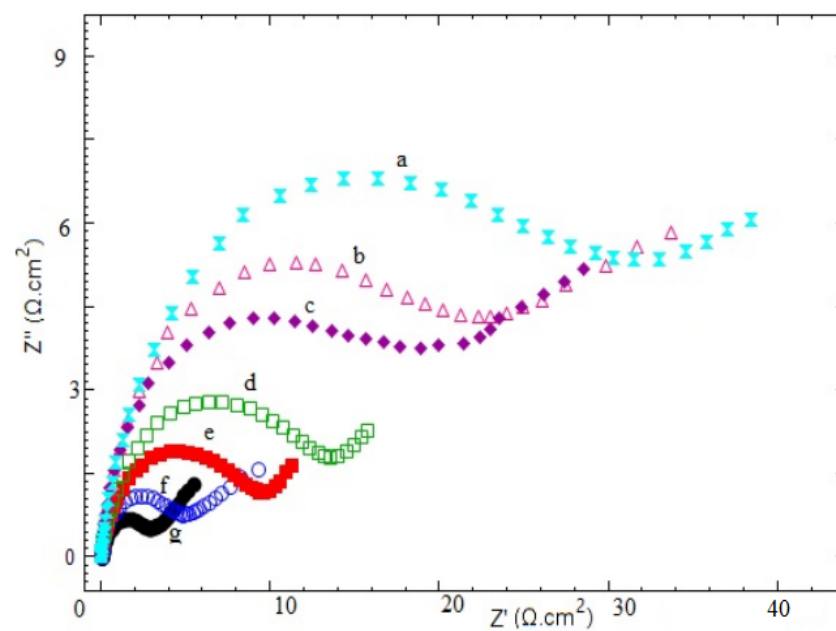
$$\varepsilon = \frac{Cd}{\varepsilon_0 A} \quad (11)$$

where  $\varepsilon_0$  is free space permittivity, C and d are capacitance and thickness of pellet, and A is electrode area.



**Figure 11.** Variation of electrical conductivity of GCSN3 on heating.

Figure 12 showed variation of impedance of GCSN3 with applied frequency.

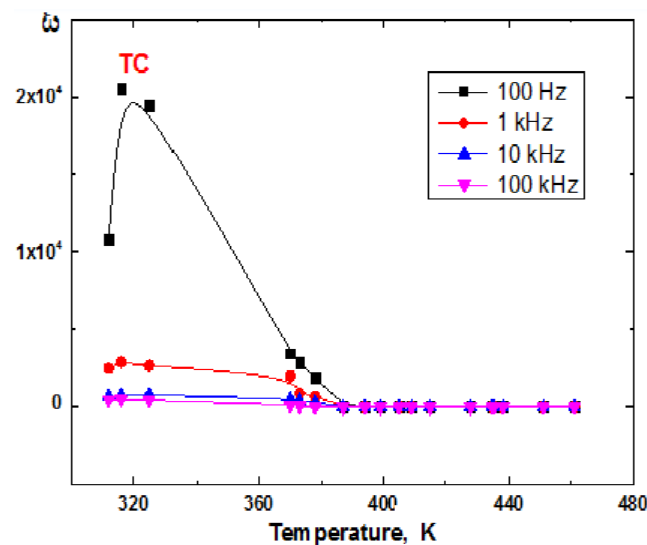


**Figure 12.** Charge transfer resistance of GCSN3 as a function of applied frequency and absolute temperature (K): (a) 200, (b) 300, (c) 350, (d) 400, (e) 450, (f) 500, (g) 550.

Real  $\epsilon'$  and imaginary  $\epsilon''$  components of  $\epsilon$  represented equals [60]:

$$\epsilon' = |\epsilon| \cos\theta, \epsilon'' = |\epsilon| \sin\theta \quad (12)$$

Figure 13 showed  $\epsilon'$  varied with temperature at a different frequency and decreased with increasing frequency, indicating an ability to dissipate incident IR radiation.  $\epsilon'$  decreased until it reached glass transition  $T_g$  at 380 K, then became limited up to 470 K. Peak at Curie TC represented phase transition from ferroelectric to paraelectric behavior.



**Figure 13.** Temperature dependent of the dielectric constant for (GCSN3) single crystal.

#### 4. Conclusions

SNPs dopant increased thermal stability of gamma and alpha glycine single crystals. Optical absorption studies revealed that cut-off wavelengths are 287, 283, 276, and 280 nm and optical band gap energy 4.58, 4.61, 4.65, and 4.67 eV for GCS, GCSN1, GCSN2, and GCSN3 single crystals, respectively. SNPs increased band gaps of crystals, hence transparency to UV radiation. AC electrical conductivity of the thin film sample of perfect crystal increased to 0.03 Siemens/cm. High thermal conductivity,  $\text{W}\cdot\text{m}^{-1}\text{K}^{-1}$  in range 1.10–1.6, confirmed efficient radiation attenuation by rapid heat dissipation due to dielectric properties of single crystals. Single crystals could be used to shield and dissipate thermal heat of IR radiation. AC confirmed the dielectric component and increased on heating due to thermal activation of charge carriers.

**Supplementary Materials:** The following supporting information can be downloaded at: <https://www.mdpi.com/article/10.3390/applnano4020007/s1>, Figure S1. Mass spectra of the single crystals (a) GCS, (b) GCSN1, (c) GCSN2 and (d) GCSN3; Figure S2. FTIR spectra of the single crystals GCS, GCSN1, GCSN2 and GCSN3; Figure S3. Plot extinction coefficient ( $k$ ) (a) and reflectance versus (b) photon energy (eV) of GCS, GCSN1, GCSN2, and GCSN3; Figure S4. DSC curves for GCS crystals: a. specific heat capacity versus  $T$ , b.  $C_p/T$  versus  $T^2$ ; Figure S5. DSC curves for GCSN1 crystal: a. specific heat capacity versus  $T$ , b.  $C_p/T$  versus  $T^2$ ; Figure S6. DSC curves for GCSN2 crystal: a. specific heat capacity versus  $T$ , b.  $C_p/T$  versus  $T^2$ ; Figure S7. DSC curves for GCSN3 crystal: a. Specific heat capacity versus  $T$ , b.  $C_p/T$  versus  $T^2$ ; Table S1. Indexed PXRD patterns of GCS, GCSN1, GCSN2 and GCSN3 single crystals.

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